

# Solving the sextic by iteration: A study in complex geometry and dynamics

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## 1 Introduction

### 1.1 Overview

Recently, [Doyle and McMullen 1989] devised an iterative solution to the fifth degree polynomial. At the method's core is a rational mapping  $f$  of  $\mathbf{CP}^1$  with the icosahedral symmetry of a general quintic. Algebraically, this means that  $f$  commutes with a group of Möbius transformations that is isomorphic to the alternating group  $\mathcal{A}_5$ . Moreover, this  $\mathcal{A}_5$ -equivariant possesses *nice* dynamics: for almost any initial point  $a \in \mathbf{CP}^1$ , the sequence of iterates  $f^k(a)$  converges to one of the periodic cycles that comprise an icosahedral orbit.<sup>1</sup> This breaking of  $\mathcal{A}_5$ -symmetry provides for a *reliable* or *generally-convergent* quintic-solving algorithm: with almost any fifth-degree equation, associate a rational mapping that has nice dynamics and whose attractor consists of a single orbit from which one computes a root.

An algorithm that solves the sixth-degree equation calls for a dynamical system with  $\mathcal{S}_6$  or  $\mathcal{A}_6$  symmetry. Since neither  $\mathcal{S}_6$  nor  $\mathcal{A}_6$  acts on  $\mathbf{CP}^1$ , attention turns to higher dimensions. Acting on  $\mathbf{CP}^2$  is an  $\mathcal{A}_6$ -isomorphic group of projective transformations found by Valentiner in the late nineteenth century. The present work exploits this 2-dimensional  $\mathcal{A}_6$  “soccer ball” in order to discover a “Valentiner-symmetric” rational mapping of  $\mathbf{CP}^2$  whose dynamics *experimentally appear* to be nice in the above sense—transferred to the  $\mathbf{CP}^2$  setting. This map provides the central feature of a conjecturally-reliable sextic-solving algorithm analogous to that employed in the quintic case.

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<sup>0</sup>For insight and inspiration, I would like thank Peter Doyle and Curt McMullen, the founders of this project.

<sup>1</sup>For a geometric description see [Doyle and McMullen 1989, p. 163].

## 1.2 Solving equations by iteration

For  $n \leq 4$ , the symmetric groups  $\mathcal{S}_n$  act faithfully on  $\mathbf{CP}^1$ . Corresponding to each action is a map whose nice dynamics provides for algorithmic convergence to roots of given  $n$ th-degree equations. For instance, Newton’s method provides a direct iterative solution to quadratic polynomials, but, due to a lack of symmetry, not to higher degree equations. My interests here are the geometric and dynamical properties of complex projective mappings rather than numerical estimates.

The search for elegant complex geometry and dynamics continues into degree five where  $\mathcal{A}_5$  is the appropriate group, since  $\mathcal{S}_5$  fails to act on the sphere. This reduction in the galois group requires the extraction of the square root of a polynomial’s discriminant. Such root-taking is itself the result of a reliable iteration, namely, Newton’s method. In practical terms, the Doyle-McMullen algorithm solves a family of fifth-degree *resolvents* the members of which possess  $\mathcal{A}_5$  symmetry. A map with icosahedral symmetry and nice dynamics plays the leading role.

Pressing on to the sixth-degree leads to the 2-dimensional  $\mathcal{A}_6$  action of the Valentiner group  $\mathcal{V}$ . Here, the problem shifts to one of finding a nice  $\mathcal{V}$ -symmetric mapping of  $\mathbf{CP}^2$  from whose attractor one calculates a given sextic’s root.<sup>2</sup> Providing the overall framework is the 2-dimensional  $\mathcal{A}_6$  analogue of the icosahedron.

## 1.3 Proofs and Computations

At the moment, many of this work’s results have only computational support. As such, I call them “Facts”. Furthermore, its conjectural nature calls for a deeper understanding of Valentiner geometry and dynamics. As the theory of complex dynamics in several dimensions develops more sophisticated weaponry, the barricades to understanding might become assailable. For now, I hope that these discoveries provide a stimulus to such development.

## 2 Valentiner’s Group: The $\mathcal{A}_6$ Action on $\mathbf{CP}^2$

If ...any equation  $f(x) = 0$  is given, we will investigate what is the smallest number of variables with which we can construct a group of linear substitutions which is isomorphic with the Galois group of  $f(x) = 0$ . [Klein 1956, p. 138]

In the wake of the mid-nineteenth century abstractionist turn in mathematics the theory of group representations began to emerge. Part of the concrete yield from work on symmetric groups was Valentiner’s discovery[Valentiner 1889] of a complex projective group that is isomorphic to  $\mathcal{A}_6$ —the alternating group of six things. Shortly thereafter [Wiman 1895] explored some of the geometric and invariant structure determined by this action on the complex projective plane. A more thorough exposition appeared in [Fricke 1926].

Here, I take a new approach to the generation of this “Valentiner” group and then explore some of its rich combinatorial geometry. The core of this work involves the development of a combinatorially sensitive description of the basic geometric structures. In so doing, I reproduce

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<sup>2</sup>The solution-procedure follows that of the quintic algorithm. See Section 4.

some of the Wiman and Fricke results. The study culminates in an account of the system of Valentiner-invariant polynomials and, thereby, lays the algebraic foundation for constructing Valentiner-symmetric mappings of  $\mathbf{CP}^2$ .

## 2.1 Basics of $\mathcal{A}_6$

The following table indicates the non-identity elements of  $\mathcal{A}_6$  by order, cycle-structure, and cardinality.

Order	Structure	Number
2	$(ab)(cd)$	$45 = \frac{1}{2} \binom{6}{4} \cdot 3!$
3	$(abc)$	$40 = \binom{6}{3} \cdot 2$
3	$(abc)(def)$	$40 = \binom{6}{3} \cdot 2$
4	$(abcd)(ef)$	$90 = \binom{6}{4} \cdot 3!$
5	$(abcde)$	$144 = 6 \cdot 4!$

Sitting inside  $\mathcal{A}_6$  are twelve versions of  $\mathcal{A}_5$  that decompose into two conjugate systems of six:

- 1) the stabilizers  $\text{Stab}\{k\}$  of one thing
- 2) the permutations of the six pairs of antipodal icosahedral vertices under rotation.

Acting by conjugation,  $\mathcal{A}_6$  permutes each of the two systems individually. A given  $\mathcal{A}_5$  subgroup fixes itself set-wise and permutes its five conjugates according to the rotational icosahedral group's action on the five cubes found in the icosahedron. Meanwhile, the other system of six subgroups undergo the permutations of the six pairs of antipodal vertices. Consequently, the intersection of two  $\mathcal{A}_5$  subgroups in the *same* system is isomorphic to  $\mathcal{A}_4$ —the tetrahedral rotations—while two in *different* systems give a dihedral group  $\mathcal{D}_5$ .

## 2.2 Generating the Valentiner group

An  $\mathcal{A}_5$  subgroup of  $\mathcal{A}_6$ , say  $\text{Stab}\{1\}$ , extends to  $\mathcal{A}_6$  by addition of the generator  $(12)(3456)$ . Furthermore, this order-four element generates an  $\mathcal{S}_4$  over the  $\mathcal{A}_4$  subgroup

$$\langle (35)(46), (456) \rangle \subset \text{Stab}\{1\}.$$

This structure suggests a method for producing an  $\mathcal{A}_6$ -isomorphic group  $\mathcal{V}$  in  $\text{PGL}_3(\mathbf{C})$ :

- 1) take a tetrahedral subgroup  $\mathcal{T}$  of an icosahedral group  $\mathcal{I}$ ;
- 2) by addition of an order-4 transformation  $Q$ , extend  $\mathcal{T}$  to an octahedral group  $\mathcal{O} = \langle \mathcal{T}, Q \rangle$ ;
- 3) generate  $\mathcal{V} = \langle \mathcal{I}, Q \rangle \simeq \mathcal{A}_6$ .

The 15 pairs of antipodal edges of the standard icosahedron decompose into five triples such that three lines joining antipodal edge-midpoints are mutually perpendicular. Stabilizing each such triple is one of the five tetrahedral subgroups of the icosahedral group. Alternatively, the lines in such a triple correspond to the two-fold rotational axes of a tetrahedron whose four vertices are face-centers of the icosahedron. With such a triple of lines as coordinate axes<sup>3</sup> in  $\{x_1, x_2, x_3\}$ , the points<sup>4</sup>

$$A = \{[1, 1, 1], [-1, -1, 1], [1, -1, -1], [-1, 1, -1]\}$$

constitute a set of tetrahedral vertices. The corresponding tetrahedral group

$$\mathcal{T} = \text{Stab}(A)$$

consists of the identity and the 11 *orthogonal* transformations:

$$\begin{aligned} Z_1 &= \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix} & Z_2 &= \begin{pmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -1 \end{pmatrix} & Z_3 &= \begin{pmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \\ T_1 &= \begin{pmatrix} 0 & 0 & 1 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix} & T_1^2 &= \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix} \\ T_2 &= \begin{pmatrix} 0 & 0 & -1 \\ 1 & 0 & 0 \\ 0 & -1 & 0 \end{pmatrix} & T_2^2 &= \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & -1 \\ -1 & 0 & 0 \end{pmatrix} \\ T_3 &= \begin{pmatrix} 0 & 0 & -1 \\ -1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix} & T_3^2 &= \begin{pmatrix} 0 & -1 & 0 \\ 0 & 0 & 1 \\ -1 & 0 & 0 \end{pmatrix} \\ T_4 &= \begin{pmatrix} 0 & 0 & 1 \\ -1 & 0 & 0 \\ 0 & -1 & 0 \end{pmatrix} & T_4^2 &= \begin{pmatrix} 0 & -1 & 0 \\ 0 & 0 & -1 \\ 1 & 0 & 0 \end{pmatrix}. \end{aligned}$$

Being orthogonal,  $\mathcal{T}$  preserves the quadratic form

$$C(x) = x_1^2 + x_2^2 + x_3^2$$

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<sup>3</sup>See Figure 1.

<sup>4</sup>Square brackets indicate points in projective space.

and hence, the conic  $\mathcal{C} = \{C = 0\}$  in  $\mathbf{CP}^2$ . This “tetrahedral” conic contains two sets of orbits<sup>5</sup> of size four: two of the points fixed projectively by each of the three-fold  $T_k$ . The third fixed point is one of the elements of the set  $A$  above. With  $\rho = e^{2\pi i/3}$  the respective points are

$$\begin{aligned} v_1 &= [\rho, \rho^2, 1] & v_{\overline{1}} &= [\rho^2, \rho, 1] \\ v_2 &= [-\rho, -\rho^2, 1] & v_{\overline{2}} &= [-\rho^2, -\rho, 1] \\ v_3 &= [-\rho, \rho^2, 1] & v_{\overline{3}} &= [-\rho^2, \rho, 1] \\ v_4 &= [\rho, -\rho^2, 1] & v_{\overline{4}} &= [\rho^2, -\rho, 1]. \end{aligned}$$

The “barred” notation  $v_{\overline{a}}$  derives from Fricke, being suggested by an antiholomorphic relationship between the two systems of tetrahedra. Indeed, in the  $x$  coordinates chosen above, the conjugation map  $x \rightarrow \overline{x}$  exchanges a tetrahedron and its “antipode”.

The order-four transformation

$$Q = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & \rho^2 \\ 0 & -\rho & 0 \end{pmatrix}$$

cyclically permutes the  $v_a$  but not the  $v_{\overline{a}}$  while

$$\overline{Q} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & \rho \\ 0 & -\rho^2 & 0 \end{pmatrix}$$

cyclically permutes the  $v_{\overline{a}}$  but not the  $v_a$ . Also,  $Q^2 = \overline{Q}^2 = Z_1$ . Since  $\mathrm{PGL}_3(\mathbf{C})$  is four-times transitive,  $Q$  and  $R$  are the unique such projective transformations. Accordingly, the groups  $\mathcal{O} = \langle \mathcal{T}, Q \rangle$  and  $\overline{\mathcal{O}} = \langle \mathcal{T}, \overline{Q} \rangle$  are octahedral, i.e., isomorphic to  $\mathcal{S}_4$ .

Extending  $\mathcal{T}$  to an icosahedral group  $\mathcal{I}$  requires a projective transformation  $P$  of order five that preserves  $\mathcal{C}$  and point-wise fixes a pair of antipodal icosahedral vertices. One way of producing such a  $P$  is to “turn” the icosahedron<sup>6</sup> of Figure 1 so that a pair of antipodal vertices corresponds to the point  $[0, 1, 0]$ , i.e., to the “affine” points  $(0, \pm 1, 0)$ . In these “icosahedral” coordinates  $\{u_1, u_2, u_3\}$  the desired transformation of order five is

$$P_u = \begin{pmatrix} \cos \frac{2\pi}{5} & 0 & -\sin \frac{2\pi}{5} \\ 0 & 1 & 0 \\ \sin \frac{2\pi}{5} & 0 & \cos \frac{2\pi}{5} \end{pmatrix}.$$

The change of basis from octahedral to icosahedral coordinates is

$$u = Ax = \begin{pmatrix} c & -s & 0 \\ s & c & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix}$$

<sup>5</sup>This is a manifestation of the pairs of “antipodal” tetrahedra situated in the icosahedron.

<sup>6</sup>See Figure 2.

where

$$c = \sqrt{\frac{5 + \sqrt{5}}{10}} \quad s = \sqrt{\frac{5 - \sqrt{5}}{10}}.$$

Thus, in icosahedral coordinates, the conic form preserved by  $\mathcal{I}$  is

$$C(u) = u_1^2 + u_2^2 + u_3^2$$

while  $P$  has the expression

$$P_x = A^{-1}P_uA = \frac{1}{2} \begin{pmatrix} 1 & \tau^{-1} & -\tau \\ \tau^{-1} & \tau & 1 \\ \tau & -1 & \tau^{-1} \end{pmatrix}, \quad \tau = \frac{1 + \sqrt{5}}{2}.$$

Finally,  $\mathcal{I} = \langle \mathcal{T}, P \rangle$ . This produces two Valentiner groups distinguished by chirality:

$$\mathcal{V} = \langle \mathcal{I}, Q \rangle \quad \overline{\mathcal{V}} = \langle \mathcal{I}, \overline{Q} \rangle.$$

Use of the terms ‘octahedral coordinates’ and ‘icosahedral coordinates’ follows that of [Fricke 1926, pp. 263ff]. His system of octahedral generators are nearly those above. In his icosahedral coordinates, the conic form is

$$C_{Fricke}(z) = z_1 z_3 + z_2^2.$$

The change of coordinates  $B$  that yields  $C(Bz) = C_{Fricke}(z)$  is

$$B = \begin{pmatrix} 1 & 0 & i \\ 0 & 1 & 0 \\ 1 & 0 & -i \end{pmatrix}.$$

With  $\epsilon = e^{2\pi i/5}$  and  $Z = Z_2$ , the generators  $Z_u = AZ_xA^{-1}$  and  $P_u = AP_xA^{-1}$  become his icosahedral generators [Fricke 1926, p. 263]  $T$  and  $S$ :

$$\begin{aligned} Z_z &= BZ_uB^{-1} = \frac{1}{\sqrt{5}} \begin{pmatrix} -\frac{1+\sqrt{5}}{2} & 2 & \frac{-1+\sqrt{5}}{2} \\ 1 & 1 & 1 \\ \frac{-1+\sqrt{5}}{2} & 2 & -\frac{1+\sqrt{5}}{2} \end{pmatrix} \\ &= \frac{1}{\sqrt{5}} \begin{pmatrix} \epsilon^2 + \epsilon^3 & 2 & \epsilon + \epsilon^4 \\ 1 & 1 & 1 \\ \epsilon + \epsilon^4 & 2 & \epsilon^2 + \epsilon^3 \end{pmatrix} \\ P_z &= BP_uB^{-1} = \begin{pmatrix} \epsilon & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & \epsilon^4 \end{pmatrix}. \end{aligned}$$

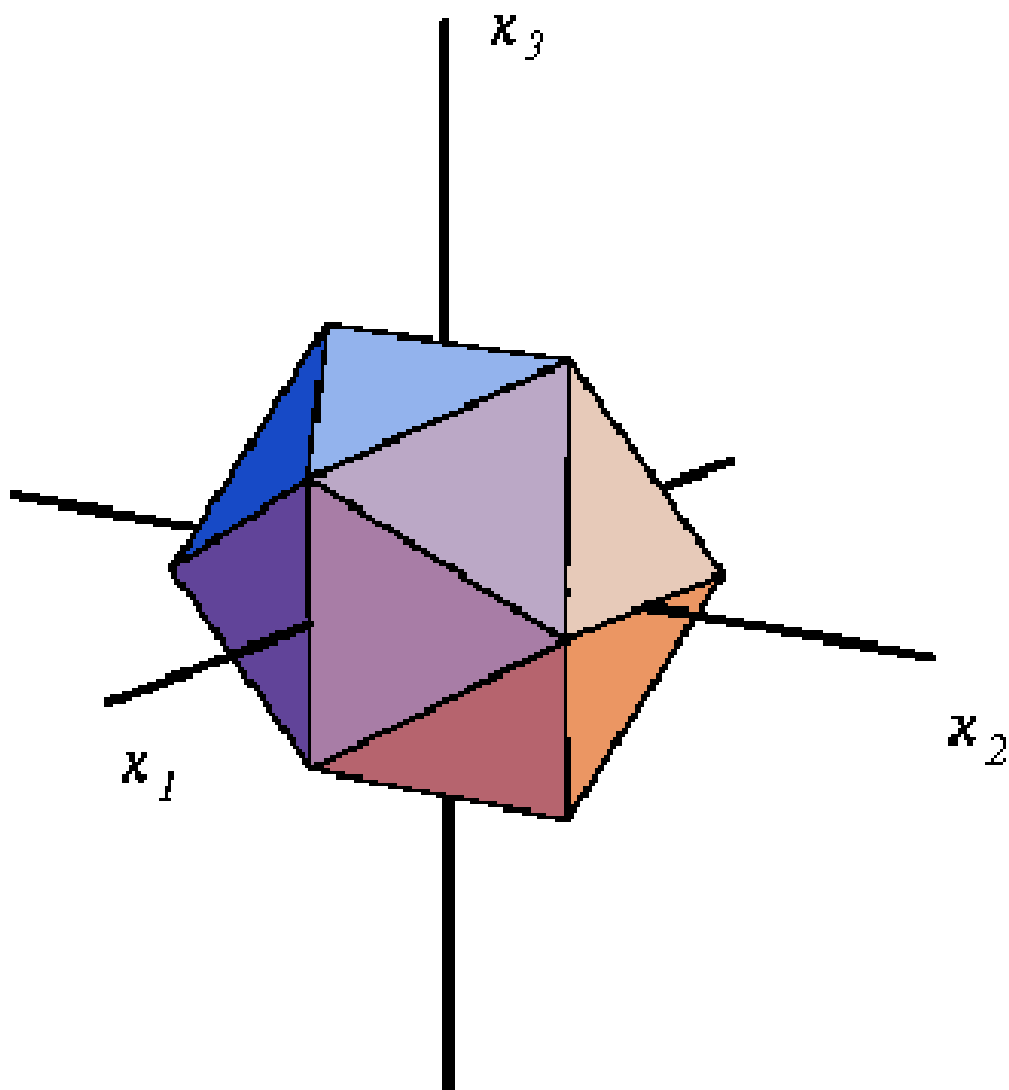


Figure 1: The icosahedron in “octahedral” coordinates.

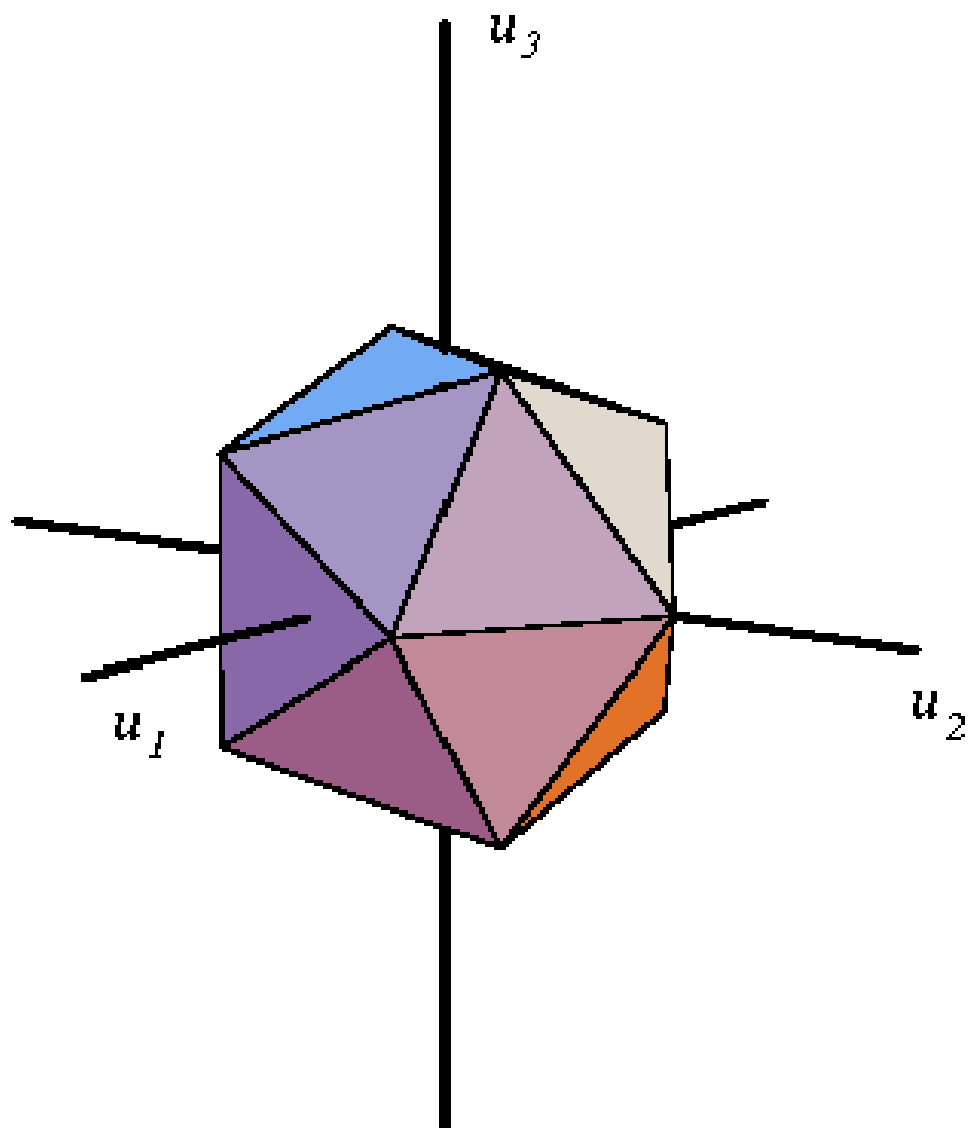


Figure 2: The icosahedron in “icosahedral” coordinates.



## 2.3 Valentiner geometry

### 2.3.1 Icosahedral conics

Since a ternary icosahedral group has an orthogonal representation on  $\mathbf{R}^3$ , its complexification stabilizes the quadratic form

$$C(x) = x_1^2 + x_2^2 + x_3^2$$

over  $\mathbf{C}$  and hence, the conic  $\mathcal{C} = \{C(x) = 0\}$  in  $\mathbf{CP}^2$ . For reasons that will become clear, the two systems of six icosahedral groups  $\mathcal{I}_{\bar{k}}$  and  $\mathcal{I}_k$  in  $\mathcal{V}$  receive the designations “barred” and “unbarred”. Corresponding to each of the twelve icosahedral subgroups of  $\mathcal{V}$  is a quadratic form  $C_{\bar{a}}$  or  $C_a$  and conic  $\mathcal{C}_{\bar{a}}$  or  $\mathcal{C}_a$  that  $\mathcal{I}_{\bar{a}}$  or  $\mathcal{I}_a$  preserves. Thereby, each such conic possesses the structure of an icosahedron.

Call the form above  $C_{\bar{1}}$ . The action of  $\mathcal{V}$  produces the five remaining forms. With  $\eta = (3 + \sqrt{15}i)/4$ ,

$$\begin{aligned} C_{\bar{1}}(x) &= x_1^2 + x_2^2 + x_3^2 \\ C_{\bar{2}}(x) &= C_{\bar{1}}(Q^{-1}x) \\ &= x_1^2 + \rho^2 x_2^2 + \rho x_3^2 \\ C_{\bar{3}}(x) &= C_{\bar{2}}(P^{-4}x) \\ &= \eta \left[ \frac{\eta}{3} (x_1^2 + \rho x_2^2 + \rho^2 x_3^2) + (\rho^2 x_1 x_2 + \rho x_1 x_3 - x_2 x_3) \right] \\ C_{\bar{4}}(x) &= C_{\bar{2}}(P^{-3}x) \\ &= \eta \left[ \frac{\eta}{3} (\rho x_1^2 + \rho^2 x_2^2 + x_3^2) + (-x_1 x_2 + \rho^2 x_1 x_3 + \rho x_2 x_3) \right] \\ C_{\bar{5}}(x) &= C_{\bar{2}}(P^{-2}x) \\ &= \eta \left[ \frac{\eta}{3} (\rho x_1^2 + \rho^2 x_2^2 + x_3^2) - (x_1 x_2 + \rho^2 x_1 x_3 + \rho x_2 x_3) \right] \\ C_{\bar{6}}(x) &= C_{\bar{2}}(P^{-1}x) \\ &= \eta \left[ \frac{\eta}{3} (x_1^2 + \rho x_2^2 + \rho^2 x_3^2) + (\rho^2 x_1 x_2 - \rho x_1 x_3 + x_2 x_3) \right]. \end{aligned}$$

As for the action of  $\mathcal{V}$  on these:

$$\begin{aligned} P : C_{\bar{1}} &\leftrightarrow C_{\bar{1}} & C_{\bar{2}} &\rightarrow C_{\bar{6}} \rightarrow C_{\bar{5}} \rightarrow C_{\bar{4}} \rightarrow C_{\bar{3}} \rightarrow C_{\bar{2}} \\ Z = Z_2 : C_{\bar{1}} &\leftrightarrow C_{\bar{1}} & C_{\bar{2}} &\leftrightarrow C_{\bar{2}} & C_{\bar{3}} &\leftrightarrow \rho^2 C_{\bar{4}} & C_{\bar{5}} &\leftrightarrow \rho C_{\bar{6}} \\ Q : C_{\bar{1}} &\leftrightarrow C_{\bar{2}} & C_{\bar{3}} &\rightarrow C_{\bar{6}} \rightarrow \rho^2 C_{\bar{5}} \rightarrow \rho^2 C_{\bar{4}} \rightarrow C_{\bar{3}}. \end{aligned}$$

Direct calculation yields

**Proposition 1** *The quadratic forms  $C_{\bar{m}}$  are linearly independent and so span the six dimensional space of ternary quadratic forms. In particular, an unbarred conic form  $C_a$  is a linear combination of the  $C_{\bar{m}}$  that is invariant under the icosahedral group  $\mathcal{I}_a$ .*

To be specific, the “5-cycle”  $P$  belongs to  $\mathcal{I}_{\overline{1}}$  and to an unbarred icosahedral group—say  $\mathcal{I}_3$  so that the indexing agrees with that of Fricke. (Recall that the intersection of two non-conjugate  $\mathcal{A}_5$  subgroups of  $\mathcal{A}_6$  is a  $\mathcal{D}_5$ .) Relative or *projective* invariance under  $P$  requires  $C_3$  to take the form

$$C_3 = \alpha C_{\overline{1}} + C_{\overline{2}} + C_{\overline{3}} + C_{\overline{4}} + C_{\overline{5}} + C_{\overline{6}}.$$

To determine the constant, apply to the  $C_{\overline{m}}$  an element  $T$  of  $\mathcal{I}_3$  that does not belong to  $\mathcal{I}_{\overline{1}}$ , e.g., one of the 20 elements of order three in  $\mathcal{I}_3$ . Recalling the association between  $\mathcal{I}_3$  and the permutations of the six pairs of antipodal icosahedral vertices labeled according to the action of  $P$ , such a transformation corresponds to a double 3-cycle of the form  $(abc)(def)$ . Specifically, the permutation  $(\overline{164})(\overline{235})$ —to be used below—corresponds to an element of  $\mathcal{I}_3$ .

The action of the generators on the *conics*  $C_{\overline{m}}$  is given by the permutation of the indices:

$$\begin{aligned} P : & (\overline{26543}) \\ Z : & (\overline{34})(\overline{56}) \\ Q : & (\overline{12})(\overline{3654}). \end{aligned}$$

Computation in  $\mathcal{A}_6$  yields the correspondence

$$T = QP^2QPQ^3 : (\overline{164})(\overline{235}).$$

Moreover, the action on the conic forms is

$$\begin{aligned} T : \quad C_{\overline{1}} & \rightarrow C_{\overline{6}} \rightarrow \rho C_{\overline{4}} \rightarrow C_{\overline{1}} \\ C_{\overline{2}} & \rightarrow \rho C_{\overline{3}} \rightarrow \rho^2 C_{\overline{5}} \rightarrow C_{\overline{2}} \end{aligned}$$

so that

$$C_3(T^{-1}x) = \rho^2 C_{\overline{1}} + \rho (C_{\overline{2}} + C_{\overline{3}} + C_{\overline{4}} + C_{\overline{5}}) + \alpha C_{\overline{6}}.$$

Projective invariance under  $T$  of the conic  $\mathcal{C}_3 = \{C_3 = 0\}$  requires  $\alpha = \rho$ . Accordingly,

$$\begin{aligned} C_3(T^{-1}x) &= \rho^2 C_{\overline{1}} + \rho (C_{\overline{2}} + C_{\overline{3}} + C_{\overline{4}} + C_{\overline{5}} + C_{\overline{6}}) \\ &= \rho C_3(x). \end{aligned}$$

Just as the barred forms stem from  $C_{\overline{1}}$ , the remaining unbarred conic forms<sup>7</sup> arise from  $C_3$ :

$$\begin{aligned}
C_1(x) &= C_2(P^{-1}x) \\
&= C_{\overline{1}} + \rho C_{\overline{2}} + C_{\overline{3}} + \rho^2 C_{\overline{4}} + C_{\overline{5}} + \rho C_{\overline{6}} \\
&= -\eta \left( \rho^2 x_1^2 + \frac{4}{3} \eta^2 x_2^2 + \rho x_3^2 + 2\rho(\rho-1)x_1x_3 \right) \\
C_2(x) &= C_3(Q^{-1}x) \\
&= C_{\overline{1}} + \rho C_{\overline{2}} + \rho C_{\overline{3}} + C_{\overline{4}} + \rho^2 C_{\overline{5}} + C_{\overline{6}} \\
&= -\eta \left( \rho^2 x_1^2 + \frac{4}{3} \eta^2 x_2^2 + \rho x_3^2 - 2\rho(\rho-1)x_1x_3 \right) \\
C_3(x) &= \rho C_{\overline{1}} + C_{\overline{2}} + C_{\overline{3}} + C_{\overline{4}} + C_{\overline{5}} + C_{\overline{6}} \\
&= -\eta \rho \left( \rho x_1^2 + \rho^2 x_2^2 + \frac{4}{3} \eta^2 x_3^2 + 2\rho(\rho-1)x_1x_2 \right) \\
C_4(x) &= C_2(P^{-3}x) \\
&= C_{\overline{1}} + \rho^2 C_{\overline{2}} + C_{\overline{3}} + \rho C_{\overline{4}} + \rho C_{\overline{5}} + C_{\overline{6}} \\
&= -\eta \left( \rho x_1^2 + \rho^2 x_2^2 + \frac{4}{3} \eta^2 x_3^2 - 2\rho(\rho-1)x_1x_2 \right) \\
C_5(x) &= C_2(P^{-2}x) \\
&= C_{\overline{1}} + C_{\overline{2}} + \rho^2 C_{\overline{3}} + C_{\overline{4}} + \rho C_{\overline{5}} + \rho C_{\overline{6}} \\
&= -\eta \left( \frac{4}{3} \eta^2 x_1^2 + \rho x_2^2 + \rho^2 x_3^2 + 2\rho(\rho-1)x_2x_3 \right) \\
C_6(x) &= C_2(P^{-4}x) \\
&= C_{\overline{1}} + C_{\overline{2}} + \rho C_{\overline{3}} + \rho C_{\overline{4}} + C_{\overline{5}} + \rho^2 C_{\overline{6}} \\
&= -\eta \left( \frac{4}{3} \eta^2 x_1^2 + \rho x_2^2 + \rho^2 x_3^2 - 2\rho(\rho-1)x_2x_3 \right)
\end{aligned}$$

Application of  $\mathcal{V}$  yields

$$\begin{aligned}
P : C_3 &\leftrightarrow C_3 & C_1 &\rightarrow C_5 \rightarrow C_4 \rightarrow C_6 \rightarrow C_2 \rightarrow C_1 \\
Z : C_1 &\leftrightarrow C_1 & C_2 &\leftrightarrow C_2 & C_3 &\leftrightarrow \rho C_4 & C_5 &\leftrightarrow C_6 \\
Q : C_5 &\leftrightarrow C_6 & C_1 &\rightarrow C_3 \rightarrow C_2 \rightarrow \rho C_4 \rightarrow C_1.
\end{aligned}$$

### 2.3.2 Antiholomorphic symmetry

The one-dimensional icosahedral group  $\mathcal{G}_{60}$  acts on two sets of five tetrahedra each of which corresponds to a quadruple of points in  $\mathbf{CP}^1$ . However, no element of the group sends the tetrahedra of one set to those of the other. Such an exchange occurs by means of anti-holomorphic maps of degree one. Of these, 15 correspond to reflections through the 15 great circles of reflective icosahedral symmetry; the remaining 45 are the various “odd” compositions of the 15 basic reflections—e.g., the antipodal map. Extending the holomorphic  $\mathcal{G}_{60}$  by such an “anti-involution” produces the group  $\overline{\mathcal{G}}_{120}$  of all 120 symmetries of the icosahedron. The 15 icosahedral

<sup>7</sup>Again, the indices are chosen to agree with Fricke’s labels.

reflections generate this extended group while their even products result in  $\mathcal{G}_{60}$ . In coordinates where one of the great circles corresponds to the real axis, the associated anti-involution is complex conjugation—in homogeneous coordinates:

$$[x_1, x_2] \rightarrow [\overline{x_1}, \overline{x_2}].$$

The Valentiner analogues of the tetrahedra are the two systems of conics. Are there ternary anti-involutions that exchange the barred and unbarred conics? If so, can they take the form

$$[x_1, x_2, x_3] \rightarrow [\overline{x_1}, \overline{x_2}, \overline{x_3}]?$$

To these questions [Fricke 1926, pp. 270-1, 286-9] provides<sup>8</sup> affirmative answers. In the current octahedral coordinates, this “bar-unbar” map is

$$\text{bub}[x_1, x_2, x_3] = [\rho^2 \overline{x_1} - \rho \overline{x_3}, -\rho(\rho + \tau) \overline{x_2}, -\rho \overline{x_1} - \overline{x_3}].$$

As for the action<sup>9</sup> on the conic forms:

$$\begin{aligned} C_1(\text{bub}(x)) &= \alpha \rho^2 C_{\overline{1}}(x) & C_2(\text{bub}(x)) &= \alpha \rho C_{\overline{2}}(x) \\ C_3(\text{bub}(x)) &= \alpha C_{\overline{3}} & C_4(\text{bub}(x)) &= \alpha C_{\overline{4}}(x) \\ C_5(\text{bub}(x)) &= \alpha \rho^2 C_{\overline{5}}(x) & C_6(\text{bub}(x)) &= \alpha \rho C_{\overline{6}}(x) \end{aligned}$$

where  $\alpha = (3 + \sqrt{15}i)/2$ .

**Proposition 2** *The group  $\overline{\mathcal{V}}_{2:360} = \langle \mathcal{V}, \text{bub} \rangle$  is a degree two extension of  $\mathcal{V}$ .*

*Proof.* For  $T \in \mathcal{V}$ ,  $T' = \text{bub} \circ T \circ \text{bub}$  is a projective transformation that permutes the conics within a system. Thereby,  $T'$  belongs to  $\mathcal{V}$ .  $\triangle$

Concerning the form of a bub map, there are coordinate systems in which its expression is conjugation of each coordinate. While interesting in their own right, such coordinates also yield computational benefits. Some of the Valentiner structure suggests a means of achieving this diagonalization. I will take up the topic once the relevant framework is in place.

### 2.3.3 Special orbits and a Valentiner nomenclature

**Intersecting conics.** Some of the special icosahedral points on a conic  $\mathcal{C}_{\overline{a}}$  occur at the intersections of  $\mathcal{C}_{\overline{a}}$  and the other 11 conics.

**Fact 1** *Within a system,  $\mathcal{C}_{\overline{a}}$  meets each  $\mathcal{C}_{\overline{m}}$  ( $\overline{m} \neq \overline{a}$ ) in four tetrahedral points; this gives the 20 face-centers on  $\mathcal{C}_{\overline{a}}$ .*

---

<sup>8</sup>See below for a combinatorial geometric computation of this additional symmetry.

<sup>9</sup>The match between  $\mathcal{C}_{\overline{a}}$  and  $\mathcal{C}_a$  is no accident. Fricke used this map to dub the unbarred conics.

The overall result is a 60 point  $\mathcal{V}$ -orbit  $\mathcal{O}_{\overline{60}}$ . Similarly, the unbarred intersections yield  $\mathcal{O}_{60}$ . Alternatively, each member of  $\mathcal{O}_{\overline{60}}$  (or  $\mathcal{O}_{60}$ ) is a point fixed by one of the 20 barred (or unbarred) cyclic subgroups<sup>10</sup> of order three in  $\mathcal{V}$ .

**Fact 2** *Across systems the intersection of  $\mathcal{C}_{\overline{a}}$  with the  $\mathcal{C}_b$  gives six pairs of antipodal icosahedral vertices  $\{p_{\overline{a}b_1}, p_{\overline{a}b_2}\}$ .*

These total to  $72 = 6 \cdot 12$  points each of which is fixed under one of 36 order-five cyclic groups  $\langle P_{\overline{a}b} \rangle$ . Since an icosahedral group is transitive on its vertices—indeed, some element  $\mathcal{I}_{\overline{a}}$  of order two exchanges  $p_{\overline{a}b_1}$  and  $p_{\overline{a}b_2}$ , the 72 points

$$\mathcal{O}_{72} = \{p_{\overline{a}b_1} | a, b = 1, \dots, 6\} \cup \{p_{\overline{a}b_2} | a, b = 1, \dots, 6\}$$

form a  $\mathcal{V}$ -orbit. A Valentiner exchange of  $p_{\overline{a}b_1}$  and  $p_{\overline{a}b_2}$  also transposes the lines<sup>11</sup>  $\mathcal{L}_{\overline{a}b_2}$  and  $\mathcal{L}_{\overline{a}b_1}$  tangent to  $\mathcal{C}_a$  and  $\mathcal{C}_{\overline{b}}$  at  $p_{\overline{a}b_1}$  and  $p_{\overline{a}b_2}$ . Hence, the intersection of these lines belongs to a 36 point orbit.<sup>12</sup> Each of these<sup>13</sup> “36-points”  $p_{\overline{a}b}$  corresponds to the “36-line”  $\mathcal{L}_{\overline{a}b} = \{L_{\overline{a}b} = 0\}$  passing through  $p_{\overline{a}b_1}$  and  $p_{\overline{a}b_2}$ . Furthermore, a dihedral  $\mathcal{D}_5$  stabilizes the “triangle”  $\{p_{\overline{a}b_1}, p_{\overline{a}b_2}, p_{\overline{a}b}\}$ :

$$\mathcal{D}_{\overline{a}b} = \text{Stab}\{p_{\overline{a}b}\} = \text{Stab}\{p_{\overline{a}b_1}, p_{\overline{a}b_2}\} = \langle P_{\overline{a}b}, Z_{\overline{a}cbd} \rangle \simeq \mathcal{D}_5.$$

An explanation of the indices attached to the element  $Z_{\overline{a}cbd}$  of order two occurs below.

As for other special orbits, each of the 45 involutions  $Z$  in  $\mathcal{V}$  is conjugate to

$$\begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix}.$$

In  $\mathbf{CP}^2$  each such  $Z$  fixes a point  $p_Z$  and point-wise fixes a line  $\mathcal{L}_Z$ . Furthermore,  $Z$  is the square of an element  $Q$  of order four. The three fixed points of  $Q$  consist of  $p_Z$  and two points  $(p_Q)_1$  and  $(p_Q)_2$  on  $\mathcal{L}_Z$ . Each of the latter points have a  $\mathbf{Z}/4$  stabilizer and so belong to a 90-point orbit  $\mathcal{O}_{90}$ . The points  $p_Z$ , having  $\mathcal{D}_4$  stabilizers, give an orbit  $\mathcal{O}_{45}$ . Since  $Z$  acts trivially on  $\mathcal{L}_Z$ , the  $\mathcal{D}_4$  action restricted to  $\mathcal{L}_Z$  reduces to that of a Klein-four group. Finally, the generic points on a “45-line” lie in four-point orbits and, overall, provide  $\mathcal{V}$ -orbits of size 180.

**The configuration of 45 lines and points.** The intersections of 45-lines yield special orbits of size less than 180. Furthermore, the number of lines meeting at such a site  $p$  corresponds to the number of involutions in the stabilizer of  $p$ . Being  $\mathcal{D}_5$ -stable, a 36-point lies on five of the

<sup>10</sup>In  $\mathcal{A}_6$ , the barred-unbarred splitting manifests itself in the two structurally distinct sets in order three:  $(abc)$  and  $(abc)(def)$ .

<sup>11</sup>The labeling of these lines in the sub-subscript is arbitrary and is done so as to agree with the natural cases in which a point does not reside on its associated line.

<sup>12</sup>See Figure 3.

<sup>13</sup>The practice of referring to special points and lines in terms of the size of their orbits will continue throughout.

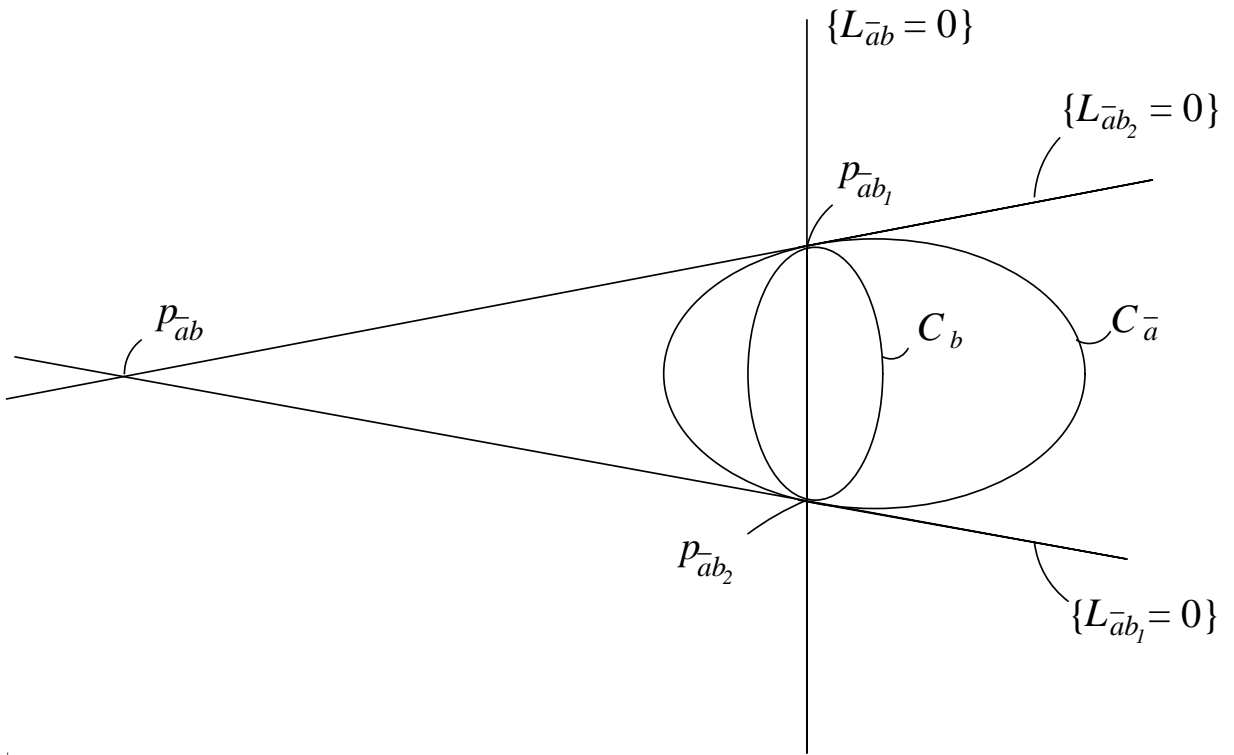


Figure 3: The triangle of one 36-point and two 72-points.

45-lines. Similarly, four of the 45-lines meet at a 45-point while three concur at each of the 60 and  $\overline{60}$ -points. This accounts for all intersections of the 45-lines:

$$36 \binom{5}{2} + 45 \binom{4}{2} + 60 \binom{3}{2} + 60 \binom{3}{2} = \binom{45}{2}.$$

On a 45-line there are four points of each type so that the corresponding clusters of lines complete the remaining set of 44 lines:

$$4 \left( \overbrace{4}^{36\text{-points}} + \underbrace{3}_{45\text{-points}} + \overbrace{2}^{60\text{-points}} + \underbrace{2}_{\overline{60}\text{-points}} \right) = 44.$$

Furthermore, only one involution fixes a 90-point. Hence, such points lie on just one of the 45-lines. Since a 72-point has a  $\mathbf{Z}/5$  stabilizer,  $\mathcal{O}_{72}$  acquires its exceptional status as the only special orbit that is not a subset of the 45-lines.

**Valentiner-speak.** Since the combinatorial relationships among special objects derive from those at the level of conics, the special orbits should admit description in a conic-based terminology. Already evident is the natural designation of 36, 72-points<sup>14</sup> in terms of one barred and one unbarred conic:

$$\{p_{\overline{ab}_1}, p_{\overline{ab}_2}\} = \mathcal{C}_{\overline{a}} \cap \mathcal{C}_b$$

and  $p_{\overline{ab}}$  is the pole of  $\{p_{\overline{ab}_1}, p_{\overline{ab}_2}\}$  with respect to  $\mathcal{C}_{\overline{a}}$  and  $\mathcal{C}_b$ . Turning to the 45-points, a given involution  $Z$  belongs to two tetrahedral subgroups of the barred and unbarred icosahedral systems alike. If  $Z$  preserves conics  $\mathcal{C}_{\overline{a}}, \mathcal{C}_{\overline{b}}, \mathcal{C}_c$ , and  $\mathcal{C}_d$ , then it does so uniquely.<sup>15</sup> Accordingly,  $Z_{\overline{abcd}}, p_{\overline{abcd}}$ , and  $\mathcal{L}_{\overline{abcd}}$  denote the transformation, corresponding point and line respectively. Alternatively, the 45-line  $\mathcal{L}_{\overline{abcd}}$  contains the 36-points  $p_{\overline{ac}}, p_{\overline{ad}}, p_{\overline{bc}}, p_{\overline{bd}}$  or by duality, the 45-point  $p_{\overline{abcd}}$  lies at the intersection of the 36-lines  $\mathcal{L}_{\overline{ac}}, \mathcal{L}_{\overline{ad}}, \mathcal{L}_{\overline{bc}}, \mathcal{L}_{\overline{bd}}$ . Furthermore,  $\mathcal{L}_{\overline{abcd}}$  is an icosahedral axis of order two for  $\mathcal{C}_{\overline{a}}, \mathcal{C}_{\overline{b}}, \mathcal{C}_c$ , and  $\mathcal{C}_d$  while

$$\mathcal{L}_{\overline{abcd}} \cap \mathcal{C}_{\overline{a}}, \mathcal{C}_{\overline{b}}, \mathcal{C}_c, \mathcal{C}_d$$

give the antipodal pairs of edge midpoints. Of course, the labels for the pair of 90-points—each of which  $Z_{\overline{abcd}}$  fixes—should be  $p_{\overline{abcd}_1}$  and  $p_{\overline{abcd}_2}$ .

Now, given a 45-line  $\mathcal{L}_{\overline{abcd}}$ , which four of the 45-points belong to it? Since there are  $15 = \binom{6}{2}$  choices each for the prefix  $\overline{ab}$  and suffix  $cd$ , a  $15 \times 15$  array with 45 distinguished entries depicts the configuration of 45-points and lines.<sup>16</sup> Being  $\mathcal{V}$ -equivalent, the rows and columns each contain

<sup>14</sup>By the Valentiner “duality” between special points and lines, whatever holds for such points also holds for the associated lines. Thus, I will usually suppress reference to one or the other.

<sup>15</sup>See below.

<sup>16</sup>See Figure 4.

three marked entries. Associated with each of the 15 barred rows  $\overline{ab}$  and its triple of “45-things” indexed by  $\{\overline{abcd}, \overline{abef}, \overline{abgh}\}$  where  $\{c, d, e, f, g, h\} = \{1, \dots, 6\}$  is a tetrahedral group

$$\mathcal{T}_{\overline{ab}} = \mathcal{I}_{\overline{a}} \cap \mathcal{I}_{\overline{b}}$$

whose three involutions are  $Z_{\overline{abcd}}, Z_{\overline{abef}}, Z_{\overline{abgh}}$ . The analogous state of affairs obtains for the unbarred columns where the tetrahedral group

$$\mathcal{T}_{cd} = \mathcal{I}_c \cap \mathcal{I}_d$$

contains involutions  $Z_{\overline{abcd}}, Z_{\overline{ijcd}}, Z_{\overline{kℓcd}}$ . Each of these 15 tetrahedral groups extends to an octahedral subgroup of  $\mathcal{V}$

$$\mathcal{O}_{\overline{ab}} = \text{Stab}\{p_{\overline{abcd}}, p_{\overline{abef}}, p_{\overline{abgh}}\} \quad \mathcal{O}_{cd} = \text{Stab}\{p_{\overline{abcd}}, p_{\overline{ijcd}}, p_{\overline{kℓcd}}\}.$$

Hence, the stabilizer of  $p_{\overline{abcd}}$  is the intersection of octahedral groups

$$\mathcal{O}_{\overline{ab}} \cap \mathcal{O}_{cd} = \text{Stab}\{p_{\overline{abcd}}\} = \text{Stab}\{\mathcal{L}_{\overline{abcd}}\} \simeq \mathcal{D}_4.$$

Furthermore, the involution  $Z_{\overline{abcd}}$  associates canonically with the pair of barred and unbarred tetrahedral groups

$$\mathcal{T}_{\overline{ab}} \cap \mathcal{T}_{cd} = \langle Z_{\overline{abcd}} \rangle \simeq \mathbf{Z}/2.$$

This “45-array” encodes a wealth of combinatorial geometry including an answer to the query of the preceding paragraph.<sup>17</sup> At a 45-point  $p_{\overline{abcd}}$  there are four concurrent 45-lines whose references have the form  $\mathcal{L}_{\overline{abef}}, \mathcal{L}_{\overline{abgh}}, \mathcal{L}_{\overline{mncd}},$  and  $\mathcal{L}_{\overline{rscd}}$ . To find these lines read along the  $\overline{ab}$  row and the  $cd$  column. By way of example,

$$p_{\overline{1234}} \in \mathcal{L}_{\overline{1212}} \cap \mathcal{L}_{\overline{1256}} \cap \mathcal{L}_{\overline{3634}} \cap \mathcal{L}_{\overline{4534}}.$$

Duality gives

$$\{p_{\overline{1212}}, p_{\overline{1256}}, p_{\overline{3634}}, p_{\overline{4534}}\} \subset \mathcal{L}_{\overline{1234}}.$$

The  $\mathcal{D}_5$  stabilizer of a 36-point  $p_{\overline{ab}}$  contains five involutions whose indices have a prefix  $\overline{a}$  and a suffix  $b$ . For  $p_{\overline{35}}$  this gives  $\overline{1335}, \overline{2315}, \overline{3445}, \overline{3556}, \overline{3625}$ . Hence,

$$p_{\overline{35}} \in \mathcal{L}_{\overline{1335}} \cap \mathcal{L}_{\overline{2315}} \cap \mathcal{L}_{\overline{3445}} \cap \mathcal{L}_{\overline{3556}} \cap \mathcal{L}_{\overline{3625}}$$

and

$$\{p_{\overline{1335}}, p_{\overline{2315}}, p_{\overline{3445}}, p_{\overline{3556}}, p_{\overline{3625}}\} \subset \mathcal{L}_{\overline{35}}.$$

The array also supplies a connection to  $\mathcal{A}_6$ . For instance, the involution  $Z_{\overline{1234}}$  fixes  $\mathcal{L}_{\overline{1234}}$  pointwise while preserving  $\mathcal{L}_{\overline{1212}}, \mathcal{L}_{\overline{1256}}, \mathcal{L}_{\overline{3634}},$  and  $\mathcal{L}_{\overline{4534}}$  as sets. Accordingly, it permutes the conics by:

$$\mathcal{C}_{\overline{3}} \leftrightarrow \mathcal{C}_{\overline{6}} \quad \mathcal{C}_{\overline{4}} \leftrightarrow \mathcal{C}_{\overline{5}}$$

$$\mathcal{C}_{\overline{1}} \leftrightarrow \mathcal{C}_{\overline{2}} \quad \mathcal{C}_{\overline{5}} \leftrightarrow \mathcal{C}_{\overline{6}}.$$

---

<sup>17</sup>A graphical version of this array appears in [Wiman 1895, p. 542].



Finally, which three involutions fix a  $\overline{60}$ , 60-point  $p$ ? (The stabilizer of  $p$  is a  $\mathcal{D}_3$ .) Unlike the other special orbits,  $\mathcal{O}_{\overline{60}}$  and  $\mathcal{O}_{60}$  have a bias toward one or the other system of conics. Recalling that  $p$  is an icosahedral face-center for two conics, say  $\mathcal{C}_{\overline{a}}$  and  $\mathcal{C}_{\overline{b}}$ , an involution  $Z$  that fixes  $p$  cannot preserve the two conics individually; such an action would have order three. Hence,  $Z$  exchanges the conics and lacks a prefix  $\overline{ab}$ . For  $\overline{ab} = \overline{25}$  the array indicates six such involutions:

$$\begin{array}{ll} Z_{\overline{1314}} : (\overline{25})(\overline{46}) & Z_{\overline{4614}} : (\overline{25})(\overline{13}) \\ Z_{\overline{3625}} : (\overline{25})(\overline{14}) & Z_{\overline{1425}} : (\overline{25})(\overline{36}) \\ Z_{\overline{3436}} : (\overline{25})(\overline{16}) & Z_{\overline{1636}} : (\overline{25})(\overline{34}). \end{array}$$

The six associated lines pass through the four points in  $\mathcal{C}_{\overline{2}} \cap \mathcal{C}_{\overline{5}}$  as edges of the tetrahedron whose stabilizer is  $\mathcal{T}_{\overline{25}}$ . They naturally fall into four sets of three lines; among these triples a given index in  $\{\overline{1}, \overline{3}, \overline{4}, \overline{6}\}$  appears twice. The intersection of the three lines occurs at the three-fold  $\overline{60}$ -points<sup>18</sup> thereby suggesting appropriate names:

$$\begin{array}{llll} \mathcal{L}_{\overline{1314}} \cap \mathcal{L}_{\overline{1425}} \cap \mathcal{L}_{\overline{1636}} & = & \{p_{\overline{1*346}}\} \\ \mathcal{L}_{\overline{1314}} \cap \mathcal{L}_{\overline{3436}} \cap \mathcal{L}_{\overline{3625}} & = & \{p_{\overline{3*146}}\} \\ \mathcal{L}_{\overline{1425}} \cap \mathcal{L}_{\overline{3436}} \cap \mathcal{L}_{\overline{4614}} & = & \{p_{\overline{4*136}}\} \\ \mathcal{L}_{\overline{1636}} \cap \mathcal{L}_{\overline{3625}} \cap \mathcal{L}_{\overline{4615}} & = & \{p_{\overline{6*134}}\}. \end{array}$$

Similarly, I will call the 60-points  $p_{a*bcd}$ .

To finish off the description of the special Valentiner points on a 45-line: Which four of the  $\overline{60}$ , 60-points lie on  $\mathcal{L}_{\overline{abcd}}$ ? Since  $\mathcal{L}_{\overline{abcd}}$  contains 60-points whose indices satisfy  $c*dst$  and  $d*cxxy$ , the matter comes down to finding values of  $s, t, x, y$  that are “Valentiner consistent.” This means that they fill out the scheme

$$\begin{array}{llll} p_{c*dst} & \in & \mathcal{L}_{\overline{cd}} \cap \mathcal{L}_{\overline{cs}} \cap \mathcal{L}_{\overline{ct}} \\ p_{c*duv} & \in & \mathcal{L}_{\overline{cd}} \cap \mathcal{L}_{\overline{cu}} \cap \mathcal{L}_{\overline{cv}} \\ p_{d*cxxy} & \in & \mathcal{L}_{\overline{cd}} \cap \mathcal{L}_{\overline{dx}} \cap \mathcal{L}_{\overline{dy}} \\ p_{d*czw} & \in & \mathcal{L}_{\overline{cd}} \cap \mathcal{L}_{\overline{dz}} \cap \mathcal{L}_{\overline{dw}} \end{array}$$

where each triple of prefixes exhausts  $\{\overline{1}, \dots, \overline{6}\}$  and

$$\{s, t, u, v\} = \{x, y, z, w\} = \{1, \dots, 6\} - \{c, d\}.$$

For  $\mathcal{L}_{\overline{1234}}$ ,

$$\begin{array}{llll} p_{3*456} & \in & \mathcal{L}_{\overline{1234}} \cap \mathcal{L}_{\overline{5635}} \cap \mathcal{L}_{\overline{3436}} \\ p_{3*124} & \in & \mathcal{L}_{\overline{1234}} \cap \mathcal{L}_{\overline{3513}} \cap \mathcal{L}_{\overline{4623}} \\ p_{4*356} & \in & \mathcal{L}_{\overline{1234}} \cap \mathcal{L}_{\overline{3445}} \cap \mathcal{L}_{\overline{5646}} \\ p_{4*123} & \in & \mathcal{L}_{\overline{1234}} \cap \mathcal{L}_{\overline{3524}} \cap \mathcal{L}_{\overline{4614}} \end{array}$$

---

<sup>18</sup>See Figure 5.

and

$$\begin{aligned}
p_{\overline{1*236}} &\in \mathcal{L}_{\overline{1234}} \cap \mathcal{L}_{\overline{1326}} \cap \mathcal{L}_{\overline{1615}} \\
p_{\overline{1*245}} &\in \mathcal{L}_{\overline{1234}} \cap \mathcal{L}_{\overline{1425}} \cap \mathcal{L}_{\overline{1516}} \\
p_{\overline{2*136}} &\in \mathcal{L}_{\overline{1234}} \cap \mathcal{L}_{\overline{2315}} \cap \mathcal{L}_{\overline{2626}} \\
p_{\overline{2*145}} &\in \mathcal{L}_{\overline{1234}} \cap \mathcal{L}_{\overline{2416}} \cap \mathcal{L}_{\overline{2525}}.
\end{aligned}$$

Collected below are the various data for  $\mathcal{L}_{\overline{1234}}$ .

Orbit	Special points on $\mathcal{L}_{\overline{1234}}$	Multiplicity of 45-lines
$\mathcal{O}_{36}$	$p_{\overline{13}} \ p_{\overline{14}} \ p_{\overline{23}} \ p_{\overline{24}}$	$\begin{pmatrix} 5 \\ 2 \end{pmatrix} = 10$
$\mathcal{O}_{45}$	$p_{\overline{1212}} \ p_{\overline{1256}} \ p_{\overline{3634}} \ p_{\overline{4534}}$	$\begin{pmatrix} 4 \\ 2 \end{pmatrix} = 6$
$\mathcal{O}_{60}$	$p_{\overline{1*236}} \ p_{\overline{2*136}} \ p_{\overline{1*245}} \ p_{\overline{2*145}}$	$\begin{pmatrix} 3 \\ 2 \end{pmatrix} = 3$
$\mathcal{O}_{60}$	$p_{3*124} \ p_{4*123} \ p_{3*456} \ p_{4*356}$	$\begin{pmatrix} 3 \\ 2 \end{pmatrix} = 3$
$\mathcal{O}_{90}$	$p_{\overline{1234_1}} \ p_{\overline{1234_2}}$	0

For subsequent reference, the following table summarizes the geometric terminology.

point/line	orbit	stabilizer
$p_{\overline{ab}}/\mathcal{L}_{\overline{ab}}$	$\mathcal{O}_{36}$	$\mathcal{D}_5$
$\{p_{\overline{ab_1}}, p_{\overline{ab_2}}\}/\{\mathcal{L}_{\overline{ab_1}}, \mathcal{L}_{\overline{ab_2}}\}$	$\mathcal{O}_{72}$	$\mathbf{Z}/5$
$p_{\overline{abcd}}/\mathcal{L}_{\overline{abcd}}$	$\mathcal{O}_{45}$	$\mathcal{D}_4$
$\{p_{\overline{abcd_1}}, p_{\overline{abcd_2}}\}/\{\mathcal{L}_{\overline{abcd_1}}, \mathcal{L}_{\overline{abcd_2}}\}$	$\mathcal{O}_{90}$	$\mathbf{Z}/4$
$p_{\overline{a*bcd}}/\mathcal{L}_{\overline{a*bcd}}$	$\mathcal{O}_{60}$	$\mathcal{D}_3$
$p_{a*bcd}/\mathcal{L}_{a*bcd}$	$\mathcal{O}_{60}$	$\mathcal{D}_3$

### 2.3.4 Computing a diagonal bub involution

One way to approach the matter of an antiholomorphic symmetry that exchanges systems of conics is to look for points that such a symmetry should fix. Given three such points  $a, b, c$  in coordinates  $y$  where

$$a = [1, 0, 0] \quad b = [0, 1, 0] \quad c = [0, 0, 1],$$

the associated bub map would have the diagonal form

$$\text{bub}[y_1, y_2, y_3] = [\alpha \overline{y_1}, \beta \overline{y_2}, \overline{y_3}].$$

Fixing a fourth point determines appropriate values for the inhomogeneous parameters  $\alpha, \beta$ .

But, which points should such a map fix? Moreover, how many such anti-involutions should there be?

	12	13	14	15	16	23	24	25	26	34	35	36	45	46	56
$\overline{12}$	•									•					•
$\overline{13}$			•						•		•				
$\overline{14}$		•						•						•	
$\overline{15}$					•	•							•		
$\overline{16}$				•			•					•			
$\overline{23}$				•		•								•	
$\overline{24}$					•		•				•				
$\overline{25}$			•					•				•			
$\overline{26}$		•							•				•		
$\overline{34}$	•											•	•		
$\overline{35}$		•					•								•
$\overline{36}$					•			•		•					
$\overline{45}$				•					•	•					
$\overline{46}$			•			•									•
$\overline{56}$	•										•			•	

Figure 4: A combinatorial scheme for the Valentiner group

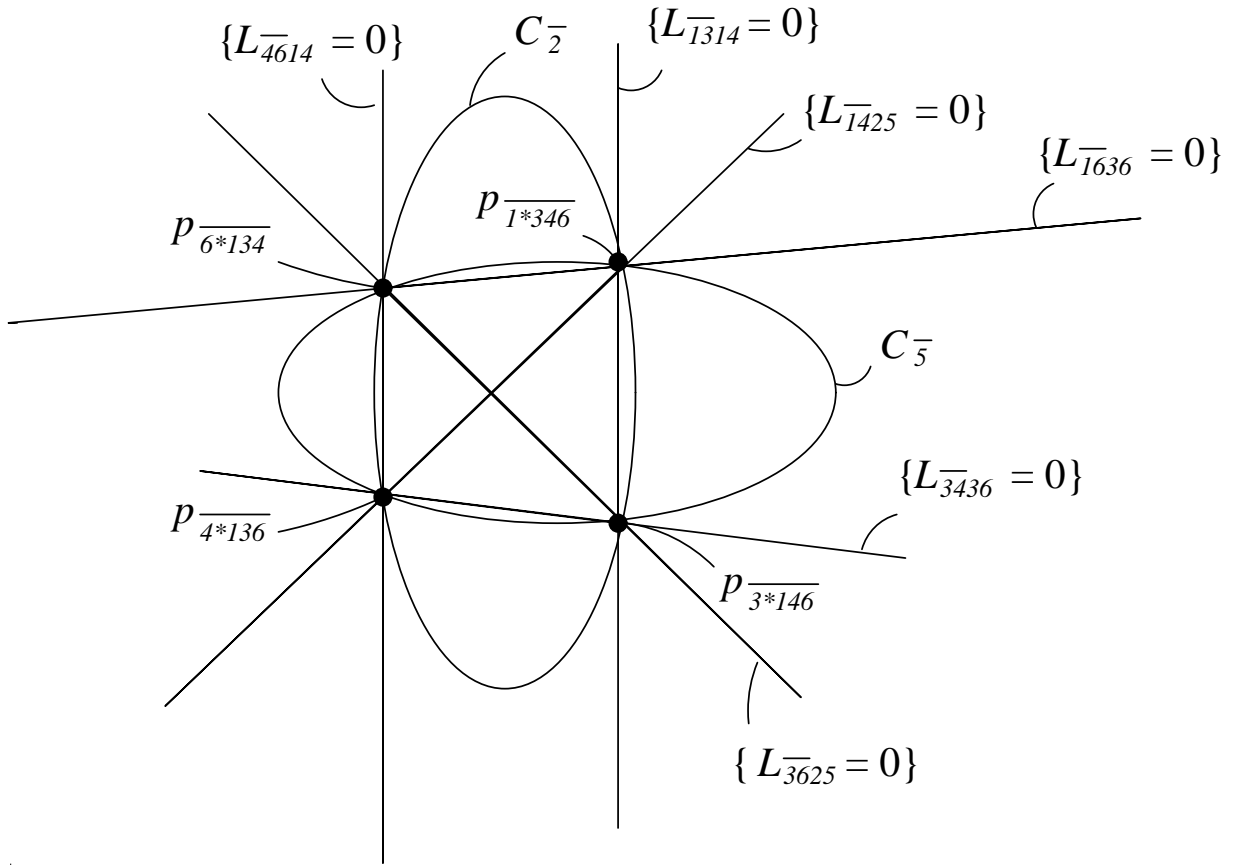


Figure 5: Intersecting conics within a system—a tetrahedral configuration

**A heuristic for bub-symmetry.** The basic Valentiner item that involves a mixture of barred and unbarred conics is the  $\mathcal{D}_5$  structure consisting of a pair of conics  $\{\mathcal{C}_{\bar{a}}, \mathcal{C}_b\}$  that intersect in the pair of 72-points  $\{p_{\bar{a}b_1}, p_{\bar{a}b_2}\}$ . A bub map that exchanges these two conics, must preserve the set  $\{p_{\bar{a}b_1}, p_{\bar{a}b_2}\}$  as well as the associated 36-point  $p_{\bar{a}b}$ . To put some flesh on the skeletal configuration of Figure 3, consider two icosahedra that 1) share a five-fold axis and 2) about this axis, are one-tenth of a revolution away from each other.<sup>19</sup> The poles where the axis passes through the icosahedra correspond to the pair of five-fold points  $p_{\bar{a}b_{1,2}}$ . A reflection through the equatorial plane preserves this arrangement while exchanging the icosahedra and the poles. The icosahedra also transpose under reflection through five planes that include the polar axis. In these cases, the two poles are fixed. This model hints that for each pair  $\mathcal{C}_{\bar{a}}$  and  $\mathcal{C}_b$ , there is a distinguished bub-involution and five of a secondary nature. This makes for a total of 36 maps  $\text{bub}_{\bar{a}b}$ .

For the primary reflection relative to the pair  $\{\mathcal{C}_{\bar{2}}, \mathcal{C}_2\}$ , this heuristic demands that  $p_{\bar{2}2_1}$  and  $p_{\bar{2}2_2}$  exchange while  $p_{\bar{2}2}$  remains fixed. What other points should “ $\text{bub}_{\bar{2}2}$ ” fix? Since five other bub maps switch  $\mathcal{C}_{\bar{2}}$  and  $\mathcal{C}_2$ , symmetry requires that  $\text{bub}_{\bar{2}2}$  provide a secondary reflection for each barred-unbarred pair of icosahedra associated with the  $\bar{2}2$  configuration. As such,  $\text{bub}_{\bar{2}2}$  fixes the corresponding poles of 72-points. Accordingly, the correspondence between the five remaining barred and unbarred conics determines these five pairs of points. Now, the five pairs of non-polar antipodal vertices on the  $\mathcal{C}_{\bar{2}}$  icosahedron correspond to the points of intersection of  $\mathcal{C}_{\bar{2}}$  with the five conics  $\mathcal{C}_1, \mathcal{C}_3, \mathcal{C}_4, \mathcal{C}_5$ , and  $\mathcal{C}_6$  while the non-polar vertices on the  $\mathcal{C}_2$  icosahedron correspond to the intersections of  $\mathcal{C}_2$  with the five conics  $\mathcal{C}_{\bar{1}}, \mathcal{C}_{\bar{3}}, \mathcal{C}_{\bar{4}}, \mathcal{C}_{\bar{5}}$ , and  $\mathcal{C}_{\bar{6}}$ . Let these sets of five pairs correspond to vertices of two pentagons that are one-tenth of a revolution away from each other<sup>20</sup> with antipodal pairs of points being  $\text{bub}_{\bar{2}2}$  symmetric. The  $\mathcal{D}_5$  action  $\mathcal{D}_{\bar{2}2} = \text{Stab}\{\mathcal{C}_{\bar{2}}, \mathcal{C}_2\}$  determines the specific arrangement.

One of the elements of order five that belongs to  $\mathcal{D}_{\bar{2}2}$  is

$$P_{\bar{2}2} = QPQ^{-1}.$$

The associated five-cyclings of conics are  $(\overline{15436})$  and  $(15436)$ . The matching of the  $\bar{2}k$  and  $\overline{m}2$  vertices depends upon the five involutions that stabilize  $\mathcal{C}_{\bar{2}}$  and  $\mathcal{C}_2$ . For example, the generator  $Z_{\bar{1}2_{12}}$  associates  $\bar{1}$  with 1 while the 5-cycles above determine the remaining matches of  $\bar{a}$  with  $a$ . This information is also readily available in the 45-array. The five entries that involve both  $\bar{2}$  and 2 are  $\bar{2}a2a$ . Indeed, the array’s symmetry about the diagonal  $\overline{a}bab$  is a combinatorial manifestation of  $\text{bub}_{\bar{2}2}$ .

<sup>19</sup>See Figure 6.

<sup>20</sup>See Figure 7 and imagine looking down, from above a pole, on the intersecting icosahedra of Figure 6.

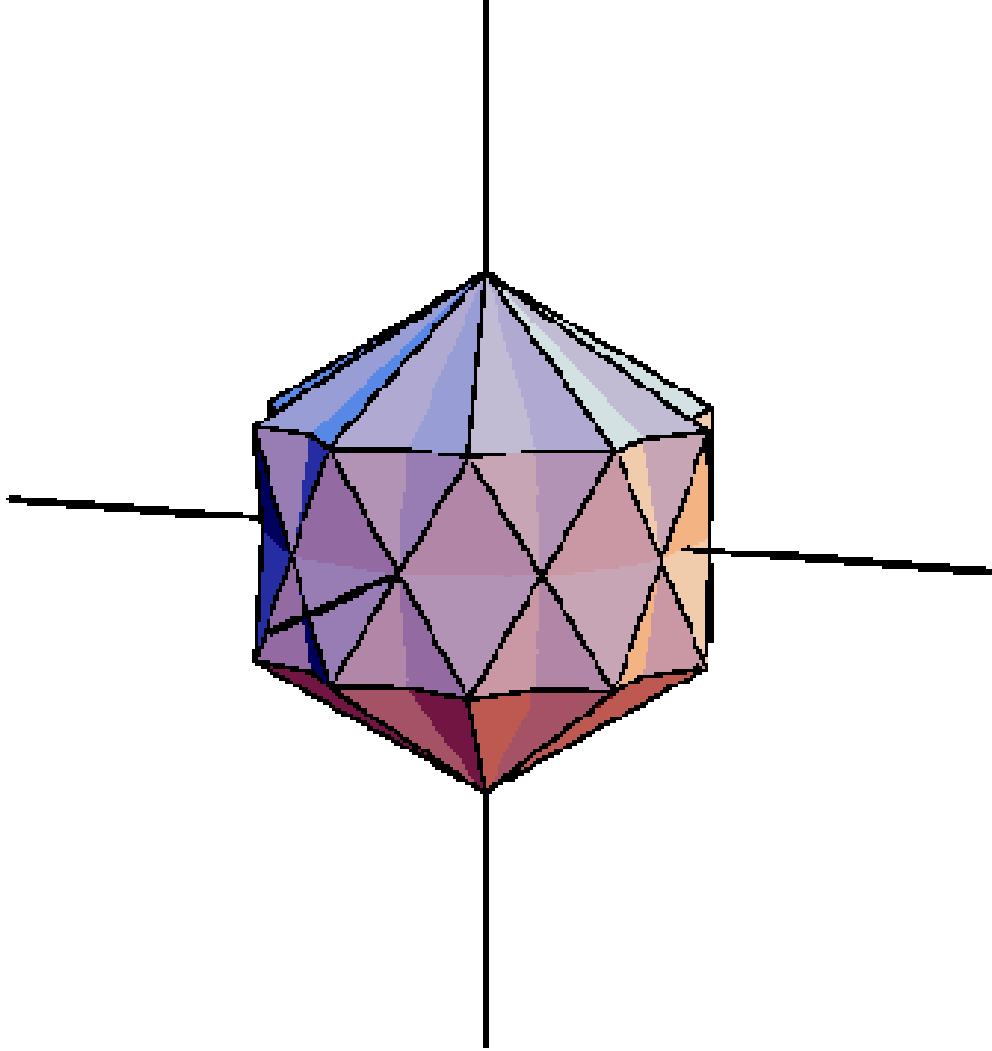


Figure 6: Intersecting icosahedra: A heuristic for bub-symmetry

The union of a barred and unbarred conic has a  $\mathcal{D}_5$  structure represented by two icosahedra that “meet” at a pair of antipodal vertices and are turned away from one another by an angle of  $\pi/5$ . The reflection through the equatorial plane exchanges the icosahedra and so suggests that for each pair  $\{\mathcal{C}_{\bar{a}}, \mathcal{C}_b\}$  there is a “primary” bub involution. Also transposing the icosahedra are “secondary” reflections through five vertical planes. These correspond to primary reflections for five other pairs of conics.

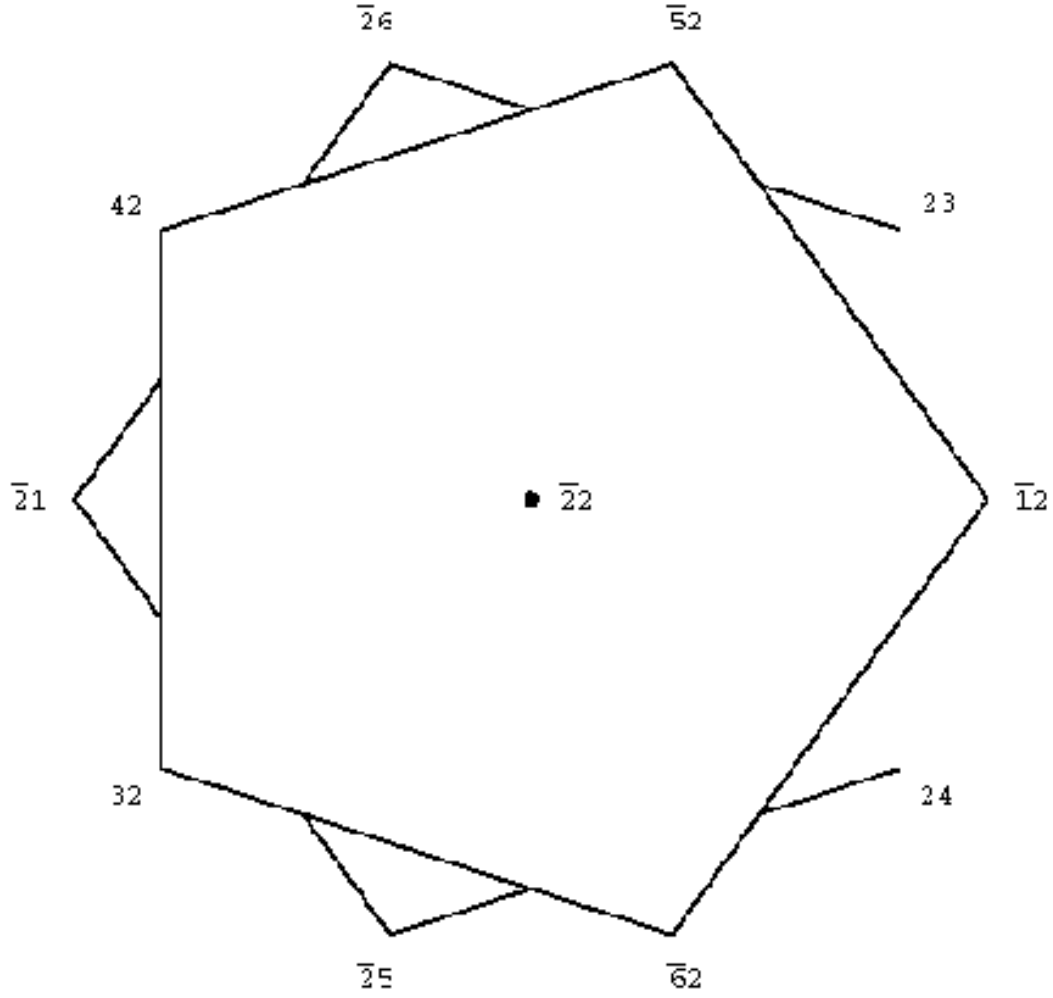


Figure 7: A heuristic for bub-symmetry

Regard each of the five pairs of antipodal vertices on the  $\mathcal{D}_5$  union of conics  $\mathcal{C}_2$  and  $\mathcal{C}_2$  as a vertex of one of two pentagons whose arrangement corresponds to that of the remaining icosahedra. The primary  $\text{bub}_{\overline{2}2}$  reflection interchanges the pentagons as well as antipodal vertices. The secondary reflections are  $\{\text{bub}_{\overline{a}a} | a \neq 2\}$  which transpose the vertices  $\overline{2}a$  and  $\overline{a}2$  respectively.

**Special coordinates and the bub- $\mathbf{RP}^2$ .** Following the clue provided by the above heuristic, make the parametrized change of icosahedral to octahedral coordinates

$$\begin{aligned} A &= \left( a p_{\bar{1}1_1}^t \mid b p_{\bar{1}1_2}^t \mid p_{\bar{1}1}^t \right) \\ &= \left( a \begin{pmatrix} \frac{1-\sqrt{5}}{2} \\ -\sqrt{\frac{5-\sqrt{5}}{2}} i \\ 1 \end{pmatrix} \mid b \begin{pmatrix} \frac{1-\sqrt{5}}{2} \\ \sqrt{\frac{5-\sqrt{5}}{2}} i \\ 1 \end{pmatrix} \mid \begin{pmatrix} \frac{1+\sqrt{5}}{2} \\ 0 \\ 1 \end{pmatrix} \right). \end{aligned}$$

In these icosahedral coordinates  $y$ , the  $\bar{1}1$  triangle is

$$p_{\bar{1}1_1} = [1, 0, 0] \quad p_{\bar{1}1_2} = [0, 1, 0] \quad p_{\bar{1}1} = [0, 0, 1].$$

The candidate bub $_{\bar{2}2}$  map  $K(y) = \bar{y}$  fixes each of these. In octahedral coordinates the  $\bar{2}2$  triangle consists of

$$\begin{aligned} p_{\bar{2}2_1} &= \left[ \frac{1+\sqrt{5}}{2} \rho^2, \sqrt{\frac{5+\sqrt{5}}{2}} i \rho, 1 \right] \\ p_{\bar{2}2_2} &= \left[ \frac{1+\sqrt{5}}{2} \rho^2, -\sqrt{\frac{5+\sqrt{5}}{2}} i \rho, 1 \right] \\ p_{\bar{2}2} &= \left[ \frac{1-\sqrt{5}}{2} \rho^2, 0, 1 \right]. \end{aligned}$$

The hope here is that, when transformed to  $y$  coordinates, some choice of  $a, b$  results in

$$K(A^{-1} p_{\bar{2}2_1}) = p_{\bar{2}2_2} \quad K(A^{-1} p_{\bar{2}2_2}) = p_{\bar{2}2_1} \quad K(A^{-1} p_{\bar{2}2}) = p_{\bar{2}2}.$$

Satisfying these conditions are the values

$$a = b = \frac{\sqrt{3} - \sqrt{5} i}{8}.$$

The change of coordinates becomes

$$A = \begin{pmatrix} \frac{(1-\sqrt{5})(3-\sqrt{5}i)}{4\sqrt{2}} & \frac{(1-\sqrt{5})(3-\sqrt{5}i)}{4\sqrt{2}} & \frac{1+\sqrt{5}}{2} \\ \frac{\sqrt{5-\sqrt{5}}(3-\sqrt{5}i)}{4} & \frac{\sqrt{5-\sqrt{5}}(3-\sqrt{5}i)}{4} & 0 \\ \frac{3-\sqrt{5}i}{2\sqrt{2}} & \frac{3-\sqrt{5}i}{2\sqrt{2}} & 1 \end{pmatrix}.$$

As for the conic forms, they satisfy the desired condition:

$$C_{\bar{k}}(y) = \overline{C_k(\bar{y})}.$$

A further change of coordinates given by a real diagonal matrix leaves bub $_{\bar{2}2}(y) = \bar{y}$  undisturbed. In  $y$  coordinates the one-point orbit  $p_{\bar{2}2}$  under the  $\mathcal{D}_5$  for  $\bar{2}2$  is<sup>21</sup>

$$\left[ 1, 1, \frac{\sqrt{6}}{\tau} \right].$$

---

<sup>21</sup>Recall that  $\tau = (1 + \sqrt{5})/2$ .



For a final simplification, the additional coordinate change

$$B = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & \frac{\sqrt{6}}{\tau} \end{pmatrix}$$

arranges for this point to be  $[1, 1, 1]$ . In these adjusted  $y$  coordinates, the  $\bar{1}1$  and  $\bar{2}2$  triangles are<sup>22</sup>

$$\begin{aligned} p_{\bar{1}1_1} &= [1, 0, 0] & p_{\bar{2}2_1} &= [3, 2\eta^2, -\eta] \\ p_{\bar{1}1_2} &= [0, 1, 0] & p_{\bar{2}2_2} &= [3, 2\bar{\eta}^2, -\bar{\eta}] \\ p_{\bar{1}1} &= [0, 0, 1] & p_{\bar{2}2} &= [1, 1, 1] \end{aligned}$$

while, in these “bub $_{\bar{2}2}$  coordinates” the normalized conic forms<sup>23</sup> for  $\bar{1}$  and  $1$  are

$$\begin{aligned} C_{\bar{1}}(y) &= \left(\frac{2}{3}\bar{\eta}\right)^2 y_1 y_2 + y_3^2 \\ C_1(y) &= \left(\frac{2}{3}\eta\right)^2 y_1 y_2 + y_3^2. \end{aligned}$$

Since bub $_{\bar{2}2}$  restricts to the identity on the  $\mathbf{RP}^2$  given by

$$\mathcal{R}_{\bar{2}2} = \{[t_1, t_2, t_3] \mid t_k \in \mathbf{R}\},$$

symmetry provides for such a fixed set  $\mathcal{R}_{\bar{a}b}$  for each of the 36 maps bub $_{\bar{a}b}$ . Figure 8 provides a geometric interpretation of these  $\mathbf{RP}^2$ s. One consequence of the extra symmetry is that  $\bar{\mathcal{V}}_{2.360}$ -invariant forms and  $\bar{\mathcal{V}}_{2.360}$ -equivariant maps are, when expressed in special bub-coordinates, given by polynomials with real and even, in special cases, rational coefficients.<sup>24</sup>

## 2.4 Invariant structure

For a group action on a vector space the Molien series provides one of the basic tools of classical invariant theory. Given a finite group  $\mathcal{G}$  acting faithfully on  $\mathbf{C}^n$ , the dimension of the space  $\mathbf{C}[x]_m^{\mathcal{G}}$  of invariant homogeneous polynomials of degree  $m$  appears as the coefficient of the  $m$ th degree term in the *Molien series* for  $\mathcal{G}$ :

$$M(\mathbf{C}[x]^{\mathcal{G}}) = \sum_{m=0}^{\infty} \left( \dim \mathbf{C}[x]_m^{\mathcal{G}} \right) t^m.$$

In the Valentiner case the space is  $\mathbf{C}^3$  while the group is a 1-to-3 lift of  $\mathcal{V}$  to a subgroup  $\mathcal{V}_{3.360}$  of  $\mathrm{SU}_3$ . As a result of the character  $\langle \rho \rangle$  that appears under  $\mathcal{V}$ ’s action on the icosahedral conic

<sup>22</sup>Recall that  $\eta = (3 + \sqrt{15}i)/4$ .

<sup>23</sup>The unwieldy expressions for the remaining forms are not recorded here.

<sup>24</sup>See Section 2.4.2. Although [Wiman 1895, pp. 548-50] and [Fricke 1926, pp. 286-9] mention these coordinates, they seem not to have made much use of them.

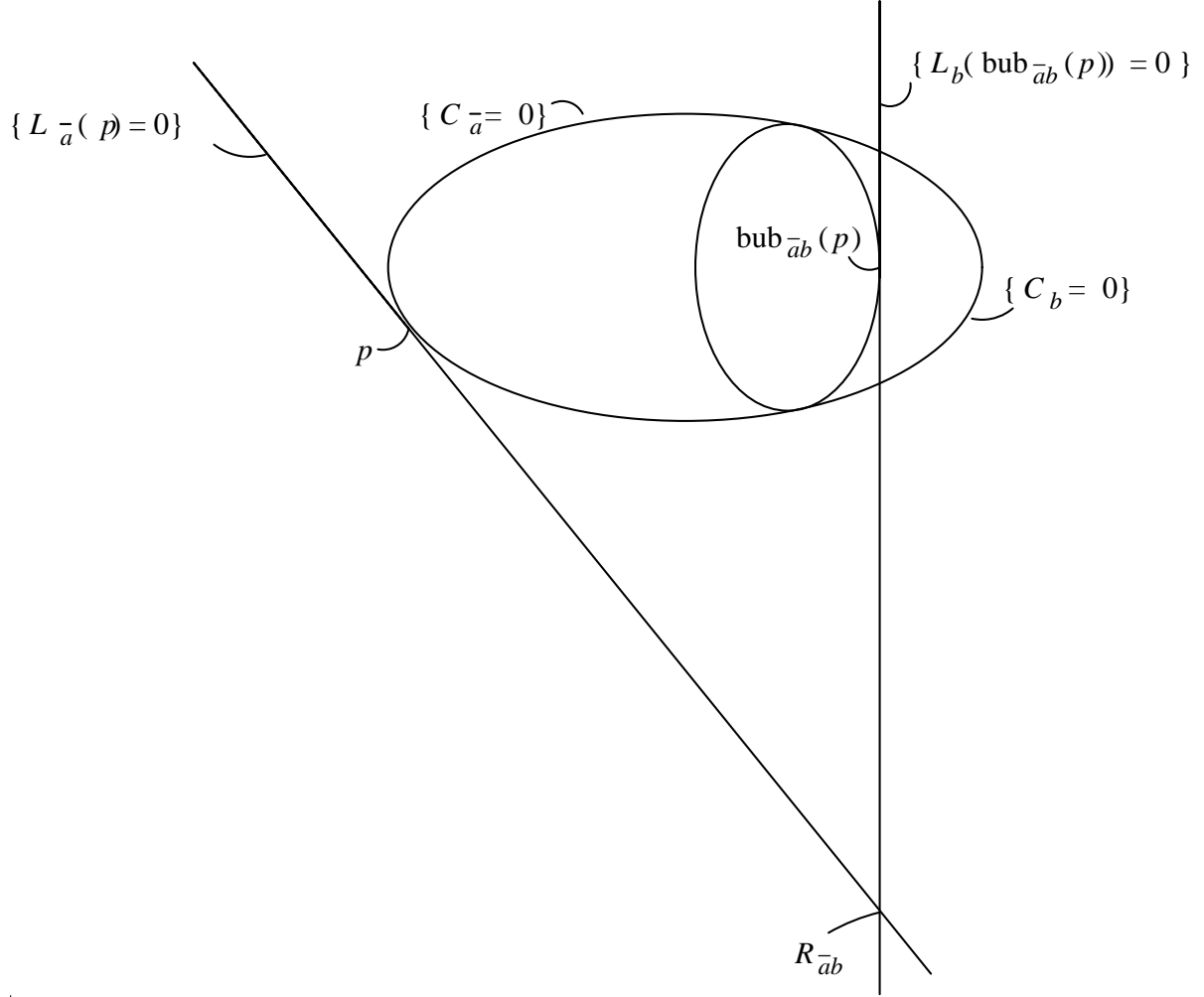


Figure 8: A geometric interpretation of bub-maps  
 Since  $\text{bub}_{\bar{a}b}$  interchanges the pair of conics  $\mathcal{C}_{\bar{a}}$  and  $\mathcal{C}_b$  as well as the lines

$$\mathcal{L}_{\bar{a}}(p) = \{L_{\bar{a}}(p) = 0\} \quad \mathcal{L}_b(\text{bub}_{\bar{a}b}(p)) = \{L_b(p) = 0\}$$

that are tangent to  $\mathcal{C}_{\bar{a}}$  at  $p$  and to  $\mathcal{C}_b$  at  $\text{bub}_{\bar{a}b}(p)$ , it also fixes the  $\mathbf{RP}^2$ 's worth of points

$$\mathcal{R}_{\bar{a}b} = \{\mathcal{L}_{\bar{a}}(p) \cap \mathcal{L}_b(\text{bub}_{\bar{a}b}(p)) | p \in \mathcal{C}_{\bar{a}}\}.$$

forms, this lift of  $\mathcal{V}$  to a linear group has minimal order.<sup>25</sup> A further consequence of minimality is that the Molien series for the  $\mathcal{V}_{3.360}$  gives complete information concerning the invariants of the projective group  $\mathcal{V}$ .

**Proposition 3** *The invariants of  $\mathcal{V}$  and  $\mathcal{V}_{3.360}$  are in one-to-one correspondence.*

*Proof.* Trivially, a  $\mathcal{V}_{3.360}$ -invariant gives a  $\mathcal{V}$ -invariant. Conversely, let  $F(x)$  be a  $\mathcal{V}$ -invariant with

$$F(T^{-1}x) = \alpha(T) F(x)$$

for  $T \in \mathcal{V}_{3.360}$ . The kernel of the multiplicative character

$$\alpha : \mathcal{V}_{3.360} \longrightarrow \mathbf{C} - \{0\}$$

is the normal subgroup  $\text{Stab}(F) \subset \mathcal{V}_{3.360}$  that stabilizes  $F$ . Since the projective image  $[\text{Stab}(F)] \simeq \text{Stab}(F)/\langle \rho \rangle$  is normal in the simple group  $\mathcal{V} \simeq \mathcal{V}_{3.360}/\langle \rho \rangle$ ,  $[\text{Stab}(F)]$  is either trivial or  $\mathcal{V}$ . In the former case,  $\text{Stab}(F)$  would be either trivial or  $\langle \rho \rangle$  so that  $\mathcal{V}_{3.360}/\text{Stab}(F)$  would be non-abelian. Thus,  $[\text{Stab}(F)] = \mathcal{V}$ . Since  $\mathcal{V}_{3.360}$  includes no subgroup of order 360,  $\text{Stab}(F) = \mathcal{V}_{3.360}$  and  $F$  is  $\mathcal{V}_{3.360}$ -invariant.  $\triangle$

Also,  $\mathcal{V}$  lifts 1-to-6 to a so-called unitary reflection group  $\mathcal{V}_{6.360}$  generated by 45 involutions on  $\mathbf{C}^3$ . [Shephard and Todd 1954, pp. 278, 287] The elements of  $\mathcal{V}_{6.360}$  satisfy  $\det T = \pm 1$  while those of  $\mathcal{V}_{3.360}$  satisfy  $\det T = 1$ .

#### 2.4.1 Molien's theorem and its application to $\mathcal{V}$

By projecting  $\mathbf{C}[x]_m$  onto  $\mathbf{C}[x]_m^{\mathcal{G}}$ , one arrives at a generating function for the Molien series.

**Theorem 1** *For a finite group action  $\mathcal{G}$  on  $\mathbf{C}^n$ ,*

$$M(\mathcal{G}) = \frac{1}{|\mathcal{G}|} \sum_{C_T \subset \mathcal{G}} \frac{|C_T|}{\det(I - tT^{-1})}$$

where  $C_T$  are conjugacy classes.

*Proof.* See [Benson 1993, pp. 21-22].  $\triangle$

**Proposition 4** *For the Valentiner groups  $\mathcal{V}_{3.360}$  and  $\mathcal{V}_{6.360}$ , the Molien series are given by*

$$\begin{aligned} M(\mathcal{V}_{3.360}) &= \frac{1 + t^{45}}{(1 - t^6)(1 - t^{12})(1 - t^{30})} \\ &= 1 + t^6 + 2t^{12} + 2t^{18} + \dots + t^{45} + \dots \\ M(\mathcal{V}_{6.360}) &= \frac{1}{(1 - t^6)(1 - t^{12})(1 - t^{30})} \\ &= 1 + t^6 + 2t^{12} + 2t^{18} + \dots \end{aligned}$$

---

<sup>25</sup>A lift of a projective group  $\mathcal{G}$  to a linear group  $\mathcal{G}'$  has *minimal* order if for every lift  $\mathcal{H}$  of  $\mathcal{G}$   $|\mathcal{H}| \geq |\mathcal{G}'|$ . See [Fricke 1926, pp. 267-8] for details.

*Proof.* With  $k = 0, 1, 2$ , the matrices

$$\pm \rho^k I, \pm \rho^k P, \pm \rho^k Z, \pm \rho^k Q, \pm \rho^k PZ, \pm \rho^k T$$

represent distinct conjugacy classes in  $\mathcal{V}_{6.360}$ . For  $\mathcal{V}_{3.360}$  the three matrices of each type corresponding to the  $+\rho^k$  do the job. Substitution into the formula of Molien's theorem produces the indicated generating functions.  $\triangle$

#### 2.4.2 The basic invariants themselves

From the theory of complex reflection groups there are three algebraically independent “basic” forms that generate the ring of  $\mathcal{V}_{6.360}$ -invariants. The generating function for the Molien series indicates that these occur in degrees 6, 12, and 30. Techniques of classical invariant theory provide for the computation of the forms in degrees 12 and 30 from that of degree 6. But, how does the latter arise? Although  $\mathcal{V}_{6.360}$  permutes the *conics*, its action on the conic forms is not “simple”—a non-trivial character appears. However, the cubes of the forms do receive simple treatment by  $\mathcal{V}_{6.360}$ . Hence, summing the cubes of either system of conic forms and normalizing the coefficients yields a  $\mathcal{V}_{6.360}$ -invariant:

$$\begin{aligned} F(x) &= \alpha \sum_{m=1}^6 C_{\overline{m}}(x)^3 \\ &= \alpha \sum_{m=1}^6 C_m(x)^3 \\ &= x_1^6 + x_2^6 + x_3^6 + 3 \left( 5 - \sqrt{15} i \right) x_1^2 x_2^2 x_3^2 \\ &\quad + \frac{3}{4} \left( 2 \sqrt{5} - \left( 5 - \sqrt{5} \right) \rho \right) \left( x_1^4 x_2^2 + x_2^4 x_3^2 + x_1^2 x_3^4 \right) \\ &\quad - \frac{3}{4} \left( 2 \sqrt{5} + \left( 5 + \sqrt{5} \right) \rho^2 \right) \left( x_1^4 x_3^2 + x_1^2 x_2^4 + x_2^2 x_3^4 \right). \end{aligned}$$

By uniqueness,  $F$  is also  $\overline{\mathcal{V}}_{2.360}$ -invariant. Expressed and normalized in  $\text{bub}_{\overline{22}}$  coordinates,

$$F(y) = 10 y_1^3 y_2^3 + 9 y_1^5 y_3 + 9 y_2^3 y_3 - 45 y_1^2 y_2^2 y_3^2 - 135 y_1 y_2 y_3^4 + 27 y_3^6.$$

The forms  $\Phi$  and  $\Psi$  of degrees 12 and 30 arise respectively from the determinants of the

Hessian  $H_F$  of  $F$  and the “bordered Hessian”  $BH(F, \Phi)$  of  $F$  and  $\Phi$ :

$$\begin{aligned}
\Phi(y) &= \alpha_\Phi |H(F(y))| \\
&= 6 y_1^{11} y_2 - 38 y_1^6 y_2^6 + 6 y_1 y_2^{11} + 90 y_1^8 y_2^3 y_3 + 90 y_1^3 y_2^8 y_3 - 9 y_1^{10} y_3^2 - \\
&\quad 468 y_1^5 y_2^5 y_3^2 - 9 y_2^{10} y_3^2 + 1080 y_1^7 y_2^2 y_3^3 + 1080 y_1^2 y_2^7 y_3^3 + \\
&\quad 3375 y_1^4 y_2^4 y_3^4 - 324 y_1^6 y_2 y_3^5 - 324 y_1 y_2^6 y_3^5 - 1080 y_1^3 y_2^3 y_3^6 + 2916 y_1^5 y_3^7 + \\
&\quad 2916 y_2^5 y_3^7 + 1215 y_1^2 y_2^2 y_3^8 + 4374 y_1 y_2 y_3^{10} + 729 y_3^{12}
\end{aligned}$$

$$\begin{aligned}
\Psi(y) &= \alpha_\Phi |BH(F(y), \Phi(y))| \\
&= \alpha_\Phi \left| \begin{array}{c|c} H(\Phi(y)) & \begin{matrix} F_{y_1} \\ F_{y_2} \\ F_{y_3} \end{matrix} \\ \hline \begin{matrix} F_{y_1} & F_{y_2} & F_{y_3} \end{matrix} & 0 \end{array} \right| \\
&= 3 y_1^{30} + \dots + 3 y_2^{30} + \dots + 57395628 y_3^{30}.
\end{aligned}$$

The constants  $\alpha_\Phi = -1/20250$  and  $\alpha_\Psi = 1/24300$  remove the highest common factor among the coefficients.

Finally, the product of the 45 linear forms that correspond to the generating involutions is a relative  $\mathcal{V}_{6,360}$ -invariant but an absolute  $\mathcal{V}_{3,360}$ -invariant and hence, a projective  $\mathcal{V}$ -invariant. As a specific instance of a general result [Shephard and Todd 1954, p. 283], this degree 45 form is given by the Jacobian determinant

$$\begin{aligned}
X(y) &= \alpha_X |J(F(y), \Phi(y), \Psi(y))| \\
&= \alpha_X \left| \begin{array}{ccc} F_{y_1} & F_{y_2} & F_{y_3} \\ \Phi_{y_1} & \Phi_{y_2} & \Phi_{y_3} \\ \Psi_{y_1} & \Psi_{y_2} & \Psi_{y_3} \end{array} \right| \\
&= \beta \prod L_{\overline{abcd}}(y) \\
&= y_1^{45} + \dots - y_2^{45} + \dots + 3570467226624 y_2^5 y_3^{40}
\end{aligned}$$

where  $\alpha_X = -1/4860$  and  $\beta$  is a constant. Being  $\mathcal{V}_{6,360}$ -invariant,  $X^2$  is a polynomial in  $F$ ,  $\Phi$ ,  $\Psi$ :

$$\begin{aligned}
3^9 \cdot X^2 &= 4 F^{13} \Phi + 80 F^{11} \Phi^2 + 816 F^9 \Phi^3 + 4376 F^7 \Phi^4 + 13084 F^5 \Phi^5 + \\
&\quad 12312 F^3 \Phi^6 + 5616 F \Phi^7 + 18 F^{10} \Psi + 198 F^8 \Phi \Psi + 954 F^6 \Phi^2 \Psi - \\
&\quad 198 F^4 \Phi^3 \Psi - 5508 F^2 \Phi^4 \Psi - 1944 \Phi^5 \Psi - 162 F^5 \Psi^2 - 1944 F^3 \Phi \Psi^2 - \\
&\quad 1458 F \Phi^2 \Psi^2 + 729 \Psi^3.
\end{aligned} \tag{1}$$

### 2.4.3 $\mathcal{V}$ -symmetric maps and the sextic

The system of invariants provides a foundation on which to construct mappings of  $\mathbf{C}^3$  or  $\mathbf{CP}^2$  that are symmetric or *equivariant* under the action of  $\mathcal{V}_{3,360}$  or  $\mathcal{V}$ . Algebraically, this means that the map commutes with the action. Given such a map that also possesses nice dynamical properties, the sixth-degree equation has an iterative solution.

## 3 Rational Maps with Valentiner Symmetry

An iterative solution to the sextic utilizes a parametrized family of dynamical systems having  $\mathcal{A}_6$  symmetry. In practice, a given sixth-degree polynomial with galois group  $\mathcal{A}_6$  specifies a projective transformation

$$S_p : \mathbf{CP}^2 \rightarrow \mathbf{CP}^2$$

and thereby “hooks-up” to a rational map

$$S_p^{-1} \circ f \circ S_p$$

that has  $\mathcal{A}_6$  symmetry. Accordingly, the fixed map  $f$  is the centerpiece of a sextic-solving algorithm.

### 3.1 Finding equivariant maps

A linear group  $\mathcal{G}$  acts on the exterior algebra  $\Lambda(\mathbf{C}^n)$  by

$$(T(\alpha))(x) = \alpha(T^{-1}x)$$

where  $T \in \mathcal{G}$ ,  $\alpha$  is a  $p$ -form, and  $x \in \mathbf{C}^n$ . As in the case of  $\mathcal{V}$ -invariant polynomials,  $\mathcal{V}$ -invariant  $p$ -forms associate one-to-one with  $\mathcal{V}_{3,360}$ -invariant  $p$ -forms. Hence, the search for symmetric maps can take place within the regime of the linear action.

When looking for equivariants, the  $\mathcal{V}_{3,360}$  action on  $\Lambda(\mathbf{C}^3)$  provides guidance. Such utility is due to a correspondence between  $\mathcal{V}_{3,360}$ -equivariants and  $\mathcal{V}_{3,360}$ -invariant 2-forms. Let  $dX_2 = (dx_2 \wedge dx_3, dx_3 \wedge dx_1, dx_1 \wedge dx_2)$  and ‘ $\cdot$ ’ signify a formal dot product.

**Proposition 5** *For a given finite action  $\mathcal{G} \subset \mathrm{U}_3$  and a  $\mathcal{G}$ -invariant 2-form*

$$\begin{aligned} \phi(x) &= f_1(x) dx_2 \wedge dx_3 + f_2(x) dx_3 \wedge dx_1 + f_3(x) dx_1 \wedge dx_2 \\ &= f(x) \cdot dX_2, \end{aligned}$$

*the map  $f = (f_1, f_2, f_3)$  is relatively  $\mathcal{G}$ -equivariant. If  $\mathcal{G} \subset \mathrm{SU}_3$ , then the equivariance of  $f$  is absolute.*

*Proof.* Given  $T \in \mathcal{G}$ , let

$$T = \begin{pmatrix} t_{11} & t_{12} & t_{13} \\ t_{21} & t_{22} & t_{23} \\ t_{31} & t_{32} & t_{33} \end{pmatrix} \quad T^{-1} = |T|^{-1} \begin{pmatrix} C_{11} & C_{12} & C_{13} \\ C_{21} & C_{22} & C_{23} \\ C_{31} & C_{32} & C_{33} \end{pmatrix}.$$

For an invariant 2-form  $\phi$ ,

$$\begin{aligned}
f(x) \cdot dX_2 &= \phi(x) \\
&= T(\phi(x)) \\
&= f(T^{-1}x) \cdot |T|^{-2} ( \\
&\quad (C_{22}C_{33} - C_{23}C_{32}) dx_2 \wedge dx_3 + (C_{23}C_{31} - C_{21}C_{33}) dx_3 \wedge dx_1 + \\
&\quad (C_{21}C_{32} - C_{22}C_{31}) dx_1 \wedge dx_2, \\
&\quad (C_{32}C_{13} - C_{33}C_{12}) dx_2 \wedge dx_3 + (C_{33}C_{11} - C_{31}C_{13}) dx_3 \wedge dx_1 + \\
&\quad (C_{31}C_{12} - C_{32}C_{11}) dx_1 \wedge dx_2, \\
&\quad (C_{12}C_{23} - C_{13}C_{22}) dx_2 \wedge dx_3 + (C_{13}C_{21} - C_{11}C_{23}) dx_3 \wedge dx_1 + \\
&\quad (C_{11}C_{22} - C_{12}C_{21}) dx_1 \wedge dx_2) \\
&= f(T^{-1}x) \cdot \\
&\quad |T|^{-2} \begin{pmatrix} C_{22}C_{33} - C_{23}C_{32} & C_{23}C_{31} - C_{21}C_{33} & C_{21}C_{32} - C_{22}C_{31} \\ C_{32}C_{13} - C_{33}C_{12} & C_{33}C_{11} - C_{31}C_{13} & C_{31}C_{12} - C_{32}C_{11} \\ C_{12}C_{23} - C_{13}C_{22} & C_{13}C_{21} - C_{11}C_{23} & C_{11}C_{22} - C_{12}C_{21} \end{pmatrix} \begin{pmatrix} dx_2 \wedge dx_3 \\ dx_3 \wedge dx_1 \\ dx_1 \wedge dx_2 \end{pmatrix} \\
&= f(T^{-1}x)^t |T|^{-2} |T| T^t dX_2^t \\
&= |T|^{-1} (T f(T^{-1}x))^t dX_2^t \\
&= |T|^{-1} T f(T^{-1}x) \cdot dX_2.
\end{aligned}$$

Hence,

$$f(x) = |T|^{-1} T f(T^{-1}x)$$

and  $f$  is relatively equivariant. In case  $T \in \text{SU}_3$ ,  $|T| = 1$  and absolute equivariance occurs.  $\triangle$

Conversely, an absolute equivariant corresponds to a relatively invariant 2-form, with absolute invariance holding in case  $\mathcal{G} \subset \text{SU}_3$ .

For invariant exterior forms, there is a 2-variable “exterior Molien series”  $M(\Lambda^{\mathcal{G}})$  in which the variables  $s$  and  $t$  index respectively the rank of the form and the polynomial degree:

$$M(\Lambda^{\mathcal{G}}) = \sum_{p=0}^n \left( \sum_{m=0}^{\infty} \left( \dim \Lambda_p(\mathbf{C}^n)_{\mathcal{G}_m}^{\mathcal{G}} \right) t^m \right) s^p$$

where  $\Lambda_p(\mathbf{C}^n)_{\mathcal{G}_m}^{\mathcal{G}}$  are the  $\mathcal{G}$ -invariant homogeneous  $p$ -forms of degree  $m$ . ([Benson 1993, p. 62] or [Smith 1995, pp. 265ff]) Projection of  $\Lambda_p(\mathbf{C}^n)_m$  onto  $\Lambda_p(\mathbf{C}^n)_{\mathcal{G}_m}^{\mathcal{G}}$  yields the analogue to Molien’s theorem.

**Theorem 2** *The exterior Molien series for a finite group action  $\mathcal{G}$  is given by the generating function*

$$M(\Lambda^{\mathcal{G}}) = \frac{1}{|\mathcal{G}|} \sum_{\mathcal{C}_T \subset \mathcal{G}} |\mathcal{C}_T| \frac{\det(I + s T^{-1})}{\det(I - t T^{-1})}$$

where  $\mathcal{C}_T$  are conjugacy classes.

As in the case of invariants, this result produces

**Proposition 6** *For  $\mathcal{V}_{3,360}$ , the exterior Molien series is given by*

$$\begin{aligned} M(\Lambda^{\mathcal{V}_{3,360}}) &= \frac{1+t^{45}+(t^5+t^{11}+t^{20}+t^{26}+t^{29}+t^{44})s+(t+t^{16}+t^{19}+t^{25}+t^{34}+t^{40})s^2+(1+t^{45})s^3}{(1-t^6)(1-t^{12})(1-t^{30})} \\ &= 1 + t^6 + 2t^{12} + 2t^{18} + 3t^{24} + 4t^{30} + \dots \\ &\quad + (t^5 + 2t^{11} + 3t^{17} + t^{20} + 4t^{23} + 2t^{26} + 6t^{29} + \dots)s \\ &\quad + (t + t^7 + 2t^{13} + t^{16} + 3t^{19} + t^{22} + 5t^{25} + 2t^{28} + \dots)s^2 \\ &\quad + (1 + t^6 + 2t^{12} + 2t^{18} + 3t^{24} + 4t^{30} + \dots)s^3. \end{aligned}$$

### 3.2 A query on finite reflection groups

For a reflection group  $\mathcal{G}$  that acts on  $\mathbf{C}^n$  the  $n$  basic invariant 0-forms are algebraically independent. [Shephard and Todd 1954, pp. 282ff] Multiplication of an invariant  $p$ -form  $\alpha$  of degree  $\ell$  by an invariant 0-form  $F$  of degree  $m$  promotes  $\alpha$  to an invariant  $p$ -form  $F \cdot \alpha$  of degree  $\ell + m$ . In the series  $M(\Lambda^{\mathcal{H}})$  for a subgroup  $\mathcal{H} \subset \mathcal{G}$  the contribution of the free algebra generated by the basic 0-forms disappears upon division of  $M(\Lambda^{\mathcal{H}})$  by  $M(\Lambda_0^{\mathcal{G}}) = M(\mathbf{C}[x]^{\mathcal{G}})$ . The resulting *polynomial* in two variables displays the degrees of the generating  $\mathcal{H}$ -invariant forms. In the cases of the 0,  $n$ -forms, which have identical series, what remains are the terms corresponding to non-constant polynomials that are  $\mathcal{H}$ -invariant but not  $\mathcal{G}$ -invariant.

**Proposition 7** *For the Valentiner group, the “exterior Molien quotient” is*

$$M(\Lambda^{\mathcal{V}_{3,360}})/M(\Lambda_0^{\mathcal{V}_{3,360}}) = (1+t^{45}) + (t^5+t^{11}+t^{20}+t^{26}+t^{29}+t^{44})s + (t+t^{16}+t^{19}+t^{25}+t^{34}+t^{40})s^2 + (1+t^{45})s^3.$$

Notice the duality in degree 45 between 0 and 3-forms:

$$\begin{aligned} s^0 : & \quad 1 = t^0 & t^{45} \\ s^3 : & \quad t^{45} & 1 = t^0 \end{aligned}$$

and between 1 and 2-forms:

$$\begin{aligned} s^1 : & \quad t^5 & t^{11} & t^{20} & t^{26} & t^{29} & t^{44} \\ s^2 : & \quad t^{40} & t^{34} & t^{25} & t^{19} & t^{16} & t. \end{aligned}$$

By uniqueness, up to scalar multiplication, of the 3-form  $X \cdot vol$  associated with the 45 complex planes of reflection, the exterior product of “dual” forms must yield a multiple of this form.

This duality between invariant  $p$  and  $(n-p)$ -forms also appears in the *ternary* icosahedral group  $\mathcal{I}_{60}$ . The generating 0-forms for the full reflection group  $\mathcal{I}_{2,60}$  have degrees<sup>26</sup> 2, 6, and 10. From the discussion of the Valentiner conics the invariant of degree 2 is familiar, while those of

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<sup>26</sup>See line 23 of the table in [Shephard and Todd 1954, p. 301].



degrees 6 and 10 are products of linear forms that  $\mathcal{I}_{60}$  preserves. Here, the duality occurs in degree 15 which is the number of reflection-planes for the icosahedron:

$$\begin{aligned} M(\Lambda^{\mathcal{I}_{60}})/M(\Lambda_0^{\mathcal{I}_{2 \cdot 60}}) &= (1 + t^{15}) + (t + t^5 + t^6 + t^9 + t^{10} + t^{14})s + \\ &\quad (t + t^5 + t^6 + t^9 + t^{10} + t^{14})s^2 + (1 + t^{15})s^3. \end{aligned}$$

Which finite reflection groups have series with this property? Is this duality connected to that described by [Orlik and Terao 1992, p. 286]?

### 3.2.1 On $\mathcal{V}$ -equivariance and special orbits

Suppose  $a$  is fixed by an element  $T \in \mathcal{G}$ . Since a  $\mathcal{G}$ -equivariant map  $f : \mathbf{CP}^n \rightarrow \mathbf{CP}^n$  satisfies

$$f(a) = f(Ta) = Tf(a),$$

$T$  also fixes  $f(a)$ . Hence, special orbits map to special orbits. For the Valentiner action, the points fixed by an involution are a 45-point and its line while the 3, 4, 5-fold fixed points come in triples. Figure 9 summarizes the matter. Thus, under a  $\mathcal{V}$ -equivariant  $f$ , a 45-line  $\mathcal{L}_{\overline{abcd}}$  maps either to itself or to its point  $p_{\overline{abcd}}$ . In the former case,  $f$  preserves the pair of 90-points  $\{p_{\overline{abcd}_1}, p_{\overline{abcd}_2}\}$ . Since  $\mathcal{O}_{45}$  cannot map to  $\mathcal{O}_{90}$ ,  $f$  must fix the 45-points.<sup>27</sup> Concerning the 36-72 triples the matter stands just as in the case of the 45-90 points so that  $f$  either fixes or exchanges the 72-points  $p_{\overline{ab}_1}, p_{\overline{ab}_2}$  and fixes the 36-points  $p_{\overline{ab}}$ . What about a triple of 60 points? Symmetry forces  $f$  to permute the three points. Since the Valentiner group does not distinguish between 60-points, an equivariant action must fix the orbit pointwise.

## 3.3 The degree 16 map

Returning to the exterior Molien series for  $\mathcal{V}_{3 \cdot 360}$ , the coefficient of  $t^m s^2$  gives the dimension of the space of degree  $m$  equivariants. The series in  $t$  begins

$$t + t^7 + 2t^{13} + t^{16} + \dots$$

The first term  $t$  is due to the identity<sup>28</sup> map while  $t^7$  occurs through promotion of the identity to the degree 7 map  $F \cdot \text{id}$ . Two dimensions worth of invariants in degree 12 account for the  $2t^{13}$  term. The occurrence in degree 16 of the first non-trivial equivariant finds explanation in exterior algebra. Since exterior differentiation and multiplication preserve invariance, the 2-form  $dF \wedge d\Phi$  is invariant and hence, corresponds to an equivariant map whose coordinate functions are given by the coefficients of the 2-form basis

$$\{dx_2 \wedge dx_3, dx_3 \wedge dx_1, dx_1 \wedge dx_2\}.$$

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<sup>27</sup>The possibility of  $f$ 's being projectively undefined at  $p_{\overline{abcd}}$  is real. See below for a case-study.

<sup>28</sup>Although the identity map is always absolutely equivariant, its 2-form counterpart

$$x_1 dx_2 \wedge dx_3 + x_2 dx_3 \wedge dx_1 + x_3 dx_1 \wedge dx_2,$$

will not be absolutely invariant if the group is not in  $\text{SU}_3$ .

**Proposition 8** *Up to scalar multiplication, the unique  $\mathcal{V}_{3,360}$ -invariant 2-form of degree 16 is*

$$dF(x) \wedge d\Phi(x) = (\nabla F(x) \times \nabla \Phi(x)) \cdot dX_2.$$

*Consequently, the unique degree 16  $\mathcal{V}$ -equivariant is*

$$\psi_{16}(x) = \nabla F(x) \times \nabla \Phi(x).$$

Here,  $\nabla$  is a “formal” gradient  $\nabla F = (\frac{\partial F}{\partial x_1}, \frac{\partial F}{\partial x_2}, \frac{\partial F}{\partial x_3})$  and ‘ $\times$ ’ is the cross-product. Geometrically,  $\psi$  associates with a point  $y \in \mathbf{CP}^2$  the intersection of the pair of lines

$$\psi(y) = \{\nabla F(y) \cdot x = 0\} \cap \{\nabla \Phi(y) \cdot x = 0\}.$$

Order	Number and Type	Notation
2-fold:	one 45-point, one 45-line	$\{p_{\overline{abcd}}, \mathcal{L}_{\overline{abcd}}\}$
3-fold:	three 60-points	$\{p_{a*def}, p_{b*def}, p_{c*def}\}$
3-fold:	three $\overline{60}$ -points	$\{p_{\overline{ade*f}}, p_{\overline{b*def}}, p_{\overline{c*def}}\}$
4-fold:	one 45-point, two 90-points	$\{p_{\overline{abcd}}, p_{\overline{abcd}_1}, p_{\overline{abcd}_2}\}$
5-fold:	one 36-point, two 72-points	$\{p_{\overline{ab}}, p_{\overline{ab}_1}, p_{\overline{ab}_2}\}.$

Figure 9: Fixed points of  $\mathcal{V}$

### 3.3.1 Dynamics of the 16-map

Given a generic point  $a$  that lies on *just one* 45-line  $\mathcal{L}_{\overline{abcd}}$  the lines

$$\{\nabla F(a) \cdot x = 0\} \text{ and } \{\nabla \Phi(a) \cdot x = 0\}$$

each pass through the point  $p_{\overline{abcd}}$ . Thus,  $\psi$  collapses  $\mathcal{L}_{\overline{abcd}}$  to its companion point. Taking  $p_{\overline{abcd}} = [1, 0, 0]$  and  $\mathcal{L}_{\overline{abcd}} = \{x_1 = 0\}$ ,

$$\psi(a) = [\psi_1(a), 0, 0]$$

where  $\psi_1(a)$  is degree 16 in the homogeneous coordinates  $a = [0, a_2, a_3]$ . The 16 roots of  $\psi_1(a)$  correspond to the 16 points where  $\mathcal{L}_{\overline{abcd}}$  intersects the 44 remaining 45-lines. These occur at the 36, 45, 60,  $\overline{60}$ -points of which there are four each on  $\mathcal{L}_{\overline{abcd}}$ . The “blowing-down” of the 45-lines forces the “blowing-up” of their intersections: To which 45-point does the intersection go?

Their collapsing behavior makes  $\psi$  critical on the 45-lines. Since the Jacobian determinant  $|J_\psi|$  has degree  $3 \cdot (16 - 1) = 45$ ,  $\{X = 0\}$  is exactly the critical set. Thus, the 45-lines are superattracting. But, in approaching  $\mathcal{L}_{\overline{abcd}}$  a trajectory inevitably gets carried near the point  $p_{\overline{abcd}}$  which is blowing up onto  $\mathcal{L}_{\overline{abcd}}$ ; this means that the  $\mathbf{CP}^1$  of directions through  $p_{\overline{abcd}}$  maps,

by the agency of the Jacobian transformation  $J_\psi$ , to points on  $\mathcal{L}_{\overline{abcd}}$ .<sup>29</sup> (Conversely,  $J_\psi$  associates a point on  $\mathcal{L}_{\overline{abcd}}$  with a direction through  $p_{\overline{abcd}}$ .) The next two iterations of  $\psi$  first return the trajectory to the vicinity of  $p_{\overline{abcd}}$  and then send it back to  $\mathcal{L}_{\overline{abcd}}$  *somewhere*.

Observation of trajectories that start near  $p_{\overline{abcd}}$  or  $\mathcal{L}_{\overline{abcd}}$  reveals a rapid attraction on every other iteration. What about the 3 and 5-fold intersections of 45-lines? Are they attracting? While these points do indeed blow-up, their “image” under  $\psi$  is not a curve that blows-down to the point; such is the sole propriety of the 45-points. Hence,  $\psi$  draws a generic point near a 60,  $\overline{60}$ , 36-point into one of three or five point-line cycles.

This “every-other” dynamics poses a problem: which every-other iterate do we watch? For some initial postions, the even iterates converge to a 45-point while the trajectory spends the odd times around the corresponding 45-line. For others, the process is reversed. Moreover, experiment indicates that the dynamics at the 45-line eventually settles down; trajectories end up at one of the 45-points on the line. Hence, the iteration outputs a pair of 45-points each of which lies on the line of the other. But, there are, for each 45-point, four possible pairs of this sort so that taking every other iterate amounts to a neglect of information.

The dynamics of  $\psi$  appears to come down to what takes place on the critical 45-lines. Given a point  $x$  on a 45-line  $\mathcal{L}_{\overline{abcd}}$ , the derivative  $J_\psi$  associates with  $x$  a  $\mathbf{CP}^1$  through  $p_{\overline{abcd}}$ , namely, the image  $\mathcal{L}_\psi(x)$  of  $J_\psi(x)$ . In turn,  $J_\psi(p_{\overline{abcd}})$  sends the line  $\mathcal{L}_\psi(x)$  to the point  $[J_\psi(p_{\overline{abcd}})]\mathcal{L}_\psi(x)$  on  $\mathcal{L}_{\overline{abcd}}$ . The degree 15 map

$$x \rightarrow [J_\psi(p_{\overline{abcd}})]\mathcal{L}_\psi(x)$$

gives  $\psi$  on  $\mathcal{L}_{\overline{abcd}}$ . In Appendix A, the dynamics on a 45-line appears in a basins-of-attraction plot.<sup>30</sup> Here, the union of the patches in a given color indicates a basin of attraction for one of the four attracting 45-points. The dark regions correspond to points that the 45-points *might* not attract. Does this set have interior or positive measure?

On a 45-line the map is not critically finite; in particular, its critical points are not periodic.<sup>31</sup> Indeed, there might be a wandering critical point there.<sup>32</sup> Hence, establishing convergence almost everywhere would be difficult should the map even possess this property. Moreover, this behavior hardly reveals the geometric elegance whose prospective discovery motivates the present enterprise.

Unlike the 1-dimensional case of the icosahedral group for which the non-trivial equivariant of lowest degree provides an elegant dynamical system<sup>33</sup> for the purposes of solving the quintic, the higher dimensional Valentiner action fails to bear similar fruit. The failure occurs in spite of the 16-map’s being obtained by a procedure analogous to that employed by

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<sup>29</sup>See below.

<sup>30</sup>See figure 10.

<sup>31</sup>On critical finiteness, see [Fornaess and Sibony 1994, pp. 223ff.].

<sup>32</sup>Correspondence with Curt McMullen.

<sup>33</sup>See [Doyle and McMullen 1989, pp. 152-3].

[Doyle and McMullen 1989] in producing the degree-11 icosahedral map:

<u>1 dimension</u>	<u>2 dimensions</u>
$F$ : $\mathcal{I}$ -invariant of degree 12	$F, \Phi$ : $\mathcal{V}$ -invariants of degrees 6, 12
$f_{11} = \begin{vmatrix} \hat{x} & \hat{y} \\ F_x & F_y \end{vmatrix}$	$\psi_{16} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ F_x & F_y & F_z \\ \Phi_x & \Phi_y & \Phi_z \end{vmatrix}$

where  $\hat{x}, \hat{y}, \hat{z}$  represent unit coordinate vectors. However, in one dimension *all*  $\mathcal{I}$ -symmetric maps arise as combinations of three others each of which are constructed in the manner of  $f_{11}$  but with different basic invariants standing in for  $F$ . A richer stock of equivariants inhabits the Valentiner waters. In contrast, the next higher degree offers promise as well as a bit of mystery.

### 3.4 A family of 19-maps

The Molien series for Valentiner equivariants

$$t + t^7 + 2t^{13} + t^{16} + 3t^{19} \dots$$

specifies three dimensions worth of maps in degree 19 of which two are due to promotion of the identity by degree 18 invariants. Hence, there are, as the exterior Molien quotient<sup>34</sup> indicates, non-trivial  $\mathcal{V}$ -symmetric maps in degree 19. How do these arise? Since there is no apparent exterior algebraic means of producing such a map, the more practical matter of computing them takes priority.<sup>35</sup>

#### 3.4.1 19 = 64 - 45

Multiplication of a degree 19 equivariant  $f$  by  $X_{45}$  elevates  $f$  to the 14 dimensional space of 64-maps. There are 14 ways of promoting the maps  $\psi_{16}$ ,  $\phi_{34}$ , and  $f_{40}$  to degree 64:

- 1) 7 dimensions of degree 48 invariants to promote  $\psi_{16} = \nabla F \times \nabla \Phi$
- 2) 4 dimensions of degree 30 invariants to promote  $\phi_{34} = \nabla F \times \nabla \Psi$
- 3) 3 dimensions of degree 24 invariants to promote  $f_{40} = \nabla \Phi \times \nabla \Psi$ .

**Proposition 9** *These 14 maps span the space of degree 64 equivariants.*

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<sup>34</sup>See Section 3.2.

<sup>35</sup>Peter Doyle and Anne Shepler appear to have found a “differential” method of generating all invariant 2-forms. The geometric and dynamic consequences remain to be explored.

*Proof.* If not, then for some  $\alpha_k$  not all 0,

$$\begin{aligned} (\alpha_1 F^8 + \dots + \alpha_7 F \Phi \Psi) dF \wedge d\Phi + (\alpha_8 F^5 + \dots + \alpha_{11} \Psi) dF \wedge d\Psi + \\ (\alpha_{12} F^4 + \alpha_{13} F^2 \Phi + \alpha_{14} \Phi^2) d\Phi \wedge d\Psi = 0. \end{aligned}$$

Since  $dF \wedge d\Phi \wedge d\Psi = \beta X \cdot \text{vol}$  for some constant  $\beta \neq 0$ , either

$$(\alpha_1 F^8 + \dots + \alpha_7 F \Phi \Psi) \beta X = 0,$$

$$(\alpha_8 F^5 + \dots + \alpha_{11} \Psi) \beta X = 0,$$

or

$$(\alpha_{12} F^4 + \alpha_{13} F^2 \Phi + \alpha_{14} \Phi^2) \beta X = 0.$$

In at least one case, there is a non-zero  $\alpha_k$ . But, then the invariants  $F, \Phi, \Psi$  are not algebraically independent, contrary to the theory of complex reflection groups. [Shephard and Todd 1954, p. 282]  $\triangle$

In this event,  $X \cdot f$  is a combination of maps whose computation is straightforward.

Reasoning in the other direction, a 64-map

$$f_{64} = F_{48} \cdot \psi_{16} + F_{30} \cdot \phi_{34} + F_{24} \cdot f_{40}$$

that “vanishes” on the 45-lines—i.e., each coordinate function of  $f_{64}$  vanishes—must have a factor of  $X$ . The quotient is a degree 19 equivariant

$$f_{19} = \frac{f_{64}}{X_{45}}.$$

Arranging for the vanishing of  $f_{64}$  on the 45-lines requires consideration of only one line; symmetry tends to the remaining 44. Forcing  $f_{64}$  to vanish at 12 “independent” points<sup>36</sup> on a 45-line yields a 2-parameter family of 64-maps each member of which vanishes on  $\{X = 0\}$ . The two *inhomogeneous* parameters reflect the three dimensions (i.e., homogeneous parameters) of degree 19  $\mathcal{V}$ -equivariants. In  $\text{bub}_{22}$  coordinates, setting these two parameters equal to 0 and normalizing the coefficients yields

$$\begin{aligned} f_{64}(y) = & \left[ 10 F(y)^6 \Phi(y) + 100 F(y)^4 \Phi(y)^2 + 45 F(y)^2 \Phi(y)^3 + \right. \\ & 156 \Phi(y)^4 + 39 F(y)^3 \Psi(y) + 51 F(y) \Phi(y) \Psi(y) \left. \right] \cdot \psi(y) - \\ & 27 \Psi(y) \cdot \phi(y) + 54 \Phi(y)^2 \cdot f(y). \end{aligned} \tag{2}$$

The 2-parameter family of non-trivial 19-maps is then

$$g_{19}(y; a, b) = f_{19}(y) + \left( a F(y)^3 + b F(y) \Phi(y) \right) \cdot y. \tag{3}$$

Are any of these maps dynamically “special”? Indeed, what might it mean to be special in this sense?

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<sup>36</sup>The points are *independent* in the sense that the 12 resulting linear conditions in the 14 undetermined coefficients of  $f_{64}$  are independent.

### 3.4.2 Extended symmetry in degree 19

Since  $f_{19}$  is a non-trivial  $\bar{\mathcal{V}}_{2,360}$ -equivariant—note the integer coefficients in (2), each member of the two *real* parameter family

$$f_{19} = F(a B_{12} + \bar{a} U_{12}) \cdot \text{id}$$

is impartial towards the two systems of conics and so, enjoys the additional symmetry. Here  $B_{12} = \prod C_{\bar{k}}$  and  $U_{12} = \prod C_k$  are the degree 12 invariants given by the product of the respective six conic forms. To honor the doubled symmetry a member of this family must preserve each  $\mathcal{R}_{\bar{a}b}$  that  $\text{bub}_{\bar{a}b}$  fixes point-wise.

## 3.5 The 19-map

### 3.5.1 The icosahedron again

An intriguing aspect of the 19-maps is the degree itself. Since 19 is one of the special equivariant numbers<sup>37</sup> for the binary icosahedral group, there arises the prospect of finding a  $\mathcal{V}$ -equivariant that restricts to self-mappings of the conics. By symmetry, an equivariant that fixes<sup>38</sup> a conic of a given system also fixes the other five. Might there be a map that preserves each of the 12 conics?

From [Doyle and McMullen 1989, p. 163] comes a geometric description of the canonical degree 19 icosahedral mapping of the round Riemann sphere: stretch each face  $F$  over the 19 faces in the complement of the face antipodal to  $F$  while making a half-turn in order to place the three vertices and edges of  $F$  on the three antipodal vertices and edges. By symmetry, the 20 face-centers are fixed and repelling. Since the resulting map is critical only at the 12 period-2 vertices, the map has nice dynamics. [Doyle and McMullen 1989, p. 156] A conic-fixing  $\mathcal{V}$ -equivariant of degree 19 would, when restricted to a conic  $\mathcal{C}_{\bar{a}}$  or  $\mathcal{C}_b$ , give the unique map in degree 19 with binary icosahedral symmetry. This would determine its effect on the special icosahedral orbits:

- 1) fix the face-centers:  $\overline{60}$ , 60-points of the appropriate system
- 2) exchange antipodal vertices: pairs of 72-points
- 3) exchange antipodal edge-midpoints: intersections of conics  $\mathcal{C}_{\bar{a}}$ ,  $\mathcal{C}_b$  with a 45-line indexed by  $\bar{a}$  and  $b$ .

General  $\mathcal{V}$ -equivariants satisfy<sup>39</sup> condition 1). Since

$$\mathcal{O}_{72} = \{F = 0\} \cap \{\Phi = 0\},$$

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<sup>37</sup>See [Doyle and McMullen 1989, p. 166].

<sup>38</sup>Here, ‘fix’ means “set-wise”.

<sup>39</sup>While some maps can be set to blow-up at the 60-points, such a circumstance is rare.

the image of a 72-point under each  $g_{19}$  in (3) is the same as that under  $f_{19}$ . Propitiously,  $f_{19}$  exchanges pairs of 72-points. Finally, a 19-map cannot blow-down a 45-line; this would make it critical there—a condition precluded by the critical set's being an invariant of degree  $54 = 3(19 - 1)$ . Consequently, each 45-line must map to itself so that arranging for 3) costs one parameter for each system of conics. Fortunately, there are two parameters to spend and their expenditure purchases a canonical  $\mathcal{V}$ -equivariant  $h_{19}$  that maps each of the 12 conics onto itself.<sup>40</sup>

**Fact 3** *There is a unique degree-19  $\mathcal{V}$ -equivariant  $h_{19}$  that preserves each of the 12 icosahedral conics.*

By favoring neither system of conics,  $h_{19}$  possesses bub-symmetry and so, self-maps each of the  $\mathcal{R}_{\bar{a}b}$ . Expressing the family of 19-maps by

$$g_{19} = h_{19} + F(a B_{12} + b U_{12}) \cdot \text{id}$$

makes evident the 1-parameter collections that fix the barred ( $b = 0$ ) and the unbarred ( $a = 0$ ) conics.

Unlike the 16-map  $\psi_{16}$ ,  $h_{19}$  does not blow up somewhere.

**Proposition 10** *The conic-preserving map  $h_{19}$  is holomorphic on  $\mathbf{CP}^2$ .*

*Proof.* By equivariance, the set of points on which  $h_{19}$  blows up is empty or a union of  $\mathcal{V}$ -orbits. Direct calculation yields that

$$h_{19}(p) \neq 0$$

for

$$p \in \mathcal{O}_{36} \cup \mathcal{O}_{45} \cup \mathcal{O}_{60} \cup \mathcal{O}_{\overline{60}} \cup \mathcal{O}_{90}.$$

The remaining possibilities are that  $h_{19} = 0$  on a 180 or 360 point orbit.

First, take the case of a 180 point orbit and recall that each such point belongs to one 45-line. Also, let

$$h_{19} = [h_1, h_2, h_3].$$

Since  $h_{19}$  preserves each 45-line  $\mathcal{L}$ , the only way that  $h_1 = h_2 = h_3 = 0$  is for the coordinates of the restriction  $h_{19}|_{\mathcal{L}}$  to have a common factor. In bub<sub>22</sub> coordinates  $\mathcal{L}_{\overline{12}12} = \{y_1 - y_2 = 0\}$  so that

$$h_{19}|_{\mathcal{L}_{\overline{12}12}} = [f, f, g].$$

But, the resultant of  $f$  and  $g$  does not vanish. Hence,  $f$  and  $g$  do not have a common factor.

Finally, suppose that  $h_{19} = 0$  at a 360 point orbit and that  $[0,0,1]$  is a 36-point  $p_{36}$ . Since

$$|\{h_1 = 0\} \cap \{h_2 = 0\} \cap \{h_3 = 0\}| \leq 19 \cdot 19 = 361,$$

there is only one member of  $h_{19}^{-1}(p_{36})$  in  $\mathbf{CP}^2$ . Of course, this holds for every 36-point of which there are four on a 45-line  $\mathcal{L}$ . Moreover,  $h_{19}(p_{36}) = p_{36}$ . Thus, the one-dimensional rational

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<sup>40</sup>See below for an explicit expression.

map  $h_{19}|_{\mathcal{L}}$  has four exceptional points—a state of affairs that requires the map to be degree one. [Beardon 1991, p. 52] Since the restricted map is not degree one,  $h_{19}$  does not vanish at a 360 point orbit. (Indeed, no degree 19 equivariant can blow up a 360 point orbit.)  $\triangle$

**Fact 4** *The conic-fixing equivariant has the  $\text{bub}_{\overline{22}}$  expression:*

$$\begin{aligned} h_{19}(y) = 1620 F(y)^3 \cdot [y_1, y_2, y_3] + f_{19}(y) = \\ [-3591 y_1^{15} y_2^4 - 5263 y_1^{10} y_2^9 + 9747 y_1^5 y_2^{14} - 81 y_2^{19} + 17955 y_1^{12} y_2^6 y_3 + 10260 y_1^7 y_2^{11} y_3 - 7695 y_1^2 y_2^{16} y_3 - \\ 107730 y_1^{14} y_2^3 y_3^2 - 74385 y_1^9 y_2^8 y_3^2 + 161595 y_1^4 y_2^{13} y_3^2 - 969570 y_1^{11} y_2^5 y_3^3 + 1292760 y_1^6 y_2^{10} y_3^3 - 46170 y_1 y_2^{15} y_3^3 - \\ 2346975 y_1^8 y_2^7 y_3^4 - 807975 y_1^3 y_2^{12} y_3^4 - 3587409 y_1^{10} y_2^4 y_3^5 + 10277442 y_1^5 y_2^9 y_3^5 + 13851 y_2^{14} y_3^5 - 969570 y_1^{12} y_2 y_3^6 - \\ 3986010 y_1^7 y_2^6 y_3^6 - 1939140 y_1^2 y_2^{11} y_3^6 - 5263380 y_1^9 y_2^3 y_3^7 - 28117530 y_1^4 y_2^8 y_3^7 + 831060 y_1^{11} y_3^8 + \\ 2423925 y_1^6 y_2^3 y_3^8 + 4363065 y_1 y_2^{10} y_3^8 - 24931800 y_1^8 y_2^2 y_3^9 + 43630650 y_1^3 y_2^7 y_3^9 - 31123197 y_1^5 y_2^4 y_3^{10} + \\ 9598743 y_2^9 y_3^{10} + 14959080 y_1^7 y_2 y_3^{11} + 23269680 y_1^4 y_2^3 y_3^{12} - 26178390 y_1^6 y_3^{13} + 52356780 y_1 y_2^5 y_3^{13} + \\ 18698850 y_1^3 y_2^2 y_3^{14} + 20194758 y_2^2 y_3^{15} + 22438620 y_1^2 y_2 y_3^{16} + 7479540 y_1 y_3^{18}, \\ -81 y_1^{19} + 9747 y_1^{14} y_2^5 - 5263 y_1^9 y_2^{10} - 3591 y_1^4 y_2^{15} - 7695 y_1^{16} y_2^2 y_3 + 10260 y_1^{11} y_2^7 y_3 + 17955 y_1^6 y_2^{12} y_3 + \\ 161595 y_1^{13} y_2^4 y_3^2 - 74385 y_1^8 y_2^9 y_3^2 - 107730 y_1^3 y_2^{14} y_3^2 - 46170 y_1^{15} y_2 y_3^3 + 1292760 y_1^{10} y_2^6 y_3^3 - 969570 y_1^5 y_2^{11} y_3^3 - \\ 807975 y_1^{12} y_2^3 y_3^4 - 2346975 y_1^7 y_2^8 y_3^4 + 13851 y_1^{14} y_3^5 + 10277442 y_1^9 y_2^5 y_3^5 - 3587409 y_1^4 y_2^{10} y_3^5 - \\ 1939140 y_1^{11} y_2^2 y_3^6 - 3986010 y_1^6 y_2^7 y_3^6 - 969570 y_1 y_2^{12} y_3^6 - 28117530 y_1^8 y_2^4 y_3^7 - 5263380 y_1^3 y_2^9 y_3^7 + \\ 4363065 y_1^{10} y_2 y_3^8 + 2423925 y_1^5 y_2^6 y_3^8 + 831060 y_2^{11} y_3^8 + 43630650 y_1^7 y_2^3 y_3^9 - 24931800 y_1^2 y_2^8 y_3^9 + \\ 9598743 y_1^9 y_3^{10} - 31123197 y_1^4 y_2^5 y_3^{10} + 14959080 y_1 y_2^7 y_3^{11} + 23269680 y_1^3 y_2^4 y_3^{12} + 52356780 y_1^5 y_2 y_3^{13} - \\ 26178390 y_2^6 y_3^{13} + 18698850 y_2^2 y_3^{14} + 20194758 y_1^4 y_3^{15} + 22438620 y_1 y_2^2 y_3^{16} + 7479540 y_2 y_3^{18}, \\ -1026 y_1^{17} y_2^2 - 3078 y_1^{12} y_2^7 - 3078 y_1^7 y_2^{12} - 1026 y_1^2 y_2^{17} - 5130 y_1^{14} y_2^4 y_3 + 113240 y_1^9 y_2^9 y_3 - 5130 y_1^4 y_2^{14} y_3 + \\ 3078 y_1^{16} y_2 y_3^2 - 272916 y_1^{11} y_2^6 y_3^2 - 272916 y_1^6 y_2^{11} y_3^2 + 3078 y_1 y_2^{16} y_3^2 + 215460 y_1^{13} y_2^3 y_3^3 + 687420 y_1^8 y_2^8 y_3^3 + \\ 215460 y_1^3 y_2^{13} y_3^3 + 4617 y_1^{15} y_3^4 + 937251 y_1^{10} y_2^5 y_3^4 + 937251 y_1^5 y_2^{10} y_3^4 + 4617 y_2^{15} y_3^4 + 290871 y_1^{12} y_2^2 y_3^5 + \\ 4813992 y_1^7 y_2^7 y_3^5 + 290871 y_2^{12} y_3^5 - 1454355 y_1^9 y_2^4 y_3^6 - 1454355 y_1^4 y_2^9 y_3^6 + 2520882 y_1^{11} y_2 y_3^7 + 8812314 y_1^6 y_2^6 y_3^7 + \\ 2520882 y_1 y_2^{11} y_3^7 + 19876185 y_1^8 y_2^3 y_3^8 + 19876185 y_1^3 y_2^8 y_3^8 - 2036097 y_1^{10} y_3^9 + 5623506 y_1^5 y_2^5 y_3^9 - 2036097 y_2^{10} y_3^9 + \\ 5235678 y_1^7 y_2^3 y_3^{10} + 5235678 y_2^2 y_3^{10} + 37813230 y_1^4 y_2^4 y_3^{11} - 2617839 y_1^6 y_2 y_3^{12} - 2617839 y_1 y_2^6 y_3^{12} - \\ 2908710 y_3^3 y_3^{13} + 6357609 y_1^5 y_3^{14} + 6357609 y_2^5 y_3^{14} - 5983632 y_2^2 y_2 y_3^{15} - 4487724 y_1 y_2 y_3^{17} - 1023516 y_3^{19}]. \end{aligned}$$

### 3.5.2 Dynamical behavior

The discovery of  $h_{19}$  supplies the unique degree-19  $\mathcal{V}$ -equivariant that self-maps, in addition to the 45-lines and the 36  $\text{bub-}\mathbf{RP}^2$ s, the 12 conics. The dynamics *on* each conic is well-understood.<sup>41</sup>

**Proposition 11** *Under  $h_{19}$ , the trajectory of almost any point on an icosahedral conic tends to an antipodal pair of the superattracting vertices.*

Moreover, the conics themselves are attracting.

**Proposition 12** *The Jacobian  $J_{h_{19}}$  has rank one at the superattracting 72-points. Thus,  $h_{19}$  attracts on a full  $\mathbf{CP}^2$  neighborhood of such a point. Furthermore, the Fatou components of the restricted map  $h_{19}|_{\mathcal{C}_{\bar{\alpha}}}$  are the intersections with  $\mathcal{C}_{\bar{\alpha}}$  of Fatou components of the map on  $\mathbf{CP}^2$ .*

<sup>41</sup>A basins-of-attraction plot appears in Figure 11.



Is this attracting behavior of the conics pervasive in the measure-theoretic sense? What about the “restricted” dynamics on the 45-lines and  $\mathbf{RP}^2$ s?

Perhaps the place to begin is at a 72-point, say  $p_{\overline{11}_1}$ , which lies at the hub of much Valentiner activity. Passing through  $p_{\overline{11}_1}$  are many special objects. To enumerate:

- 1) the pair  $\{\mathcal{C}_1, \mathcal{C}_{\overline{1}}\}$  of conics, which meet tangentially
- 2) the 36-line  $\mathcal{L}_{\overline{11}}$ , which gives a 10-fold  $\mathcal{D}_5$  axis about which  $\mathcal{C}_1 \cup \mathcal{C}_{\overline{1}}$  “turns”
- 3) the 72-line  $\mathcal{L}_{\overline{11}_2}$ , which is stable under the cyclic half of the  $\mathcal{D}_5$  stabilizer of  $\mathcal{L}_{\overline{11}}$  and thereby tangent to  $\mathcal{C}_{\overline{1}}$  and  $\mathcal{C}_1$
- 4) the sixth-degree curve  $\{F = 0\}$ , which is tangent to  $\mathcal{L}_{\overline{11}}$
- 5) the twelfth-degree curve  $\{\Phi = 0\}$ , which is tangent to  $\mathcal{L}_{\overline{11}_2}$
- 6) the five  $\mathbf{RP}^2$ s  $\{\mathcal{R}_{\overline{22}}, \mathcal{R}_{\overline{34}}, \mathcal{R}_{\overline{43}}, \mathcal{R}_{\overline{56}}, \mathcal{R}_{\overline{65}}\}$ , each of which intersect  $\mathcal{C}_1$  and  $\mathcal{C}_{\overline{1}}$  only at  $p_{\overline{11}_1}$  and  $p_{\overline{11}_2}$ .

In addition, a 72-point situates itself at the intersection of two components  $\{F_6 = 0\}$  and  $\{G_{48} = 0\}$  of  $h_{19}$ ’s critical set.

**Fact 5** *The Jacobian determinant  $|J_{h_{19}}| = F_6 G_{48}$  where the invariant  $G_{48}$  distinguishes itself by lacking a  $\Psi$  term when decomposed into an expression in the basic invariants  $F, \Phi, \Psi$ :*

$$G_{48}(y) = -13718 \left[ 14 F(y)^8 + 180 F(y)^6 \Phi(y) + 1701 F(y)^4 \Phi(y)^2 + 3402 F(y)^2 \Phi(y)^3 + 5103 \Phi(y)^4 \right].$$

Since  $G_{48}$  is a polynomial in  $F$  and  $\Phi$  alone, its curve satisfies

$$\{F = 0\} \cap \{G_{48} = 0\} = \{F = 0\} \cap \{\Phi = 0\} = \mathcal{O}_{72}.$$

Furthermore, the special invariant structure of  $G_{48}$  has an alternative expression in terms of  $B_{12}$  and  $U_{12}$  alone. From the identities

$$\begin{aligned} B_{12}(y) &= 2(5 - \sqrt{15}i) \rho F(y)^2 - 2\sqrt{15}(1 + \sqrt{15}i) \rho^2 \Phi(y) \\ U_{12}(y) &= 2(5 + \sqrt{15}i) \rho^2 F(y)^2 - 2\sqrt{15}(1 - \sqrt{15}i) \rho \Phi(y) \end{aligned}$$

it follows that

$$\begin{aligned} G_{48}(y) &= 2^3 3^{24} 5^8 19^3 \left( -6(3 - \sqrt{15}i) \rho B_{12}(y)^4 + \right. \\ &\quad 4(32 + 3\sqrt{15}i) \rho^2 B_{12}(y)^3 U_{12}(y) - 333 B_{12}(y)^2 U_{12}(y)^2 + \\ &\quad \left. 4(32 - 3\sqrt{15}i) \rho B_{12}(y) U_{12}(y)^3 - 6(3 + \sqrt{15}i) \rho^2 U_{12}(y)^4 \right). \end{aligned}$$

Consequently, the degree 48 component of the critical set meets the conics *only* at the 72-points:

$$\begin{aligned}\{G_{48} = 0\} \cap \{B_{12} = 0\} &= \{G_{48} = 0\} \cap \{U_{12} = 0\} \\ &= \{B_{12} = 0\} \cap \{U_{12} = 0\} \\ &= \mathcal{O}_{72}\end{aligned}$$

Accounting for the multiplicity at these eighth-order intersections is the singularity of  $\{G_{48} = 0\}$  at  $\mathcal{O}_{72}$ , a result that follows directly from the invariant decomposition.

### 3.5.3 $\mathbf{RP}^2$ dynamics

On each of the five bub- $\mathbf{RP}^2$ s that are mutually tangent at  $p_{\overline{11}1}$  and  $p_{\overline{11}2}$ , these 72-points are superattracting for the restricted maps

$$h_{19}|_{\mathcal{R}_{\overline{ab}}}, \quad \overline{ab} = \overline{22}, \overline{34}, \overline{43}, \overline{56}, \overline{65}.$$

Are there attracting sites on  $\mathcal{R}_{\overline{ab}}$  other than the five pairs of 72-points? Are there sets of positive measure or open sets on which  $h_{19}|_{\mathcal{R}_{\overline{ab}}}$  fails to converge to a pair of 72-points? The experimental evidence strongly suggests that

- 1) the 72-points are the only attractors
- 2) there is no region with thickness or positive measure that remains outside of their influence
- 3) the set of 45-lines  $\{X = 0\}$  is repelling.

Appendix A exhibits basin plots of  $h_{19}|_{\mathcal{R}_{\overline{22}}}$  which, of course, is dynamically equivalent to each  $h_{19}|_{\mathcal{R}_{\overline{ab}}}$ . What significance does the  $\mathbf{RP}^2$ -dynamics hold for that on  $\mathbf{CP}^2$ ? Extensive trials<sup>42</sup> in  $\mathbf{CP}^2$  have not revealed behavior contrary to that observed on the  $\mathbf{RP}^2$ s.

### 3.5.4 Conjectures

**Conjecture 1** *The only attracting periodic points for  $h_{19}$  are the elements of  $\mathcal{O}_{72}$ . Moreover, the union of the basins of attraction for  $\mathcal{O}_{72}$  has full measure in  $\mathbf{CP}^2$ .*

The 45-lines present a problem in that they map to themselves but do not contain the 72-points.<sup>43</sup> The basin plot in Figure 12 reveals repelling behavior along the  $\mathbf{RP}^1$  where the  $\mathbf{RP}^2$  meets one of the 45-lines that bub $_{\overline{22}}$  fixes set-wise.

**Conjecture 2** *On the 45-lines  $h_{19}$  is repelling and, hence,  $\{X = 0\}$  resides in the Julia set  $J_{h_{19}}$ .*

Is  $J_{h_{19}}$  the closure of the backward orbit of the 45-lines?

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<sup>42</sup>At <http://math.ucsd.edu/~scrass> there are *Mathematica* notebooks and supporting files with which to iterate  $h_{19}$  and, from the output, to approximate a solution to a sixth-degree equation.

<sup>43</sup>Again, this is a feature peculiar to the 72-points. They form the only special  $\mathcal{V}$ -orbit that does not lie on the 45-lines.

## 4 Solving the Sextic

By means of various algebraic manipulations, a general sixth-degree polynomial reduces to a member of a 2-parameter family of “Valentiner resolvents”. Such a reduction requires the extraction of square and cube roots. Furthermore, a certain set of sextics transforms into a special 1-parameter collection of resolvents. These resolvents are especially suited for solution by an iterative algorithm that exploits Valentiner symmetry and symmetry-breaking.

### 4.1 General sixth-degree Valentiner resolvents

At the core of Klein’s program for equation-solving is the “form problem” relative to a particular action of a given equation’s symmetry group: for prescribed values  $a_1, \dots, a_n$  of the generating invariants  $F_1, \dots, F_n$  find a point  $p$  common to the inverse images  $F_1^{-1}(a_1), \dots, F_n^{-1}(a_n)$ . As with the quintic, solving the general sextic is tantamount to solving the corresponding form-problem. ([Fricke 1926, pp. 308-10] and [Coble 1911]) This circumstance has a “projectively-equivalent” formulation in terms of rational functions in the basic invariants. In the Valentiner setting this concerns

$$Y_1 = \alpha \frac{\Phi}{F^2} \quad Y_2 = \beta \frac{\Psi}{F^5}$$

where  $\alpha$  and  $\beta$  are chosen so that  $Y_1 = 1$  and  $Y_2 = 1$  at a 36-point.<sup>44</sup> Given values  $a_1$  and  $a_2$  of  $Y_1$  and  $Y_2$ , the task is to find a point  $z$  in  $\mathbf{CP}^2$  that belongs to the  $\mathcal{V}$ -orbit

$$Y_1^{-1}(a_1) \cap Y_2^{-1}(a_2).$$

Accordingly, the general 6-parameter sextic  $p(x)$  reduces to a resolvent that depends on the two parameters  $Y_1$  and  $Y_2$ . Such a reduction requires the extraction of a cube root [Fricke 1926, p. 285] in addition to the square root of  $p$ ’s discriminant. This cube root is a so-called “accessory irrationality”—its adjunction to the coefficient field does not reduce the galois group. The 1-to-3 correspondence between the projective and linear Valentiner groups  $\mathcal{V}$  and  $\mathcal{V}_{3,360}$  accounts for its appearance. In the 1-dimensional icosahedral case, the projective group lifts 1-to-2 to a linear group thereby producing the need for an accessory square root. [Klein 1956, pp. 172-3]

As for the *derivation* of a 2-parameter resolvent, the map

$$Y : \mathbf{CP}^2 - \{F(z) = 0\} \rightarrow (\mathbf{CP}^2 - \{F(z) = 0\})/\mathcal{V}$$

given by

$$Y(z) = [F(z)^3 \Phi(z), \Psi(z), F(z)^5] = [Y_1(z), Y_2(z), 1]$$

provides the  $\mathcal{V}$ -quotient of  $\mathbf{CP}^2 - \{F(z) = 0\}$  in that the fibers are  $\mathcal{V}$ -orbits. The exceptional status of the sixth-degree curve is due to its being the fiber above the single point  $[0, 1, 0]$ . Furthermore, under the icosahedral function

$$U_1(z) = \frac{C_1(z)^3}{F(z)}$$

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<sup>44</sup>In  $\text{bub}_{\overline{22}}$  coordinates  $\alpha = 1$  and  $\beta = 1/4$ .

a fiber  $Y^{-1}[a_1, a_2, 1]$  maps to six points

$$\left\{ U_{\bar{n}}(z) = \frac{C_{\bar{n}}(z)^3}{F(z)} \mid n = 1, \dots, 6 \right\}$$

where  $z \in Y^{-1}[a_1, a_2, 1]$ . The  $U_{\bar{n}}(z)$  are the roots of the sixth-degree polynomial

$$R_z(u) = \prod_{n=1}^6 (u - U_{\bar{n}}(z)).$$

As  $z$  varies in  $\mathbf{CP}^2 - \{F(z) = 0\}$ ,  $R_z(u)$  yields a family of sextic resolvents. Since  $\mathcal{V}_{3.360}$  permutes the  $C_{\bar{n}}(z)^3$  simply—no multiplicative character appears,  $R_z$  is  $\mathcal{V}$ -invariant in  $z$  and hence, so is each  $u$ -coefficient. Expressing the coefficients in terms of the basic invariants  $F(z)$ ,  $\Phi(z)$ ,  $\Psi(z)$  and then converting to  $Y_1$  and  $Y_2$  yields, in  $\text{bub}_{\overline{22}}$  coordinates, the resolvents

$$\begin{aligned} R_Y(u) &= R_{(Y_1, Y_2)}(u) \\ &= u^6 + \frac{-5 + \sqrt{15}i}{90} u^5 + \\ &\quad \frac{11(1 - \sqrt{15}i) - 3(3 + \sqrt{15}i)Y_1}{2^2 3^5 5^2} u^4 + \frac{(100 + 57\sqrt{15}i) + 9(30 + \sqrt{15}i)Y_1}{3^9 5^4} u^3 + \\ &\quad \frac{-(152 + 17\sqrt{15}i) + 18(-21 + 4\sqrt{15}i)Y_1 + 27(-4 + \sqrt{15}i)Y_1^2}{2^2 3^{11} 5^5} u^2 + \\ &\quad \frac{(425 + 103\sqrt{15}i) + 6(75 + 193\sqrt{15}i)Y_1 + 27(-25 + 33\sqrt{15}i)Y_1^2 - 7776\sqrt{15}iY_2}{2^3 3^{14} 5^8} u + \\ &\quad \frac{-(5 + 3\sqrt{15}i) + 9(15 - 7\sqrt{15}i)Y_1 + 81(25 - \sqrt{15}i)Y_1^2 + 81(45 + 11\sqrt{15}i)Y_1^3}{2^4 3^{18} 5^8}. \end{aligned}$$

This makes explicit the fact that the solution of  $R_Y(u)$  follows from inversion of  $Y$ .

For the unbarred functions

$$U_n(z) = \frac{C_n(z)^3}{F(z)},$$

one obtains the associated resolvents  $S_Y$  from  $R_Y$  by complex conjugation of the  $u$ -coefficients:

$$S_Y(u) = \overline{R_Y(\overline{u})}.$$

## 4.2 Special sixth-degree resolvents

For the resolvents  $R_Y$  the parameter space is an affine plane  $[Y_1, Y_2, 1]$  that lifts to  $\mathbf{CP}^2 - \{F(z) = 0\}$ . There is a complementary set of resolvents parametrized by a  $\mathbf{CP}^1$  that lifts to  $\{F(z) = 0\}$ .

The sixth-degree  $\mathcal{V}$ -invariant curve  $\{F(z) = 0\}$  is a genus 10 surface that contains three special Valentiner orbits:

$$\begin{aligned} \mathcal{O}_{72} &= \{F(z) = 0\} \cap \{\Phi(z) = 0\} \\ \mathcal{O}_{90} &= \{F(z) = 0\} \cap \{\Psi(z) = 0\} \\ \mathcal{O}_{180} &= \{F(z) = 0\} \cap \{X(z) = 0\} - \mathcal{O}_{90}. \end{aligned}$$

On  $\{F(z) = 0\}$  the rational map

$$V(z) = \alpha \frac{\Phi(z)^5}{\Psi(z)^2}$$

gives<sup>45</sup> the 2-4-5 quotient of  $\{F(z) = 0\}$  under  $\mathcal{V}$ :

$$V : \{F(z) = 0\} \rightarrow \{F(z) = 0\}/\mathcal{V}.$$

Furthermore, the icosahedral function

$$S_{\overline{1}}(z) = \frac{\Phi(z)^2}{\Psi(z)} C_{\overline{1}}(z)^3$$

divides  $\{F(z) = 0\}$  by the icosahedral subgroup  $\mathcal{I}_{\overline{1}}$ :

$$S_{\overline{1}} : \{F(z) = 0\} \rightarrow \{F(z) = 0\}/\mathcal{I}_{\overline{1}}.$$

A value  $V_0 \neq 0, 1, \infty$  of  $V$  has an inverse image on  $\{F(z) = 0\}$  that consists of a  $\mathcal{V}$ -orbit of size 360 while the image of  $V^{-1}(V_0) \cap \{F(z) = 0\}$  under  $S_{\overline{1}}$  is the set of six points

$$\left\{ S_{\overline{n}}(z) = \frac{\Phi(z)^2 C_{\overline{n}}(z)^3}{\Psi(z)} \mid n = 1, \dots, 6 \right\}$$

where  $z \in V^{-1}(V_0) \cap \{F(z) = 0\}$ . The  $S_{\overline{n}}(z)$  supply roots of a  $\mathcal{V}$ -parametrized family of sextic resolvents

$$T_z(s) = \prod_{n=1}^6 (s - S_{\overline{n}}(z)).$$

As above,  $T_z$  is  $\mathcal{V}$ -invariant in  $z$  and hence, so is each  $s$ -coefficient. Expressing the coefficients in terms of the basic invariants  $F(z)$ ,  $\Phi(z)$ , and  $\Psi(z)$ , restricting to  $\{F(z) = 0\}$ , and converting to  $V$  gives the one-parameter resolvents

$$T_V(s) = s^6 - \frac{-3+\sqrt{15}i}{2^5 3^3 5^2} V s^4 - \frac{4+\sqrt{15}i}{2^8 3^6 5^5} V^2 s^2 + \frac{\sqrt{15}i}{2^6 3^7 5^8} V^2 s + \frac{45-11\sqrt{15}i}{2^{13} 3^{11} 5^8} V^3.$$

Again, with the unbarred functions

$$S_n(z) = \frac{\Phi(z)^2}{\Psi(z)} C_n(z)^3,$$

conjugation of the coefficients of  $T_V$  yields the unbarred resolvents.

### 4.3 Parametrized families of Valentiner groups

The algorithms that solve given resolvents  $R_Y$  or  $T_V$  employ an iteration of a dynamical system  $h_Y(w)$  or  $h_V(w)$  that belongs to a family of maps

- 1) parametrized by  $Y = (Y_1, Y_2)$  or  $V$
- 2) each member of which is conjugate to  $h_{19}(y)$ .

The first task is to parametrize by  $Y$  and  $V$  families of Valentiner groups. Each such group supports a conic-fixing 19-map the computation of which follows that of its conjugate,  $h_{19}(y)$ .

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<sup>45</sup>The constant  $\alpha$  is chosen so that  $V = 1$  at a 180-point. In  $\text{bub}_{\overline{22}}$  coordinates,  $\alpha = 3/8$ .

### 4.3.1 Invariant building-blocks

Success in finding a 19-map for almost every value<sup>46</sup> of  $Y$  or  $V$  requires provision only of basic invariant forms

$$F_Y, \Phi_Y, \Psi_Y, X_Y \text{ or } F_V, \Phi_V, \Psi_V, X_V$$

that are parametrized by  $Y$  and  $V$ . In turn, the latter three of each type depend on the single forms  $F_Y$  and  $F_V$ .

Much of this development amounts to keeping track of coordinates. The Valentiner actions  $\mathcal{V}_z$  and  $\mathcal{V}_y$  on the respective planes  $\mathbf{CP}_z^2$  and  $\mathbf{CP}_y^2$  are the same—the parameter  $z$  merely replaces  $y$ . Think of these as a parameter and reference space respectively.

To obtain a parametrized sixth-degree form in the general case:

- 1) Compose  $F(z)$  with a certain family of maps  $\tau_z(w)$  each of which is  $\mathcal{V}_z$ -equivariant and linear in  $w$ . (The  $w$ -space  $\mathbf{CP}_w^2$  is the iteration space.)
- 2) Express the coefficients of the  $w$  monomials in terms of  $F(z)$ ,  $\Phi(z)$ , and  $\Psi(z)$ .
- 3) Convert these coefficients to expressions in  $Y_1$  and  $Y_2$ —as in the derivation of  $R_Y$ .

The special case requires more care.

- 1) Compose  $F(z)$  with a select family of maps  $\sigma_z(w)$  each of which is  $\mathcal{V}_z$ -equivariant and linear in  $w$ .
- 2) Restrict to  $\{F(z) = 0\}$  and express the coefficients of the  $w$  monomials in terms<sup>47</sup> of  $\Phi(z)$ ,  $\Psi(z)$ .
- 3) Divide through by any overall factors in  $\Phi$  and  $\Psi$  to obtain a polynomial whose degree in  $z$  is a multiple of 60—the degree of  $V$ —and then express the result in terms of  $V$ .

### 4.3.2 Sixth-degree forms in the 2-parameter case.

Consider the family of maps

$$y = \tau_z(w) = [F(z)^4 \cdot z] w_1 + [F(z) \cdot h_{19}(z)] w_2 + k_{25}(z) w_3$$

that are degree 25 in  $z$  and projective transformations in  $w$ . Here,  $k_{25}(z)$  is the equivariant whose expression in “Hermitian<sup>48</sup> coordinates” is

$$k_{25}(z) = \overline{\nabla F(\overline{\nabla F(z)})}.$$

With  $\tau_z$ , one constructs a  $z$ -parametrized family of Valentiner groups  $\mathcal{V}_w = \tau_z^{-1} V_y \tau_z$  each member of which acts on  $\mathbf{CP}_w^2$ . By construction, this family possesses an equivariance property: for  $T \in \mathcal{V}_z, \mathcal{V}_y$ ,

$$\tau_{Tz}(w) = T \tau_z(w).$$

<sup>46</sup>The singular values  $Y(\{X(z) = 0\})$  of  $Y$  and  $0, 1, \infty$  of  $V$  provide exceptions.

<sup>47</sup>The choice of  $\sigma_z(w)$  becomes significant at this stage; see below.

<sup>48</sup>The elements  $T \in \mathcal{V}_{3.360}$  satisfy  $T\overline{T^t} = I$ .

Hence,

$$F(\tau_{Tz}(w)) = F(\tau_z(w))$$

so that the  $w$ -coefficients of  $F(\tau_z(w))$  are  $\mathcal{V}_z$ -invariant and thereby expressible in terms of the basic forms  $F(z)$ ,  $\Phi(z)$ ,  $\Psi(z)$ , and  $X(z)$ . However, since the degree *in*  $z$  of each  $w$ -coefficient is  $6 \cdot 25 = 150$ , an odd power of  $X(z)$  cannot appear in the decomposition of these coefficients into polynomials in the basic invariants. Being  $\mathcal{V}_{3 \cdot 360}$ -invariant,  $X^2$  decomposes<sup>49</sup> into a polynomial in  $F$ ,  $\Phi$ , and  $\Psi$ . Thus, each coefficient is a combination of the forms of degrees 6, 12, and 30. After division by an appropriate power of  $F(z)$  as well as a simplifying numerical factor  $\alpha$ , the result is expressible in terms of  $Y_1$  and  $Y_2$ :

$$F_Y(w) = \frac{F(\tau_z(w))}{\alpha F(z)^{25}}. \quad (4)$$

The coefficients of the  $Y$  monomials in a  $w$ -coefficient are solutions to a system of linear equations. An expression for this fundamental form appears in Appendix B.

An important matter concerns the degeneration of  $\tau_z(w)$  where the determinant  $|\tau_z|$  vanishes. Taking  $z^t$ ,  $h_{19}(z)^t$ , and  $k_{25}(z)^t$  to be column vectors

$$\begin{aligned} |\tau_z| &= \begin{vmatrix} F(z)^4 \cdot z^t & F(z) \cdot h_{19}(z)^t & k_{25}(z)^t \end{vmatrix} \\ &= F(z)^5 \begin{vmatrix} z^t & h_{19}(z)^t & k_{25}(z)^t \end{vmatrix} \\ &= -1458 F(z)^5 X(z). \end{aligned}$$

The final equality follows by uniqueness of  $X$  as a degree 45 invariant and evaluation of  $|\tau_z|$ ,  $F(z)$ , and  $X(z)$  at a single point. Thus, the square of  $|\tau_z|$  is expressible in  $F(z)$ ,  $\Phi(z)$ , and  $\Psi(z)$  alone. In terms of  $Y$ ,

$$\begin{aligned} |\tau_z|^2 &= 1458^2 F^{10} X^2 \\ &= 432 F^{25} (Y_1 + 20 Y_1^2 + 204 Y_1^3 + 1094 Y_1^4 + 3271 Y_1^5 + 3078 Y_1^6 + \\ &\quad 1404 Y_1^7 + 18 Y_2 + 198 Y_1 Y_2 + 954 Y_1^2 Y_2 - 198 Y_1^3 Y_2 - 5508 Y_1^4 Y_2 + \\ &\quad 1944 Y_1^5 Y_2 - 648 Y_2^2 - 7776 Y_1 Y_2^2 - 5832 Y_1^2 Y_2^2 + 11664 Y_2^3). \end{aligned} \quad (5)$$

#### 4.3.3 The special case.

Take the family of maps

$$\sigma_z(w) = [72 \Phi(z)^4 \cdot z] w_1 + [\Psi(z) \cdot h_{19}(z)] w_2 + [24 \Phi(z)^2 \cdot k_{25}(z)] w_3.$$

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<sup>49</sup>See (1).

with  $z$ -degree 49 and  $w$ -degree one. The integer coefficients have been chosen so that, in  $\text{bub}_{22}$ -coordinates, the point

$$[w_1, w_2, w_3] = [1, 1, 1]$$

corresponds to the map

$$\nabla F(z) \times \nabla X(z)$$

associated with the 2-form  $dF \wedge dX$ . As in the general situation,

$$\sigma_{Tz}(w) = T\sigma_z(w)$$

so that

$$F(\sigma_{Tz}(w)) = F(\sigma_z(w)).$$

Thus, the  $w$ -coefficients are  $\mathcal{V}_z$ -invariant and thereby expressible in terms of the basic forms  $F(z)$ ,  $\Phi(z)$ ,  $\Psi(z)$ , and  $X(z)$ . Since the degree in  $z$  of  $F(\sigma_z(w))$  is  $6 \cdot 49 = 294$ , odd powers of  $X(z)$  cannot take part in the basic invariant decomposition of these coefficients. Furthermore, restriction of the parameter space  $\mathbf{CP}_z^2$  to  $\{F(z) = 0\}$  yields coefficients in  $\Phi(z)$  and  $\Psi(z)$  alone. Finally, since  $294 = 22 \cdot 12 + 30 = 2 \cdot 12 + 9 \cdot 30$ ,  $F(\sigma_z(w))|_{\{F(z)=0\}}$  is divisible by  $\Phi(z)^2 \Psi(z)$ . Hence,

$$F(\sigma_z(w))|_{\{F(z)=0\}} = \eta \Phi(z)^2 \Psi(z)^9 F_V(w)$$

where  $\eta$  is a simplifying numerical factor and  $F_V(w)$  is a polynomial that is degree four in  $V$  and degree six in  $w$ . The expression for  $F_V(w)$  also appears in Appendix B.

Since the parameter space gets restricted to  $\{F(z) = 0\}$ ,  $|\sigma_z|$  should not vanish there. In fact, 49 is the lowest degree in which this fails to occur for three (projectively) distinct maps. Explicitly,

$$\begin{aligned} |\sigma_z| &= \left| \begin{array}{c|c|c} 72 \Phi^4(z) \cdot z & \Psi(z) \cdot h_{19}(z) & 24 \Phi^2(z) \cdot k_{25}(z) \end{array} \right| \\ &= 24 \cdot 72 \Phi^4(z) \Psi(z) \Phi^2(z) \left| \begin{array}{c|c|c} z & h_{19}(z) & k_{25}(z) \end{array} \right| \\ &= 2^7 3^9 \Phi(z)^6 \Psi(z) X(z). \end{aligned}$$

Furthermore, on  $\{F(z)=0\}$ , the expression (1) for  $X^2$  reduces to

$$\begin{aligned} X^2 &= -\frac{1}{81} (8 \Phi^5 \Psi - 3 \Psi^3) \\ &= -\frac{\Psi^3}{81} (8 \frac{\Phi^5}{\Psi^2} - 3) \\ &= -\frac{\Psi^3}{81} (8 \frac{3V}{8} - 3) \end{aligned}$$



$$= -\frac{\Psi^3}{27}(V-1).$$

Consequently,

$$\begin{aligned} |\sigma_z|^2 &= -2^{14}3^{15}\Phi^{12}\Psi^5(V-1) \\ &= -2^{14}3^{15}\Phi^2\Psi^9\frac{\Phi^{10}}{\Psi^4}(V-1) \\ &= -2^{14}3^{15}\Phi^2\Psi^9\left(\frac{3V}{8}\right)^2(V-1) \\ &= -2^83^{17}\Phi^2\Psi^9V^2(V-1). \end{aligned}$$

#### 4.3.4 The remaining basic invariants.

The forms of degrees 12, 30, and 45 arise from the sixth-degree invariant as before. However, a parametrized change of coordinates requires special handling. Under  $y = Ax$  the Hessian  $H_x(F(x))$ , Bordered Hessian  $BH_x(F(x), G(x))$ , and Jacobian  $J_x(F(x), G(x), K(x))$  transform as

$$\begin{aligned} H_x(F(y)) &= A^t H_y(F(y)) A \\ BH_x(F(y), G(y)) &= \left( \begin{array}{c|c} A^t & 0 \\ \hline 0 & 1 \end{array} \right) BH_y(F(y), G(y)) \left( \begin{array}{c|c} A & 0 \\ \hline 0 & 1 \end{array} \right) \\ J_x(F(y), G(y), K(y)) &= J_y(F(y), G(y), K(y)) A \end{aligned}$$

where the subscript indicates the differentiation variable. As for transformation of the respective determinants:

$$\begin{aligned} |H_x(F(y))| &= |A|^2 |H_y(F(y))| \\ |BH_x(F(y), G(y))| &= |A|^2 |BH_y(F(y), G(y))| \\ |J_x(F(y), G(y), K(y))| &= |A| |J_y(F(y), G(y), K(y))|. \end{aligned}$$

For the parametrized change of coordinates  $y = \tau_z(w)$ , let<sup>50</sup>

$$\begin{aligned} \Phi_Y(w) &= \alpha_\Phi |H_w(F_Y(w))| \\ \Psi_Y(w) &= \alpha_\Psi |BH_w(F_Y(w), \Phi_Y(w))| \\ X_Y(w) &= \alpha_X |J_w(F_Y(w), \Phi_Y(w), \Psi_Y(w))|. \end{aligned}$$

Then

$$\begin{aligned} \Phi(y) &= \alpha_\Phi |H_y(F(y))| \\ &= \alpha_\Phi |H_y(F(\tau_z(w)))| \end{aligned}$$

---

<sup>50</sup>Recall the constants  $\alpha_\Phi, \alpha_\Psi, \alpha_X$  from Section 2.4.2.

$$\begin{aligned}
&= \alpha_\Phi |\tau_z|^{-2} |H_w(\alpha F(z)^{25} F_Y(w))| \\
&= \frac{\alpha_\Phi}{|\tau_z|^2} (\alpha F(z)^{25})^3 |H_w(F_Y(w))| \\
&= \frac{(\alpha F(z)^{25})^3}{|\tau_z|^2} \Phi_Y(w)
\end{aligned} \tag{6}$$

$$\begin{aligned}
\Psi(y) &= \alpha_\Psi |BH_y(F(y), \Phi(y))| \\
&= \alpha_\Psi |BH_y(F(\tau_z(w)), \Phi(\tau_z(w)))| \\
&= \alpha_\Psi |\tau_z|^{-2} |BH_w(\alpha F(z)^{25} F_Y(w), \frac{(\alpha F(z)^{25})^3}{|\tau_z|^2} \Phi_Y(w))| \\
&= \frac{(\alpha F(z)^{25})^8}{|\tau_z|^6} \Psi_Y(w)
\end{aligned} \tag{7}$$

$$\begin{aligned}
X(y) &= \alpha_X |J_y(F(y), \Phi(y), \Psi(y))| \\
&= \frac{(\alpha F(z)^{25})^{12}}{|\tau_z|^9} X_Y(w).
\end{aligned} \tag{8}$$

With  $y = \sigma_z(w)$ , similar calculations lead to the one-parameter forms:

$$\begin{aligned}
\Phi(y) &= \frac{(\eta \Phi(z)^2 \Psi(z)^9)^3}{|\sigma_z|^2} \Phi_V(w) \\
\Psi(y) &= \frac{(\eta \Phi(z)^2 \Psi(z)^9)^8}{|\sigma_z|^6} \Psi_V(w) \\
X(y) &= \frac{(\eta \Phi(z)^2 \Psi(z)^9)^{12}}{|\sigma_z|^9} X_V(w).
\end{aligned}$$

#### 4.3.5 The 19-maps

With an invariant system in place for each (non-singular) value of  $Y$  (or  $V$ ), production of the degree 19 map that preserves all 12 of the conics proceeds as before:

- 1) determine a 64-map  $(f_{64})_Y(w)$  that vanishes at  $\{X_Y(w) = 0\}$
- 2) compute  $(f_{19})_Y(w) = (f_{64})_Y(w)/X_Y(w)$
- 3) compute the conic-fixing map  $h_Y(w)$ .

In fact, the previous calculations in the coordinates  $\{y_1, y_2, y_3\}$  provide a framework for those at hand. Into the  $y$ -expressions involving  $F(y)$ ,  $\Phi(y)$ ,  $\Psi(y)$ , and  $X(y)$  as well as the maps

$\psi_{16}(y)$ ,  $\phi_{34}(y)$ , and  $f_{40}(y)$  substitute the appropriate ones in  $Y$  and  $w$ , namely,  $F_Y(w)$ ,  $\Phi_Y(w)$ ,  $\Psi_Y(w)$ ,  $X_Y(w)$ ,  $\psi_Y(w)$ ,  $\phi_Y(w)$ , and  $f_Y(w)$ . The substitutions for the invariants appear above.<sup>51</sup> Concerning maps, they transform as the coefficients of 2-forms. In terms of the cross-product, for  $u = Ax$ ,

$$\begin{aligned}\nabla_x(F(u)) \times \nabla_x(G(u)) &= A^t \nabla_u F(u) \times A^t \nabla_u G(u) \\ &= |A^t| ((A^t)^{-1})^t [\nabla_u F(u) \times \nabla_u G(u)] \\ &= |A| A^{-1} (\nabla_u F(u) \times \nabla_u G(u)).\end{aligned}$$

Accordingly,

$$\begin{aligned}\psi(y) &= \nabla_y F(y) \times \nabla_y \Phi(y) \\ &= \frac{1}{|\tau_z|} \tau_z \left[ \nabla_w \left( \alpha F(z)^{25} F_Y(w) \right) \times \nabla_w \left( \frac{(\alpha F(z)^{25})^3}{|\tau_z|^2} \Phi_Y(w) \right) \right] \\ &= \frac{(\alpha F(z)^{25})^4}{|\tau_z|^3} \tau_z \nabla_w F_Y(w) \times \nabla_w \Phi_Y(w) \\ &= \frac{(\alpha F(z)^{25})^4}{|\tau_z|^3} \tau_z (\psi_Y(w)) \\ \\ \phi(y) &= \nabla_y F(y) \times \nabla_y \Psi(y) \\ &= \frac{1}{|\tau_z|} \tau_z \left[ \nabla_w \left( \alpha F(z)^{25} F_Y(w) \right) \times \nabla_w \left( \frac{(\alpha F(z)^{25})^8}{|\tau_z|^6} \Psi_Y(w) \right) \right] \\ &= \frac{(\alpha F(z)^{25})^9}{|\tau_z|^7} \tau_z (\phi_Y(w)) \\ \\ f(y) &= \nabla_y \Phi(y) \times \nabla_y \Psi(y) \\ &= \frac{1}{|\tau_z|} \tau_z \left[ \nabla_w \left( \frac{(\alpha F(z)^{25})^3}{|\tau_z|^2} \Phi_Y(w) \right) \times \nabla_w \left( \frac{(\alpha F(z)^{25})^8}{|\tau_z|^6} \Psi_Y(w) \right) \right] \\ &= \frac{(\alpha F(z)^{25})^{11}}{|\tau_z|^9} \tau_z (f_Y(w)).\end{aligned}$$

The one-parameter maps transform as follows:

$$\begin{aligned}\psi(y) &= \frac{(\eta \Phi(z)^2 \Psi(z)^9)^4}{|\sigma_z|^3} \sigma_z (\psi_Y(w)) \\ \\ \phi(y) &= \frac{(\eta \Phi(z)^2 \Psi(z)^9)^9}{|\sigma_z|^7} \sigma_z (\phi_Y(w))\end{aligned}$$

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<sup>51</sup>See (4), (6), (7), and (8).

$$f(y) = \frac{(\eta \Phi(z)^2 \Psi(z)^9)^{11}}{|\sigma_z|^9} \sigma_z(f_Y(w)).$$

Making the substitutions into the expression for  $f_{64}(y)$  yields a collection of maps  $(f_{64})_Y(w)$  each of which is divisible by  $X_Y(w)$ . Finally, substitution into the formula for the canonical map  $h_{19}(y)$  supplies the conic-fixing family  $h_Y(w)$ . The explicit calculations appear in Appendix B.

#### 4.3.6 General expressions for the canonical 19-map.

Each  $w$ -coefficient in

$$h_z(w) = \tau_z^{-1}(h_{19}(\tau_z(w)))$$

is a polynomial in  $Y_1$  and  $Y_2$  with  $z$ -degree

$$\deg_z h_z(w) \leq 21 \cdot 25 = 19 \cdot 25 + 2 \cdot 25 = 12 \cdot 40 + 45 = 30 \cdot 16 + 45.$$

The factor of  $|\tau_z|^{-1}$  due to  $\tau_z^{-1}$  does not affect the map on  $\mathbf{CP}^2$  and so, is neglected. After dividing away a factor of  $X(z)$  the polynomial map  $h_Y(w)$  satisfies

$$\deg_{Y_1} h_Y(w) \leq 40 \quad \deg_{Y_2} h_Y(w) \leq 16.$$

Hence, one finds the coefficients of the  $w$  monomials by solving, for each term, 353 linear equations.

In the 1-parameter case,

$$h_z(w) = \sigma_z^{-1}(h_{19}(\sigma_z(w)))$$

and

$$\deg_z h_z(w) \leq 21 \cdot 49 = 12 \cdot 82 + 45 = 12 \cdot 2 + 30 \cdot 32 + 45.$$

Thus, on  $\{F(z) = 0\}$ ,  $h_z(w)$  is divisible by  $\Phi(z)^2 X(z)$  so that

$$\deg_z \left( \frac{h_z(w)}{\Phi(z)^2 X(z)} \right) = 60 \cdot 16 \quad \deg_V h_V(w) \leq 16.$$

### 4.4 Symmetry lost, a root found

Under Conjecture 1, the trajectory  $\{h_{19}^k(y_0)\}$  converges to a pair of 72-points for almost any  $y_0 \in \mathbf{CP}_y^2$ . Being conjugate to  $h_{19}(y)$ , the maps  $h_Y(w)$  share this property for points in  $\mathbf{CP}_w^2$ . Breaking the  $\mathcal{A}_6$  symmetry of  $R_Y(u) = 0$  qualifies  $h_Y(w)$  for a role in root-finding.<sup>52</sup>

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<sup>52</sup>An analogous treatment applies in the one-parameter case.

#### 4.4.1 Root-selection

Consider the rational function

$$\overline{J}_z(w) = \frac{\overline{\Gamma}_z(w)^3}{F(z) \Psi(\tau_z(w))}$$

where

$$\overline{\Gamma}_z(w) = \sum_{m=1}^6 \prod_{n \neq m} C_{\overline{n}}(\tau_z(w)) \cdot C_{\overline{m}}(z).$$

At a 72-point pair  $\{q_1, q_2\} = \tau_z^{-1}(\{p_{\overline{a}b_1}, p_{\overline{a}b_2}\})$  in  $w$  space, five of the six terms in  $\overline{\Gamma}_z(w)$  vanish. The result is a “selection” of one of the six roots  $U_{\overline{a}}(z)$  of  $R_Y(u)$ :

$$\begin{aligned} \overline{J}_z(\{q_1, q_2\}) &= \frac{[\prod_{n \neq a} C_{\overline{n}}(\{p_{\overline{a}b_1}, p_{\overline{a}b_2}\})]^3}{\Psi(\{p_{\overline{a}b_1}, p_{\overline{a}b_2}\})} \frac{C_{\overline{a}}(z)^3}{F(z)} \\ &= \frac{U_{\overline{a}}(z)}{\mu}. \end{aligned}$$

Here,  $\mu$  is the value of

$$\frac{\Psi(\{p_{\overline{a}b_1}, p_{\overline{a}b_2}\})}{[\prod_{n \neq a} C_{\overline{n}}(\{p_{\overline{a}b_1}, p_{\overline{a}b_2}\})]^3}$$

which is constant on  $\mathcal{O}_{72}$ :

$$\mu = \frac{6561 (279 + 145\sqrt{15}i)}{2}.$$

In light of the  $\mathcal{V}_z$ -invariance of  $\overline{\Gamma}_z(w)$ ,  $\overline{J}_z$  also enjoys this property and so, presents a second face which expresses itself in  $Y$  and  $w$ . Since

$$\deg_z \overline{\Gamma}_z(w) = 252 = 10 \cdot 25 + 2 = 6 \cdot 42,$$

it transforms by

$$\overline{\Gamma}_Y(w) = \frac{\overline{\Gamma}_z(w)}{\beta F(z)^{42}}$$

where  $\beta$  is a simplifying factor for the coefficients over  $w$ . Each coefficient is a polynomial in  $Y_1$  and  $Y_2$  whose coefficients satisfy a system of linear equations.<sup>53</sup> Now, let  $T_Y$  be the polynomial<sup>54</sup> in  $Y_1$  and  $Y_2$  that satisfies

$$|\tau_z|^2 = F(z)^{25} T_Y.$$

Then, from (7),

$$\begin{aligned} \Psi(\tau_z(w)) &= \frac{(\alpha F(z)^{25})^8}{(F(z)^{25} T_Y)^3} \Psi_Y(w) \\ &= \frac{\alpha^8 F(z)^{125}}{T_Y^3} \Psi_Y(w). \end{aligned}$$

---

<sup>53</sup>The polynomial  $\overline{\Gamma}_Y(w)$  is rather large and takes an explicit albeit partial form in Appendix B.

<sup>54</sup>See (5).

Finally,

$$\overline{J}_Y(w) = \frac{\beta^3}{\alpha^8} \frac{(T_Y \overline{\Gamma}_Y(w))^3}{\Psi_Y(w)}.$$

In the one-parameter case, the above development goes through with  $\overline{J}_z(w)$  and  $\overline{\Gamma}_z(w)$  replaced by

$$\overline{K}_z(w) = \frac{\Phi(z)^2 \Theta_z(w)^3}{\Psi(z) \Psi(\tau_z(w))}$$

and

$$\overline{\Theta}_z(w) = \sum_{m=1}^6 \prod_{n \neq m} C_{\overline{n}}(\sigma_z(w)) \cdot C_{\overline{m}}(z).$$

First of all, there is the transformation<sup>55</sup> of  $\overline{\Theta}_z(w)$  under  $\sigma_z$  on  $\{F(z) = 0\}$ :

$$\overline{\Theta}_V(w) = \frac{\overline{\Theta}_z(w)|_{\{F(z)=0\}}}{\gamma \Phi(z) \Psi(z)^{16}}.$$

Furthermore, restricting to  $\{F(z) = 0\}$  yields, on the one hand, a root of  $T_V(s)$

$$\overline{K}_z(\{q_1, q_2\}) = \frac{S_{\overline{a}}(z)}{\mu}$$

and, on the other,

$$\begin{aligned} \overline{K}_V(w) &= \frac{\Phi(z)^2 [\gamma \Phi(z) \Psi(z)^{16} \overline{\Theta}_V(w)]^3 |\sigma_z|^6}{(\eta \Phi(z)^2 \Psi(z)^9)^8 \Psi_V(w)} \\ &= -\frac{2^9 3^{50} \gamma^3 V^5 (V-1)^3 \overline{\Theta}_V(w)}{\eta^8 \Psi_V(w)}. \end{aligned}$$

#### 4.4.2 The algorithm

Now within reach are the ingredients required for preparation of a root-finding algorithm.<sup>56</sup> To summarize the procedure:

- 1) Select a value  $A = (A_1, A_2)$  of  $Y = (Y_1, Y_2)$  and, thereby, a sixth-degree resolvent  $R_A(u)$ . (For sake of description, let  $z \in Y_1^{-1}(A_1) \cap Y_2^{-1}(A_2)$ . The algorithm actually finds a root without explicitly inverting  $Y_1$  or  $Y_2$ .)
- 2) From an initial point  $w_0 \in \mathbf{CP}_w^2$ , iterate the map  $h_A(w)$  to convergence:

$$h_A^n(w_0) \longrightarrow \{q_1, q_2\} \in (\mathcal{O}_{72})_w \subset \mathbf{CP}_w^2.$$

As output take the pair of approximate 72-points in  $\mathbf{CP}_w^2$

$$\{p_1, p_2\} \approx \{q_1, q_2\} = \{\tau_z^{-1}(p_{\overline{a}b_1}), \tau_z^{-1}(p_{\overline{a}b_2})\}.$$

---

<sup>55</sup>In Appendix B there appears a partial expression for  $\overline{\Theta}_V(w)$ .

<sup>56</sup>Available at <http://math.ucsd.edu/~scrass> are *Mathematica* notebooks and supporting files that implement this method.

3) Using either  $p_1$  or  $p_2$  approximate a root of  $R_A(u)$ :

$$\begin{aligned}
U_{\bar{a}}(z) &\approx \mu \bar{J}_A(p_1) \\
&= \frac{\mu \beta^3 (T_A \bar{\Gamma}_A(p_1))^3}{\alpha^8 \Psi_A(p_1)} \\
&= 2^{79} 3^{94} (11 + 3\sqrt{15}i) \frac{T_A \bar{\Gamma}_A(p_1)^3}{\Psi_A(p_1)}.
\end{aligned}$$

Performing the corresponding steps in the special case produces an approximate solution to the  $V$ -resolvent  $T_{V_0}(s)$  given an initial choice  $V_0$  of  $V$ :

$$\begin{aligned}
S_{\bar{a}}(z) &\approx \mu \bar{K}_{V_0}(p_1) \\
&= -\frac{2^9 3^{50} \mu \gamma^3 V_0^5 (V_0 - 1)^3 \bar{\Theta}_{V_0}(p_1)}{\eta^8 \Psi_{V_0}(p_1)} \\
&= -2^5 5^{-96} (11 + 3\sqrt{15}i) \frac{V_0^5 (V_0 - 1)^3 \bar{\Theta}_{V_0}(p_1)}{\Psi_{V_0}(p_1)}.
\end{aligned}$$

#### 4.5 Getting *all* the roots

Because a pair of 72-points lies on *one* barred conic, the algorithm above determines just one of the six roots of the selected resolvent. This manifests the iteration's failure to break the galois symmetry completely; the stabilizer of a 72-point is a  $\mathcal{D}_5$ . From the coefficient field of the resolvent  $R_A(u)$ , the algorithm leads to an extended field  $K$  whose galois group  $G(\Sigma/K)$  is  $\mathcal{D}_5$ . Of course,  $\Sigma$  is the splitting field for  $R_A$ .

Finding all six roots calls for a dynamical system that converges to 6-cycles in a 360 point orbit and a root-selector function that gives equations in all the roots. Needless to say, this would likely complicate the associated formulas.

## A Seeing is Believing

The following gallery of pictures provides empirical dynamical information for some of the special maps discussed in the text. The program *Dynamics*, running on a Silicon Graphics Indigo-2, created the basin plots.<sup>57</sup>

Complete basin images are the product of the “BAS” routine. In the plots for  $h_{19}$ , this procedure colors a grid-cell in the event that the trajectory of the cell’s center gets close enough to a 72-point to guarantee ultimate attraction to the associated period-2 cycle. The color depends upon the destination of its center. If, in a specified number of iterations, the center’s trajectory fails to converge to a pair of 72-points, the cell’s color is black. The presence of such points would compromise Conjecture 1. In the case of  $\psi_{16}$  “restricted” to a 45-line, the “test” attractor consists of the 45-points. Partial plots come by way of the “BA” algorithm’s coloring of whole “trajectories” of cells, thereby manifesting some aspects of the dynamics.<sup>58</sup> All plots have the maximum resolution available: a  $720 \times 720$  grid of cells. *Mathematica* produced the sketches of the various curves that appear.

In the basin plots for  $h_{19}$  restricted to an affine plane in one of the bub- $\mathbf{RP}^2$ s, the chosen coordinates make evident the map’s  $D_5$  symmetry. Specifically selected here is  $\mathcal{R}_{\overline{22}}$  with the 1-point orbit  $p_{\overline{22}}$  at the origin and the 1-line orbit  $\mathcal{L}_{\overline{22}}$  at infinity. Distributed along the unit circle is the 10-point orbit of 72-points  $p_{\overline{aa}1,2}$  with  $a = 1, 3, 4, 5, 6$ . The five lines of reflective symmetry passing through (0,0) are affine lines in the five  $\mathbf{RP}^1$  intersections with  $\mathcal{R}_{\overline{22}}$  of both the 45-lines  $\mathcal{L}_{\overline{2a2a}}$  and the  $\mathcal{R}_{\overline{aa}}$ ,  $a = 1, 3, 4, 5, 6$ .

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<sup>57</sup>The software is the work of H. Nusse and J. Yorke while E. Kostelich is responsible for the Unix implementation. See their manual [Nusse and Yorke 1994].

<sup>58</sup>See [Nusse and Yorke 1994, pp. 269ff] for a more thorough description.



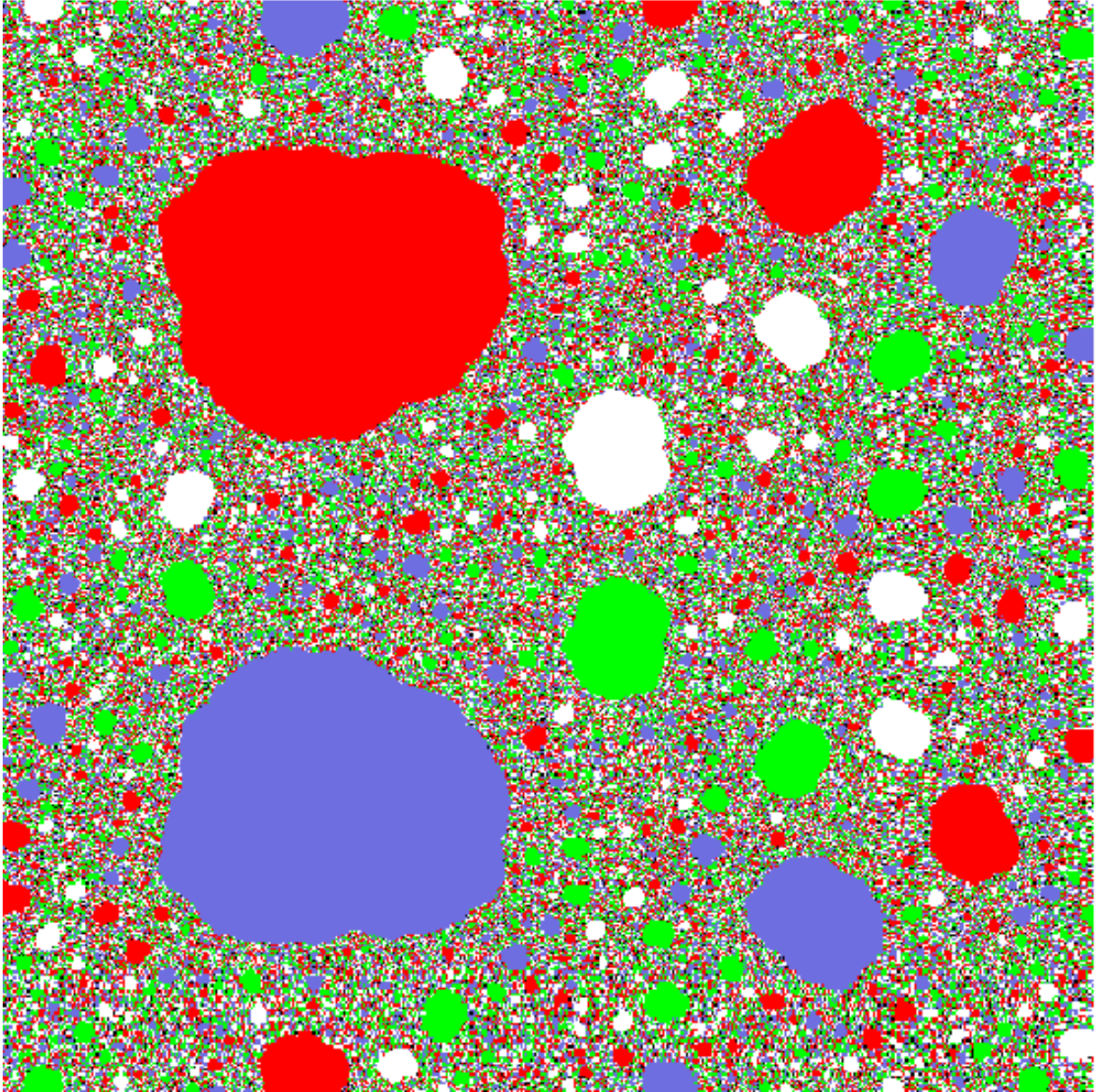


Figure 10: Dynamics of the 16-map

When “restricted” to a 45-line, the degree 16 map  $\psi_{16}$  “mostly” converges to one of the 45-points on the line. Does this occur for almost every point on the line? Do the black specks contain sets of positive measure whose forward orbits fail to converge to one of the four attractive 45-points that lie in the large “central” basins? The BAS algorithm checked 60 iterates before concluding that a trajectory did not converge.

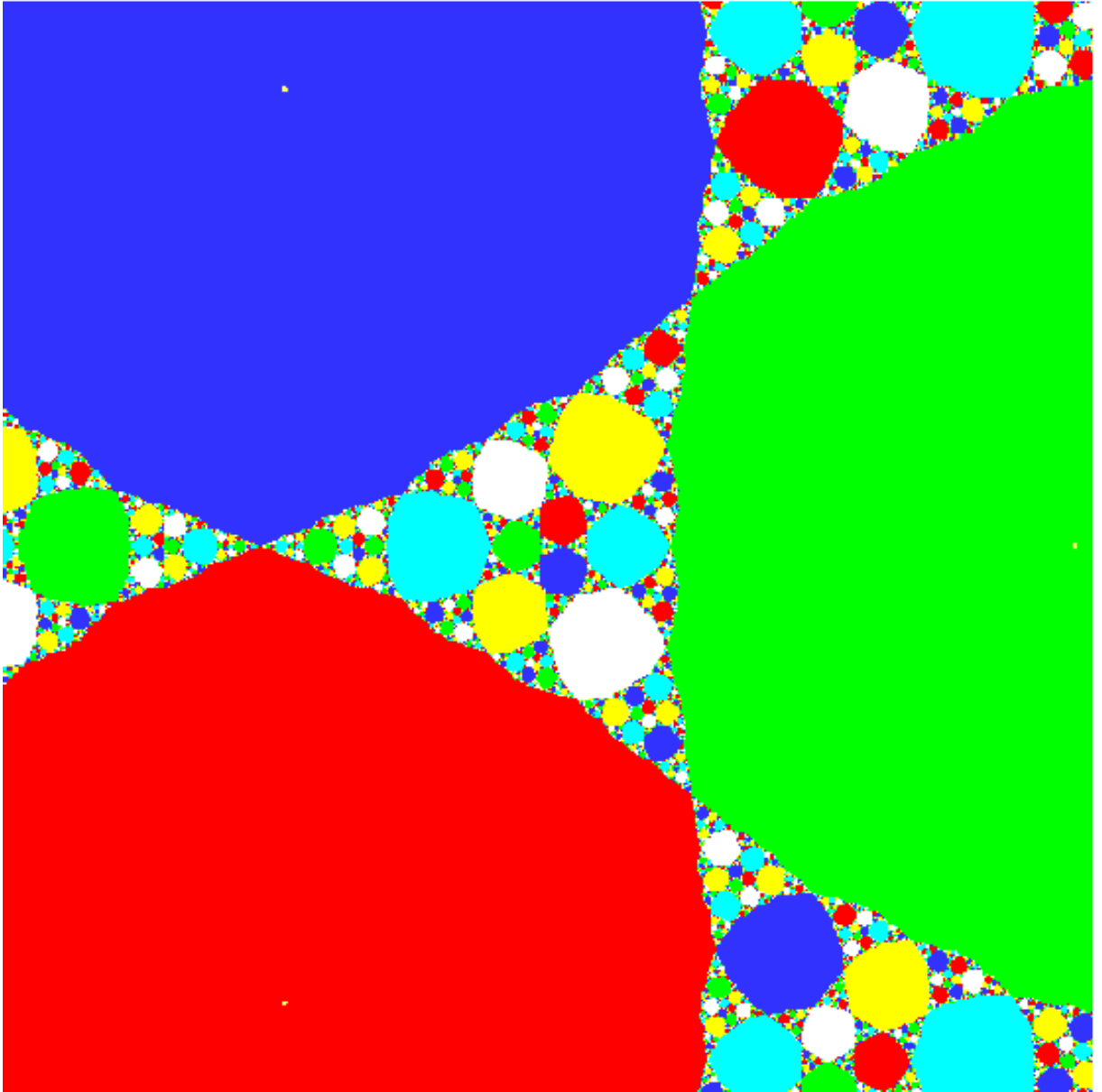


Figure 11: Icosahedral dynamics

The degree 19 map with icosahedral symmetry attracts almost all points in  $\mathbf{C} \cup \{\infty\}$  to an antipodal pair of vertices. Each of the six colors corresponds to such a pair and the three large basins each contain a vertex. For the conic-fixing  $h_{19}$ , the basin plot on each conic is conjugate to this one. Moreover, each basin is the 1-dimensional intersection of a 2-dimensional basin in  $\mathbf{CP}^2$ . Does the backward orbit of these basins fill out  $\mathbf{CP}^2$  in measure?





Figure 12:  $\mathbf{RP}^2$  dynamics of  $h_{19}$

For this plot the vertical and horizontal scales are roughly from -2 to 2. The large “radial” basins are immediate, i.e., each contains one of the 72-points and come in pairs as do the period-2 attractors. Notice the repelling behavior along the 45-lines and particularly at their intersection in the 36-point  $p_{\overline{2}2}$ .

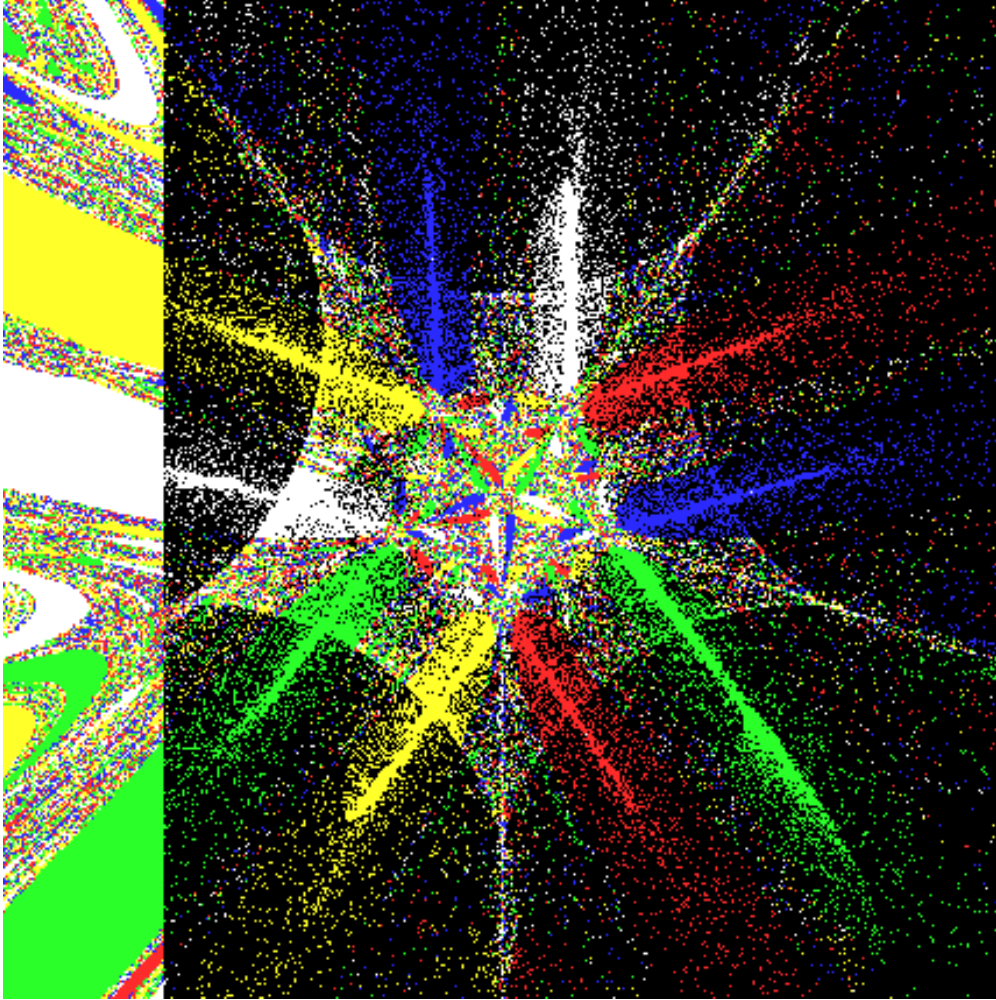


Figure 13:  $\mathbf{RP}^2$  dynamics  $h_{19}$

Again, the vertical and horizontal scales are roughly from -2 to 2. Shown here are trajectories, colored according to their destinations, of the points in the vertical strip on the left. Once more, only the five period-2 cycles of 72-points appear as attractors. Many of the points in the strip map inside the “hazy pentagon” whose vertices lie on the 45-lines—the inner “star” is nearly filled. “Circumscribing” this pentagon is the outer star-like piece of the critical set shown in Figure 16. Furthermore, the pentagon seems to be the image of the inner pentagonal oval. Accordingly, the map folds the plane along the pentagon’s edges just outside of which the 72-points make their presence seen in the dense streaks at  $p_{\bar{a}a_{1,2}}$ . Compare this pattern of streaks to that of the 72-lines given in Figure 14. Since  $\mathcal{C}_{\bar{a}}$  and  $\mathcal{C}_a$  are tangent to  $\mathcal{R}_{\bar{a}a}$  at  $p_{\bar{a}a_{1,2}}$ , the icosahedral 19-map opens up a triangular angle of  $\frac{\pi}{3}$  to  $\frac{4\pi}{3}$ . Thus, the behavior at a 72-point consists of “fourth-powering”. Figure 15 displays this local “squeezing”.

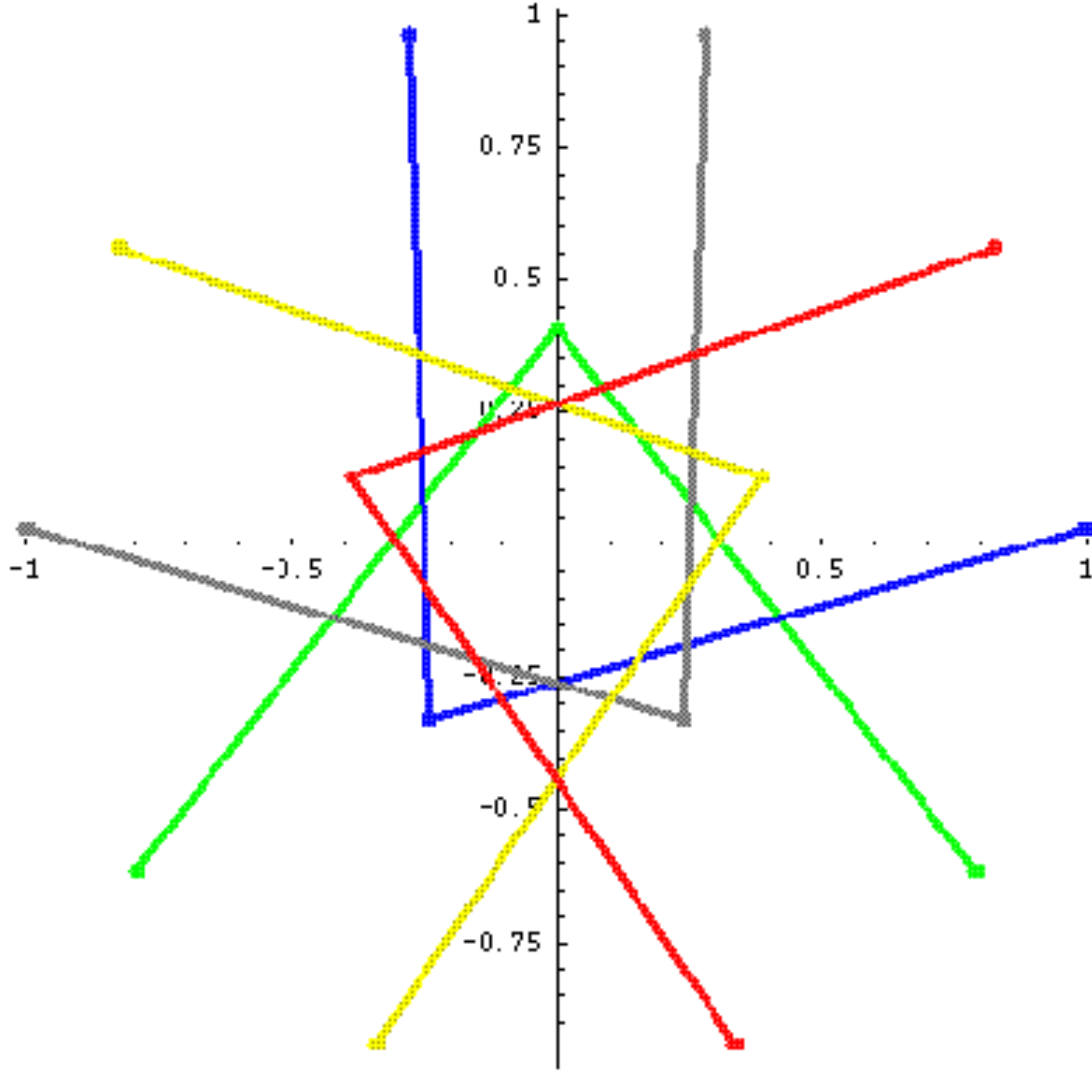


Figure 14: Configuration of 72-lines on a bub- $\mathbf{RP}^2$

Under  $\text{bub}_{\overline{22}}$ , five pairs of 72-lines

$$\{\{\mathcal{L}_{\overline{a}a_1}, \mathcal{L}_{\overline{a}a_2}\} | a = 1, \dots, 6\}$$

map to themselves. Accordingly, each line meets  $\mathcal{R}_{\overline{22}}$  in an  $\mathbf{RP}^1$ . The picture shows their configuration in the affine plane of Figures 12 and 13. Each pair receives a single color according to the scheme of the basin plots. A given pair  $\mathcal{L}_{\overline{a}a_1}$  and  $\mathcal{L}_{\overline{a}a_2}$  passes through the 72-points  $p_{\overline{a}a_2}$  and  $p_{\overline{a}a_1}$  respectively; they intersect in the corresponding repelling and fixed 36-point  $p_{\overline{a}a}$ .

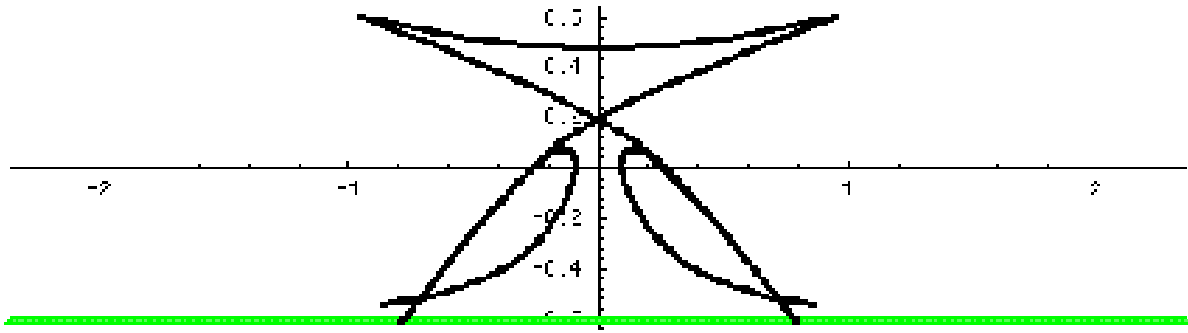


Figure 15: Image on a  $\text{bub-}\mathbf{RP}^2$  of a 36-line under  $h_{19}$

The green horizontal line corresponds to the  $\mathbf{RP}^1$  intersection of  $\mathcal{R}_{\overline{2}2}$  and the 36-line associated with the pair of green 72-points from the basin plots. The dark curve is the image of the line under  $h_{19}$ . Sitting at the sharp cusps are the 72-points which the map exchanges. As indicated in the caption to Figure 13, the line folds over at these critical points. The upper two sharp turns are not critical values; they occur where the line passes through the yellow and red “streak” that approximate 72-lines.

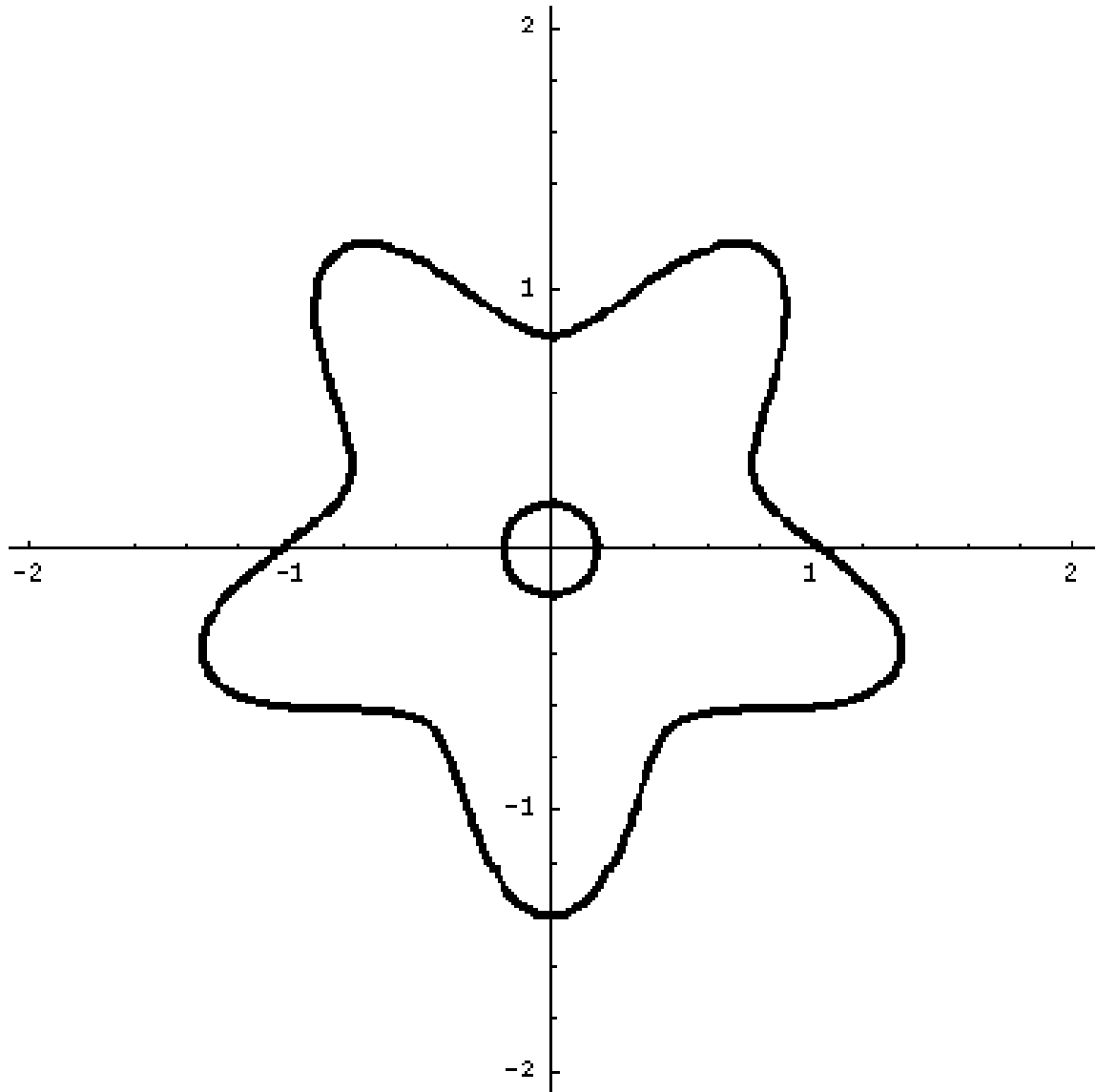


Figure 16: Critical set of  $h_{19}$  on a bub- $\mathbf{RP}^2$

Here is a *Mathematica* contour plot of the sixth degree curve  $\{F = 0\}$  on  $\mathcal{R}_{\overline{22}}$ . At the 10 inflection points are the superattracting 72-points. Further computation suggests that  $\{G_{48} = 0\}$  hits the  $\mathbf{RP}^2$  in a discrete set. Might this set consist only of the singular 72-points?

## B Parametrized Expressions

### B.1 Sixth-degree forms

After solving a system of linear equations for each coefficient of the degree 6 monomials in  $w$ , the simplifying change of coordinates

$$\begin{aligned} w_1 &\rightarrow 4(2w_1 - 23w_2 + 200w_3) \\ w_2 &\rightarrow 2(w_1 - 52w_2 + 64w_3) \\ w_3 &\rightarrow 6w_3 \end{aligned}$$

leaves the following expression for the 2-parameter family of sixth-degree Valentiner invariants:

$$\begin{aligned} F_Y(w) = (2^{10} \cdot 3^{12}) \frac{F(\tau_z(w))}{F(z)^{25}} = \\ 11664(-w_1^2 + 2w_1w_2 - w_2^2 - 14w_1w_3 + 14w_2w_3 - 39w_3^2)(-w_1^2 + 2w_1w_2 - w_2^2 - 4w_1w_3 + \\ 4w_2w_3 + 11w_3^2)(w_1^2 - 2w_1w_2 + w_2^2 + 12w_1w_3 - 12w_2w_3 + 45w_3^2) + \\ 7776w_3(-127w_1^5 + 650w_1^4w_2 - 1285w_1^3w_2^2 + 1225w_1^2w_2^3 - 560w_1w_2^4 + 97w_2^5 - 3790w_1^4w_3 + \\ 15760w_1^3w_2w_3 - 23820w_1^2w_2^2w_3 + 15520w_1w_2^3w_3 - 3670w_2^4w_3 - 44385w_1^3w_3^2 + \\ 140715w_1^2w_2w_3^2 - 144360w_1w_2^2w_3^2 + 48030w_2^3w_3^2 - 260580w_1^2w_3^3 + 558960w_1w_2w_3^3 - \\ 290820w_2^2w_3^3 - 756816w_1w_3^4 + 822741w_2w_3^4 - 831054w_3^5)Y_1 + \\ 54(-1211w_1^6 + 10374w_1^5w_2 - 31275w_1^4w_2^2 + 42880w_1^3w_2^3 - 26160w_1^2w_2^4 + 4176w_1w_2^5 + \\ 1216w_2^6 - 79538w_1^5w_3 + 490810w_1^4w_2w_3 - 1044200w_1^3w_2^2w_3 + 939920w_1^2w_2^3w_3 - \\ 307120w_1w_2^4w_3 + 128w_2^5w_3 - 1619325w_1^4w_3^2 + 8179200w_1^3w_2w_3^2 - 13473000w_1^2w_2^2w_3^2 + \\ 8657280w_1w_2^3w_3^2 - 1728360w_2^4w_3^2 - 14229000w_1^3w_3^3 + 58644000w_1^2w_2w_3^3 - \\ 68856480w_1w_2^2w_3^3 + 23945760w_2^3w_3^3 - 64072440w_1^2w_3^4 + 193407120w_1w_2w_3^4 - \\ 119268000w_2^2w_3^4 - 116140176w_1w_3^5 + 210347712w_2w_3^5 + 49729896w_3^6)Y_1^2 + \\ (48469w_1^6 - 743244w_1^5w_2 + 4444890w_1^4w_2^2 - 13073120w_1^3w_2^3 + 19723200w_1^2w_2^4 - \\ 14608896w_1w_2^5 + 4205056w_2^6 + 1051200w_1^5w_3 - 21117960w_1^4w_2w_3 + 118315440w_1^3w_2^2w_3 - \\ 259434720w_1^2w_2^3w_3 + 249145920w_1w_2^4w_3 - 88047360w_2^5w_3 + 21207690w_1^4w_3^2 - \\ 76688640w_1^3w_2w_3^2 + 441644400w_1^2w_2^2w_3^2 - 809291520w_1w_2^3w_3^2 + 439814880w_2^4w_3^2 + \\ 815030640w_1^3w_3^3 - 868838400w_1^2w_2w_3^3 + 529260480w_1w_2^2w_3^3 - 665167680w_2^3w_3^3 + \\ 8239760640w_1^2w_3^4 - 10091148480w_1w_2w_3^4 + 4723181280w_2^2w_3^4 + 50452492032w_1w_3^5 - \\ 45443177280w_2w_3^5 + 145956110976w_3^6)Y_1^3 + \\ 6(43274w_1^6 - 667704w_1^5w_2 + 4028205w_1^4w_2^2 - 11998480w_1^3w_2^3 + 18432480w_1^2w_2^4 - \\ 13967616w_1w_2^5 + 4111616w_2^6 + 1805328w_1^5w_3 - 19051740w_1^4w_2w_3 + 75568680w_1^3w_2^2w_3 - \\ 143512560w_1^2w_2^3w_3 + 134468640w_1w_2^4w_3 - 49479552w_2^5w_3 + 23063490w_1^4w_3^2 - \\ 97315920w_1^3w_2w_3^2 + 127881720w_1^2w_2^2w_3^2 - 77891040w_1w_2^3w_3^2 + 44389440w_2^4w_3^2 + \\ 409273920w_1^3w_3^3 - 1276901280w_1^2w_2w_3^3 + 1247520960w_1w_2^2w_3^3 - 650304000w_2^3w_3^3 + \\ 3052420200w_1^2w_3^4 - 12009740400w_1w_2w_3^4 + 10525836600w_2^2w_3^4 + 26559204048w_1w_3^5 - \\ 45364278096w_2w_3^5 + 101461763160w_3^6)Y_1^4 + \\ 2(105901w_1^6 - 1755600w_1^5w_2 + 11616675w_1^4w_2^2 - 38949440w_1^3w_2^3 + 69116640w_1^2w_2^4 - \\ 60716544w_1w_2^5 + 19859200w_2^6 - 6063192w_1^5w_3 + 82652760w_1^4w_2w_3 - 412432200w_1^3w_2^2w_3 + \\ 898113600w_1^2w_2^3w_3 - 786384000w_1w_2^4w_3 + 220981248w_2^5w_3 - 146310030w_1^4w_3^2 + \\ 1634297040w_1^3w_2w_3^2 - 6098117400w_1^2w_2^2w_3^2 + 8499487680w_1w_2^3w_3^2 - 3567715200w_2^4w_3^2 + \end{aligned}$$



$$\begin{aligned}
& 880886880 w_1^3 w_3^3 - 8797422960 w_1^2 w_2 w_3^3 + 21711123360 w_1 w_2^2 w_3^3 - 19644958080 w_2^3 w_3^3 + \\
& 16531620480 w_1^2 w_3^4 - 146923709760 w_1 w_2 w_3^4 + 177054446160 w_2^2 w_3^4 + 281903296896 w_1 w_3^5 - \\
& 514798256592 w_2 w_3^5 + 1116497698848 w_3^6) Y_1^5 + \\
& 2(-60989 w_1^6 + 1013970 w_1^5 w_2 - 6509550 w_1^4 w_2^2 + 19983520 w_1^3 w_2^3 - 29628480 w_1^2 w_2^4 + \\
& 21098496 w_1 w_2^5 - 11333120 w_2^6 - 10304730 w_1^5 w_3 + 141141510 w_1^4 w_2 w_3 - 711217440 w_1^3 w_2^2 w_3 + \\
& 1580051520 w_1^2 w_2^3 w_3 - 1444147200 w_1 w_2^4 w_3 + 433497600 w_2^5 w_3 + 7910865 w_1^4 w_3^2 - \\
& 615917520 w_1^3 w_2 w_3^2 + 4948801920 w_1^2 w_2^2 w_3^2 - 12996339840 w_1 w_2^3 w_3^2 + 10198068480 w_2^4 w_3^2 + \\
& 6077720520 w_1^3 w_3^3 - 63144917280 w_1^2 w_2 w_3^3 + 192593116800 w_1 w_2^2 w_3^3 - 183921926400 w_2^3 w_3^3 + \\
& 58576870440 w_1^2 w_3^4 - 408475787760 w_1 w_2 w_3^4 + 511734905280 w_2^2 w_3^4 + 680591915568 w_1 w_3^5 - \\
& 892873018080 w_2 w_3^5 + 1956383853480 w_3^6) Y_1^6 + \\
& 9(14893 w_1^6 - 291876 w_1^5 w_2 + 2263200 w_1^4 w_2^2 - 8574080 w_1^3 w_2^3 + 15233280 w_1^2 w_2^4 - \\
& 7590912 w_1 w_2^5 - 6127616 w_2^6 + 1600536 w_1^5 w_3 - 28552200 w_1^4 w_2 w_3 + 197250240 w_1^3 w_2^2 w_3 - \\
& 650334720 w_1^2 w_2^3 w_3 + 995343360 w_1 w_2^4 w_3 - 529618944 w_2^5 w_3 + 150441690 w_1^4 w_3^2 - \\
& 1985417760 w_1^3 w_2 w_3^2 + 9512081280 w_1^2 w_2^2 w_3^2 - 19309670400 w_1 w_2^3 w_3^2 + 13599045120 w_2^4 w_3^2 + \\
& 2795123760 w_1^3 w_3^3 - 27544011840 w_1^2 w_2 w_3^3 + 87759394560 w_1 w_2^2 w_3^3 - 89221309440 w_2^3 w_3^3 + \\
& 10326191040 w_1^2 w_3^4 - 42916236480 w_1 w_2 w_3^4 + 17753904000 w_2^2 w_3^4 + 100671672960 w_1 w_3^5 + \\
& 161852587776 w_2 w_3^5 - 169776974592 w_3^6) Y_1^7 + \\
& 8(16238 w_1^6 - 291297 w_1^5 w_2 + 1928820 w_1^4 w_2^2 - 5038240 w_1^3 w_2^3 - 631680 w_1^2 w_2^4 + \\
& 26204928 w_1 w_2^5 - 34266112 w_2^6 + 2811906 w_1^5 w_3 - 49346640 w_1^4 w_2 w_3 + 339641280 w_1^3 w_2^2 w_3 - \\
& 1138037760 w_1^2 w_2^3 w_3 + 1835020800 w_1 w_2^4 w_3 - 1115172864 w_2^5 w_3 + 134656560 w_1^4 w_3^2 - \\
& 1881693720 w_1^3 w_2 w_3^2 + 9653294880 w_1^2 w_2^2 w_3^2 - 21377139840 w_1 w_2^3 w_3^2 + 17012160000 w_2^4 w_3^2 + \\
& 719438220 w_1^3 w_3^3 - 6652711440 w_1^2 w_2 w_3^3 + 19475976960 w_1 w_2^2 w_3^3 - 17504570880 w_2^3 w_3^3 - \\
& 24811199100 w_1^2 w_3^4 + 218557854720 w_1 w_2 w_3^4 - 442688885280 w_2^2 w_3^4 - 132608756664 w_1 w_3^5 + \\
& 960112922112 w_2 w_3^5 - 1682221113576 w_3^6) Y_1^8 + \\
& 108(-171 w_1^6 + 4104 w_1^5 w_2 - 41040 w_1^4 w_2^2 + 218880 w_1^3 w_2^3 - 656640 w_1^2 w_2^4 + 1050624 w_1 w_2^5 - \\
& 700416 w_2^6 + 47962 w_1^5 w_3 - 959240 w_1^4 w_2 w_3 + 7673920 w_1^3 w_2^2 w_3 - 30695680 w_1^2 w_2^3 w_3 + \\
& 61391360 w_1 w_2^4 w_3 - 49113088 w_2^5 w_3 - 394570 w_1^4 w_3^2 + 7285120 w_1^3 w_2 w_3^2 - 49542720 w_1^2 w_2^2 w_3^2 + \\
& 147665920 w_1 w_2^3 w_3^2 - 163217920 w_2^4 w_3^2 - 153027360 w_1^3 w_3^3 + 1762456320 w_1^2 w_2 w_3^3 - \\
& 6754337280 w_1 w_2^2 w_3^3 + 8611799040 w_2^3 w_3^3 - 2805037200 w_1^2 w_3^4 + 22007304000 w_1 w_2 w_3^4 - \\
& 43148620800 w_2^2 w_3^4 - 15525586656 w_1 w_3^5 + 67663897344 w_2 w_3^5 - 163361647584 w_3^6) Y_1^9 + \\
& 432 w_3(-6075 w_1^5 + 121500 w_1^4 w_2 - 972000 w_1^3 w_2^2 + 3888000 w_1^2 w_2^3 - 7776000 w_1 w_2^4 + \\
& 6220800 w_2^5 - 617700 w_1^4 w_3 + 9883200 w_1^3 w_2 w_3 - 59299200 w_1^2 w_2^2 w_3 + 158131200 w_1 w_2^3 w_3 - \\
& 158131200 w_2^4 w_3 - 20973400 w_1^3 w_3^2 + 251680800 w_1^2 w_2 w_3^2 - 1006723200 w_1 w_2^2 w_3^2 + \\
& 1342297600 w_2^3 w_3^2 - 201064680 w_1^2 w_3^3 + 1584217440 w_1 w_2 w_3^3 - 3119834880 w_2^2 w_3^3 - \\
& 1374418368 w_1 w_3^4 + 4253513472 w_2 w_3^4 - 25378959744 w_3^5) Y_1^{10} + \\
& 23328000 w_3^3(-5 w_1^3 + 60 w_1^2 w_2 - 240 w_1 w_2^2 + 320 w_2^3 + 575 w_1^2 w_3 - 4600 w_1 w_2 w_3 + 9200 w_2^2 w_3 - \\
& 1202 w_1 w_3^2 + 4808 w_2 w_3^2 - 124918 w_3^3) Y_1^{11} + \\
& 2332800000 w_3^5(3 w_1 - 12 w_2 + 176 w_3) Y_1^{12} + \\
& 139968 w_3(23 w_1^5 - 145 w_1^4 w_2 + 395 w_1^3 w_2^2 - 545 w_1^2 w_2^3 + 370 w_1 w_2^4 - 98 w_2^5 + 620 w_1^4 w_3 - \\
& 2960 w_1^3 w_2 w_3 + 5880 w_1^2 w_2^2 w_3 - 5360 w_1 w_2^3 w_3 + 1820 w_2^4 w_3 + 6855 w_1^3 w_3^2 - 23355 w_1^2 w_2 w_3^2 + \\
& 30060 w_1 w_2^2 w_3^2 - 13560 w_2^3 w_3^2 + 39120 w_1^2 w_3^3 - 84720 w_1 w_2 w_3^3 + 53160 w_2^2 w_3^3 + 112674 w_1 w_3^4 - \\
& 114654 w_2 w_3^4 + 126636 w_3^5) Y_2 + \\
& 1944(115 w_1^6 - 1470 w_1^5 w_2 + 7245 w_1^4 w_2^2 - 17660 w_1^3 w_2^3 + 22890 w_1^2 w_2^4 - 15120 w_1 w_2^5 + \\
& 4000 w_2^6 + 5770 w_1^5 w_3 - 61550 w_1^4 w_2 w_3 + 239260 w_1^3 w_2^2 w_3 - 421840 w_1^2 w_2^3 w_3 +
\end{aligned}$$

$$\begin{aligned}
& 348920 w_1 w_2^4 w_3 - 110560 w_2^5 w_3 + 48435 w_1^4 w_3^2 - 503340 w_1^3 w_2 w_3^2 + 1702980 w_1^2 w_2^2 w_3^2 - \\
& 2154480 w_1 w_2^3 w_3^2 + 922200 w_2^4 w_3^2 + 7020 w_1^3 w_3^3 - 1034640 w_1^2 w_2 w_3^3 + 3974400 w_1 w_2^2 w_3^3 - \\
& 3058560 w_2^3 w_3^3 - 540630 w_1^2 w_3^4 + 1480320 w_1 w_2 w_3^4 + 2145600 w_2^2 w_3^4 - 2453976 w_1 w_3^5 + \\
& 11969856 w_2 w_3^5 - 10703448 w_3^6) Y_1 Y_2 + \\
& 54 (2599 w_1^6 - 44844 w_1^5 w_2 + 308760 w_1^4 w_2^2 - 1072400 w_1^3 w_2^3 + 1947840 w_1^2 w_2^4 - 1718016 w_1 w_2^5 + \\
& 572416 w_2^6 + 59640 w_1^5 w_3 - 638760 w_1^4 w_2 w_3 + 2621760 w_1^3 w_2^2 w_3 - 5860320 w_1^2 w_2^3 w_3 + \\
& 8221440 w_1 w_2^4 w_3 - 4462080 w_2^5 w_3 - 4497300 w_1^4 w_3^2 + 25488720 w_1^3 w_2 w_3^2 - 63547200 w_1^2 w_2^2 w_3^2 + \\
& 81408960 w_1 w_2^3 w_3^2 - 34493760 w_2^4 w_3^2 - 49122720 w_1^3 w_3^3 + 119024640 w_1^2 w_2 w_3^3 - \\
& 177344640 w_1 w_2^2 w_3^3 + 143251200 w_2^3 w_3^3 - 169406640 w_1^2 w_3^4 + 94141440 w_1 w_2 w_3^4 - \\
& 51114240 w_2^2 w_3^4 - 2646390528 w_1 w_3^5 + 3953878272 w_2 w_3^5 - 11168202240 w_3^6) Y_1^2 Y_2 + \\
& 216 (1364 w_1^6 - 24924 w_1^5 w_2 + 182055 w_1^4 w_2^2 - 670960 w_1^3 w_2^3 + 1287840 w_1^2 w_2^4 - 1180416 w_1 w_2^5 + \\
& 386816 w_2^6 - 112209 w_1^5 w_3 + 1371000 w_1^4 w_2 w_3 - 6057780 w_1^3 w_2^2 w_3 + 11491440 w_1^2 w_2^3 w_3 - \\
& 8506560 w_1 w_2^4 w_3 + 1787136 w_2^5 w_3 - 739980 w_1^4 w_3^2 + 8601300 w_1^3 w_2 w_3^2 - 39277980 w_1^2 w_2^2 w_3^2 + \\
& 79248960 w_1 w_2^3 w_3^2 - 49596480 w_2^4 w_3^2 + 16542000 w_1^3 w_3^3 - 94577760 w_1^2 w_2 w_3^3 + \\
& 143164800 w_1 w_2^2 w_3^3 - 46719360 w_2^3 w_3^3 - 150956460 w_1^2 w_3^4 + 1419294960 w_1 w_2 w_3^4 - \\
& 1405900800 w_2^2 w_3^4 - 5687381952 w_1 w_3^5 + 12549919680 w_2 w_3^5 - 19620243792 w_3^6) Y_1^3 Y_2 + \\
& 72 (-2891 w_1^6 + 50565 w_1^5 w_2 - 353910 w_1^4 w_2^2 + 1272640 w_1^3 w_2^3 - 2556480 w_1^2 w_2^4 + 3005184 w_1 w_2^5 - \\
& 1902080 w_2^6 - 3213 w_1^5 w_3 + 371250 w_1^4 w_2 w_3 - 4551120 w_1^3 w_2^2 w_3 + 21029760 w_1^2 w_2^3 w_3 - \\
& 40711680 w_1 w_2^4 w_3 + 25892352 w_2^5 w_3 + 3293460 w_1^4 w_3^2 - 25503660 w_1^3 w_2 w_3^2 + 30025080 w_1^2 w_2^2 w_3^2 + \\
& 140849280 w_1 w_2^3 w_3^2 - 254689920 w_2^4 w_3^2 - 209584800 w_1^3 w_3^3 + 2077646760 w_1^2 w_2 w_3^3 - \\
& 5395870080 w_1 w_2^2 w_3^3 + 3266213760 w_2^3 w_3^3 - 7122067560 w_1^2 w_3^4 + 42669011520 w_1 w_2 w_3^4 - \\
& 50244896160 w_2^2 w_3^4 - 80632905408 w_1 w_3^5 + 157744600848 w_2 w_3^5 - 205830059040 w_3^6) Y_1^4 Y_2 + \\
& 36 (-8317 w_1^6 + 192264 w_1^5 w_2 - 1849200 w_1^4 w_2^2 + 9470720 w_1^3 w_2^3 - 27237120 w_1^2 w_2^4 + \\
& 41699328 w_1 w_2^5 - 26546176 w_2^6 - 616698 w_1^5 w_3 + 10877580 w_1^4 w_2 w_3 - 75369600 w_1^3 w_2^2 w_3 + \\
& 254874240 w_1^2 w_2^3 w_3 - 416540160 w_1 w_2^4 w_3 + 258665472 w_2^5 w_3 - 44838090 w_1^4 w_3^2 + \\
& 611442000 w_1^3 w_2 w_3^2 - 3022349760 w_1^2 w_2^2 w_3^2 + 6308133120 w_1 w_2^3 w_3^2 - 4528673280 w_2^4 w_3^2 - \\
& 1959857640 w_1^3 w_3^3 + 17825326560 w_1^2 w_2 w_3^3 - 50015905920 w_1 w_2^2 w_3^3 + 40289287680 w_2^3 w_3^3 - \\
& 29008110420 w_1^2 w_3^4 + 165895011360 w_1 w_2 w_3^4 - 207940737600 w_2^2 w_3^4 - 203187801456 w_1 w_3^5 + \\
& 216601319808 w_2 w_3^5 - 165452312016 w_3^6) Y_1^5 Y_2 + \\
& 18 (-22553 w_1^6 + 541272 w_1^5 w_2 - 5412720 w_1^4 w_2^2 + 28867840 w_1^3 w_2^3 - 86603520 w_1^2 w_2^4 + \\
& 138565632 w_1 w_2^5 - 92377088 w_2^6 - 2102004 w_1^5 w_3 + 37238400 w_1^4 w_2 w_3 - 259493760 w_1^3 w_2^2 w_3 + \\
& 884321280 w_1^2 w_2^3 w_3 - 1461335040 w_1 w_2^4 w_3 + 923222016 w_2^5 w_3 - 132943140 w_1^4 w_3^2 + \\
& 1818849600 w_1^3 w_2 w_3^2 - 9063653760 w_1^2 w_2^2 w_3^2 + 19237893120 w_1 w_2^3 w_3^2 - 14306042880 w_2^4 w_3^2 - \\
& 2851640640 w_1^3 w_3^3 + 28046165760 w_1^2 w_2 w_3^3 - 87490575360 w_1 w_2^2 w_3^3 + 83728650240 w_2^3 w_3^3 - \\
& 2171428560 w_1^2 w_3^4 - 43665816960 w_1 w_2 w_3^4 + 185471596800 w_2^2 w_3^4 + 257835472704 w_1 w_3^5 - \\
& 1771101884160 w_2 w_3^5 + 2944386630144 w_3^6) Y_1^6 Y_2 + \\
& 324 (-513 w_1^6 + 12312 w_1^5 w_2 - 123120 w_1^4 w_2^2 + 656640 w_1^3 w_2^3 - 1969920 w_1^2 w_2^4 + 3151872 w_1 w_2^5 - \\
& 2101248 w_2^6 - 3772 w_1^5 w_3 + 75440 w_1^4 w_2 w_3 - 603520 w_1^3 w_2^2 w_3 + 2414080 w_1^2 w_2^3 w_3 - \\
& 4828160 w_1 w_2^4 w_3 + 3862528 w_2^5 w_3 - 753840 w_1^4 w_3^2 + 12061440 w_1^3 w_2 w_3^2 - 72368640 w_1^2 w_2^2 w_3^2 + \\
& 192983040 w_1 w_2^3 w_3^2 - 192983040 w_2^4 w_3^2 + 79293120 w_1^3 w_3^3 - 934669440 w_1^2 w_2 w_3^3 + \\
& 3671285760 w_1 w_2^2 w_3^3 - 4805191680 w_2^3 w_3^3 + 3802053600 w_1^2 w_3^4 - 29682374400 w_1 w_2 w_3^4 + \\
& 57896640000 w_2^2 w_3^4 + 40244283648 w_1 w_3^5 - 167285958912 w_2 w_3^5 + 282944332224 w_3^6) Y_1^7 Y_2 + \\
& 7776 w_3^2 (59025 w_1^4 - 944400 w_1^3 w_2 + 5666400 w_1^2 w_2^2 - 15110400 w_1 w_2^3 + 15110400 w_2^4 + \\
& 4227400 w_1^3 w_3 - 50728800 w_1^2 w_2 w_3 + 202915200 w_1 w_2^2 w_3 - 270553600 w_2^3 w_3 + 86580180 w_1^2 w_3^2 - \\
& 692641440 w_1 w_2 w_3^2 + 1385282880 w_2^2 w_3^2 + 747341424 w_1 w_3^3 - 2882445696 w_2 w_3^3 +
\end{aligned}$$

$$\begin{aligned}
& 7024285008 w_3^4) Y_1^8 Y_2 + \\
& 69984000 w_3^4 (-185 w_1^2 + 1480 w_1 w_2 - 2960 w_2^2 - 1812 w_1 w_3 + 7248 w_2 w_3 + 171480 w_3^2) Y_1^9 Y_2 - \\
& 629856000000 w_3^6 Y_1^{10} Y_2 + \\
& 17496 (-35 w_1^6 + 510 w_1^5 w_2 - 2835 w_1^4 w_2^2 + 7480 w_1^3 w_2^3 - 9600 w_1^2 w_2^4 + 5760 w_1 w_2^5 - 1280 w_2^6 - \\
& 1370 w_1^5 w_3 + 16810 w_1^4 w_2 w_3 - 75680 w_1^3 w_2^2 w_3 + 151760 w_1^2 w_2^3 w_3 - 131200 w_1 w_2^4 w_3 + \\
& 39680 w_2^5 w_3 - 8925 w_1^4 w_3^2 + 114360 w_1^3 w_2 w_3^2 - 479880 w_1^2 w_2^2 w_3^2 + 751200 w_1 w_2^3 w_3^2 - \\
& 360960 w_2^4 w_3^2 + 14400 w_1^3 w_3^3 + 108000 w_1^2 w_2 w_3^3 - 855360 w_1 w_2^2 w_3^3 + 1005120 w_2^3 w_3^3 + \\
& 159480 w_1^2 w_3^4 - 835200 w_1 w_2 w_3^4 + 147600 w_2^2 w_3^4 + 363744 w_1 w_3^5 - 2877984 w_2 w_3^5 + 873072 w_3^6) Y_2^2 + \\
& 972 (109 w_1^6 - 2004 w_1^5 w_2 + 14730 w_1^4 w_2^2 - 54560 w_1^3 w_2^3 + 104640 w_1^2 w_2^4 - 93696 w_1 w_2^5 + 27136 w_2^6 + \\
& 840 w_1^5 w_3 - 14640 w_1^4 w_2 w_3 + 99840 w_1^3 w_2^2 w_3 - 330240 w_1^2 w_2^3 w_3 + 522240 w_1 w_2^4 w_3 - 307200 w_2^5 w_3 + \\
& 694170 w_1^4 w_3^2 - 6246720 w_1^3 w_2 w_3^2 + 17943120 w_1^2 w_2^2 w_3^2 - 17210880 w_1 w_2^3 w_3^2 + 3836160 w_2^4 w_3^2 - \\
& 2509920 w_1^3 w_3^3 - 2617920 w_1^2 w_2 w_3^3 + 53239680 w_1 w_2^2 w_3^3 - 35631360 w_2^3 w_3^3 - 80261280 w_1^2 w_3^4 + \\
& 220086720 w_1 w_2 w_3^4 - 40214880 w_2^2 w_3^4 + 17200512 w_1 w_3^5 - 151103232 w_2 w_3^5 + 977186592 w_3^6) Y_1 Y_2^2 + \\
& 1944 (-91 w_1^6 + \\
& 1986 w_1^5 w_2 - 17880 w_1^4 w_2^2 + 84800 w_1^3 w_2^3 - 222720 w_1^2 w_2^4 + 305664 w_1 w_2^5 - 169984 w_2^6 + \\
& 28158 w_1^5 w_3 - 381450 w_1^4 w_2 w_3 + 1889520 w_1^3 w_2^2 w_3 - 4076160 w_1^2 w_2^3 w_3 + 3521280 w_1 w_2^4 w_3 - \\
& 978432 w_2^5 w_3 + 328455 w_1^4 w_3^2 - 3800520 w_1^3 w_2 w_3^2 + 14949360 w_1^2 w_2^2 w_3^2 - 21254400 w_1 w_2^3 w_3^2 + \\
& 4976640 w_2^4 w_3^2 - 3570480 w_1^3 w_3^3 + 17496000 w_1^2 w_2 w_3^3 - 14074560 w_1 w_2^2 w_3^3 + 4872960 w_2^3 w_3^3 + \\
& 192588840 w_1^2 w_3^4 - 980372160 w_1 w_2 w_3^4 + 793877760 w_2^2 w_3^4 + 2910144672 w_1 w_3^5 - \\
& 6921487584 w_2 w_3^5 + 6774579504 w_3^6) Y_1^2 Y_2^2 + \\
& 324 (-1777 w_1^6 + 38868 w_1^5 w_2 - 350880 w_1^4 w_2^2 + 1669760 w_1^3 w_2^3 - 4404480 w_1^2 w_2^4 + 6079488 w_1 w_2^5 - \\
& 3407872 w_2^6 + 5328 w_1^5 w_3 - 145440 w_1^4 w_2 w_3 + 1474560 w_1^3 w_2^2 w_3 - 7142400 w_1^2 w_2^3 w_3 + \\
& 16773120 w_1 w_2^4 w_3 - 15409152 w_2^5 w_3 - 1347030 w_1^4 w_3^2 + 20308320 w_1^3 w_2 w_3^2 - 114384960 w_1^2 w_2^2 w_3^2 + \\
& 285120000 w_1 w_2^3 w_3^2 - 265213440 w_2^4 w_3^2 + 141959520 w_1^3 w_3^3 - 1184232960 w_1^2 w_2 w_3^3 + \\
& 2735389440 w_1 w_2^2 w_3^3 - 1079239680 w_2^3 w_3^3 + 5862792960 w_1^2 w_3^4 - 31292412480 w_1 w_2 w_3^4 + \\
& 34249703040 w_2^2 w_3^4 + 54443259648 w_1 w_3^5 - 118126460160 w_2 w_3^5 + 73898741664 w_3^6) Y_1^3 Y_2^2 + \\
& 17496 (19 w_1^6 - 456 w_1^5 w_2 + 4560 w_1^4 w_2^2 - 24320 w_1^3 w_2^3 + 72960 w_1^2 w_2^4 - 116736 w_1 w_2^5 + \\
& 77824 w_2^6 + 676 w_1^5 w_3 - 13520 w_1^4 w_2 w_3 + 108160 w_1^3 w_2^2 w_3 - 432640 w_1^2 w_2^3 w_3 + \\
& 865280 w_1 w_2^4 w_3 - 692224 w_2^5 w_3 + 47970 w_1^4 w_3^2 - 601920 w_1^3 w_2 w_3^2 + 2617920 w_1^2 w_2^2 w_3^2 - \\
& 4331520 w_1 w_2^3 w_3^2 + 1681920 w_2^4 w_3^2 + 4475280 w_1^3 w_3^3 - 40034880 w_1^2 w_2 w_3^3 + \\
& 105465600 w_1 w_2^2 w_3^3 - 67722240 w_2^3 w_3^3 + 91656720 w_1^2 w_3^4 - 536624640 w_1 w_2 w_3^4 + \\
& 679991040 w_2^2 w_3^4 + 118710144 w_1 w_3^5 + 969224832 w_2 w_3^5 - 3877256160 w_3^6) Y_1^4 Y_2^2 + \\
& 23328 w_3 (199 w_1^5 - 3980 w_1^4 w_2 + 31840 w_1^3 w_2^2 - 127360 w_1^2 w_2^3 + 254720 w_1 w_2^4 - 203776 w_2^5 - \\
& 5760 w_1^4 w_3 + 92160 w_1^3 w_2 w_3 - 552960 w_1^2 w_2^2 w_3 + 1474560 w_1 w_2^3 w_3 - 1474560 w_2^4 w_3 + 43200 w_1^3 w_3^2 - \\
& 518400 w_1^2 w_2 w_3^2 + 2073600 w_1 w_2^2 w_3^2 - 2764800 w_2^3 w_3^2 - 48576240 w_1^2 w_3^3 + 381028320 w_1 w_2 w_3^3 - \\
& 746893440 w_2^2 w_3^3 - 1294311312 w_1 w_3^4 + 5203372608 w_2 w_3^4 - 8225713296 w_3^5) Y_1^5 Y_2^2 + \\
& 1119744 w_3^3 (-25775 w_1^3 + 309300 w_1^2 w_2 - 1237200 w_1 w_2^2 + 1649600 w_2^3 - 1342845 w_1^2 w_3 + \\
& 10742760 w_1 w_2 w_3 - 21485520 w_2^2 w_3 - 18125964 w_1 w_3^2 + 72503856 w_2 w_3^2 - 114744438 w_3^3) Y_1^6 Y_2^2 + \\
& 15116544000 (11 w_1 - 44 w_2 - 972 w_3) w_3^5 Y_1^7 Y_2^2 + \\
& 5832 (-41 w_1^6 + 876 w_1^5 w_2 - 7680 w_1^4 w_2^2 + 35200 w_1^3 w_2^3 - 88320 w_1^2 w_2^4 + 113664 w_1 w_2^5 - 57344 w_2^6 - \\
& 134460 w_1^4 w_3^2 + 1257120 w_1^3 w_2 w_3^2 - 3576960 w_1^2 w_2^2 w_3^2 + 2695680 w_1 w_2^3 w_3^2 + 414720 w_2^4 w_3^2 + \\
& 432000 w_1^3 w_3^3 + 1347840 w_1^2 w_2 w_3^3 - 14722560 w_1 w_2^2 w_3^3 + 9676800 w_2^3 w_3^3 + 13433040 w_1^2 w_3^4 - \\
& 40979520 w_1 w_2 w_3^4 - 622080 w_2^2 w_3^4 - 5971968 w_1 w_3^5 + 17169408 w_2 w_3^5 - 148039488 w_3^6) Y_2^3 +
\end{aligned}$$

$$\begin{aligned}
& 629856 w_3 (59 w_1^5 - 1090 w_1^4 w_2 + 8000 w_1^3 w_2^2 - 29120 w_1^2 w_2^3 + 52480 w_1 w_2^4 - 37376 w_2^5 + \\
& 920 w_1^4 w_3 - 14000 w_1^3 w_2 w_3 + 79680 w_1^2 w_2^2 w_3 - 200960 w_1 w_2^3 w_3 + 189440 w_2^4 w_3 + 5000 w_1^3 w_3^2 - \\
& 57840 w_1^2 w_2 w_3^2 + 222720 w_1 w_2^2 w_3^2 - 285440 w_2^3 w_3^2 - 1573920 w_1^2 w_3^3 + 7381440 w_1 w_2 w_3^3 - \\
& 5276160 w_2^2 w_3^3 - 17284752 w_1 w_3^4 + 45336672 w_2 w_3^4 - 29611008 w_3^5) Y_1 Y_2^3 + \\
& 52488 (19 w_1^6 - 456 w_1^5 w_2 + 4560 w_1^4 w_2^2 - 24320 w_1^3 w_2^3 + 72960 w_1^2 w_2^4 - 116736 w_1 w_2^5 + \\
& 77824 w_2^6 + 244 w_1^5 w_3 - 4880 w_1^4 w_2 w_3 + 39040 w_1^3 w_2^2 w_3 - 156160 w_1^2 w_2^3 w_3 + 312320 w_1 w_2^4 w_3 - \\
& 249856 w_2^5 w_3 + 18300 w_1^4 w_3^2 - 292800 w_1^3 w_2 w_3^2 + 1756800 w_1^2 w_2^2 w_3^2 - 4684800 w_1 w_2^3 w_3^2 + \\
& 4684800 w_2^4 w_3^2 - 708960 w_1^3 w_3^3 + 5863680 w_1^2 w_2 w_3^3 - 12879360 w_1 w_2^2 w_3^3 + 3072000 w_2^3 w_3^3 - \\
& 36048240 w_1^2 w_3^4 + 196007040 w_1 w_2 w_3^4 - 207256320 w_2^2 w_3^4 - 308650176 w_1 w_3^5 + 759456000 w_2 w_3^5 + \\
& 447550272 w_3^6) Y_1^2 Y_2^3 + \\
& 1259712 w_3^2 (125 w_1^4 - 2000 w_1^3 w_2 + 12000 w_1^2 w_2^2 - 32000 w_1 w_2^3 + 32000 w_2^4 + 1000 w_1^3 w_3 - \\
& 12000 w_1^2 w_2 w_3 + 48000 w_1 w_2^2 w_3 - 64000 w_2^3 w_3 - 60840 w_1^2 w_3^2 + 486720 w_1 w_2 w_3^2 - \\
& 973440 w_2^2 w_3^2 + 17240256 w_1 w_3^3 - 68416704 w_2 w_3^3 + 152725392 w_3^4) Y_1^3 Y_2^3 + \\
& 90699264 w_3^4 (11215 w_1^2 - 89720 w_1 w_2 + 179440 w_2^2 + 329676 w_1 w_3 - 1318704 w_2 w_3 + \\
& 1922532 w_3^2) Y_1^4 Y_2^3 + \\
& 5714053632000 w_3^6 Y_1^5 Y_2^3 + \\
& 45349632 w_3 (-w_1^5 + 20 w_1^4 w_2 - 160 w_1^3 w_2^2 + 640 w_1^2 w_2^3 - 1280 w_1 w_2^4 + 1024 w_2^5 - 10 w_1^4 w_3 + \\
& 160 w_1^3 w_2 w_3 - 960 w_1^2 w_2^2 w_3 + 2560 w_1 w_2^3 w_3 - 2560 w_2^4 w_3 - 40 w_1^3 w_3^2 + 480 w_1^2 w_2 w_3^2 - \\
& 1920 w_1 w_2^2 w_3^2 + 2560 w_2^3 w_3^2 + 16560 w_1^2 w_3^3 - 80640 w_1 w_2 w_3^3 + 57600 w_2^2 w_3^3 + 168048 w_1 w_3^4 - \\
& 485568 w_2 w_3^4 + 235872 w_3^5) Y_2^4 - \\
& 81498730659840 w_3^6 Y_1 Y_2^4 + \\
& 78364164096 (-201 w_1 + 804 w_2 - 1682 w_3) w_3^5 Y_1^2 Y_2^4 + \\
& 45137758519296 w_3^6 Y_2^5.
\end{aligned}$$

In the case of the special 1-parameter family, the coordinate change

$$\begin{aligned}
3 w_1 & \rightarrow w_2 \\
18 w_2 & \rightarrow (3 w_1 + 5 w_2 - 5 w_3) \\
12 w_3 & \rightarrow (w_1 - w_3)
\end{aligned}$$

yields:

$$F_V(w) = \frac{F(\sigma_z(w))|_{\{F(z)=0\}}}{3^{18} \Phi(z)^2 \Psi(z)^9} =$$

$$\begin{aligned}
& 9 w_2^5 (-24 w_1 + 19 w_2 + 24 w_3) + \\
& w_2^4 (365 w_1^2 - 54 w_1 w_2 - 171 w_2^2 - 640 w_1 w_3 - 16 w_2 w_3 + 275 w_3^2) V + \\
& (128 w_1^6 - 328 w_1^5 w_2 + 165 w_1^4 w_2^2 - 20 w_1^3 w_2^3 - 365 w_1^2 w_2^4 + 270 w_1 w_2^5 - 768 w_1^5 w_3 + 2120 w_1^4 w_2 w_3 - \\
& 2320 w_1^3 w_2^2 w_3 + 2000 w_1^2 w_2^3 w_3 - 360 w_1 w_2^4 w_3 + 1920 w_1^4 w_3^2 - 5200 w_1^3 w_2 w_3^2 + 5850 w_1^2 w_2^2 w_3^2 - \\
& 3800 w_1 w_2^3 w_3^2 + 675 w_2^4 w_3^2 - 2560 w_1^3 w_3^3 + 6160 w_1^2 w_2 w_3^3 - 5400 w_1 w_2^2 w_3^3 + 1800 w_2^3 w_3^3 + 1920 w_1^2 w_3^4 - \\
& 3560 w_1 w_2 w_3^4 + 1705 w_2^2 w_3^4 - 768 w_1 w_3^5 + 808 w_2 w_3^5 + 128 w_3^6) V^2 + \\
& 5 (-25 w_1^6 + 66 w_1^5 w_2 - 33 w_1^4 w_2^2 + 4 w_1^3 w_2^3 + 112 w_1^5 w_3 - 224 w_1^4 w_2 w_3 + 64 w_1^3 w_2^2 w_3 - 193 w_1^4 w_3^2 + \\
& 280 w_1^3 w_2 w_3^2 - 30 w_1^2 w_2^2 w_3^2 + 152 w_1^3 w_3^3 - 152 w_1^2 w_2 w_3^3 - 43 w_1^2 w_3^4 + 30 w_1 w_2 w_3^4 - 8 w_1 w_3^5 + 5 w_3^6) V^3 +
\end{aligned}$$

$$w_1^4 (-3 w_1^2 - 2 w_1 w_2 + 8 w_1 w_3 - 5 w_3^2) V^4.$$

## B.2 Root-selectors

In the coordinates given above, the expressions for the respective factors in the root-selecting functions  $\overline{J}_Y(w)$  and  $\overline{K}_V(w)$  take the forms

$$\begin{aligned} \overline{\Gamma}_Y(w) &= \frac{5184 F(\tau_z(w))}{(39 - \sqrt{15} i) F(z)^{42}} = \\ &((3 + \sqrt{15} i) w_1^{10} - 10 i (-3 i + \sqrt{15}) w_1^9 w_2 + 45 (3 + \sqrt{15} i) w_1^8 w_2^2 - 120 i (-3 i + \sqrt{15}) w_1^7 w_2^3 + 210 (3 + \sqrt{15} i) w_1^6 w_2^4 - 252 i (-3 i + \sqrt{15}) w_1^5 w_2^5 + 210 (3 + \sqrt{15} i) w_1^4 w_2^6 - 120 i (-3 i + \sqrt{15}) w_1^3 w_2^7 + 45 (3 + \sqrt{15} i) w_1^2 w_2^8 - 10 i (-3 i + \sqrt{15}) w_1 w_2^9 + (3 + \sqrt{15} i) w_2^{10} + 10 (27 + 5 \sqrt{15} i) w_1^9 w_3 - 90 i (-27 i + 5 \sqrt{15}) w_1^8 w_2 w_3 + 360 (27 + 5 \sqrt{15} i) w_1^7 w_2^2 w_3 - 840 i (-27 i + 5 \sqrt{15}) w_1^6 w_2^3 w_3 + \\ &+ \dots + \end{aligned}$$

$$4608 i (5 (9 i + 17 \sqrt{15}) w_1^2 w_3^8 - 40 i (9 - 17 \sqrt{15} i) w_1 w_2 w_3^8 + 80 (9 i + 17 \sqrt{15}) w_2^2 w_3^8 + 60 (15 i + 7 \sqrt{15}) w_1 w_3^9 - 240 i (15 - 7 \sqrt{15} i) w_2 w_3^9 + 36 (9 i + 17 \sqrt{15}) w_3^{10}) Y_2^8 + 663552 (3 + 5 \sqrt{15} i) w_3^{10} Y_1 Y_2^8.$$

$$\overline{\Theta}_V(w) = \frac{2^7 3^4 \overline{\Theta}_z(w)|_{\{F(z)=0\}}}{(39 - \sqrt{15} i) \Phi(z) \Psi(z)^{16}} =$$

$$1944 w_2^{10} - i (12960 (-3 i + \sqrt{15}) w_1^3 w_2^7 + 135 i (765 + 203 \sqrt{15} i) w_1^2 w_2^8 + 90 (-1033 i + 228 \sqrt{15}) w_1 w_2^9 + i (18837 + 5500 \sqrt{15} i) w_2^{10} + 38880 i (3 + \sqrt{15} i) w_1^2 w_2^7 w_3 + 270 (-693 i + 227 \sqrt{15}) w_1 w_2^8 w_3 + 450 i (167 + 57 \sqrt{15} i) w_2^9 w_3 + 38880 (-3 i + \sqrt{15}) w_1 w_2^7 w_3^2 + 135 i (621 + 251 \sqrt{15} i) w_2^8 w_3^2 + 12960 i (3 + \sqrt{15} i) w_2^7 w_3^3) V +$$

$$+ \dots +$$

$$5 (2 (46 + \sqrt{15} i) w_1^{10} + 2 i (38 i + \sqrt{15}) w_1^9 w_2 + (8 + \sqrt{15} i) w_1^8 w_2^2 - 2 i (-171 i + 5 \sqrt{15}) w_1^9 w_3 + 6 (25 - \sqrt{15} i) w_1^8 w_2 w_3 + (473 + 19 \sqrt{15} i) w_1^8 w_3^2 + 4 i (19 i + \sqrt{15}) w_1^7 w_2 w_3^2 - 16 i (-18 i + \sqrt{15}) w_1^7 w_3^3 + 5 (13 + \sqrt{15} i) w_1^6 w_3^4) V^7 + w_1^{10} V^8.$$

## B.3 Degree 19 maps

Using (2), the identities

$$\begin{aligned} \alpha &= \frac{1}{2^{10} 3^{12}} \\ |\tau_z|^2 &= F(z)^{25} T_Y, \\ |\sigma_z|^2 &= -2^8 3^{17} \Phi(z)^2 \Psi(z)^9 V^2 (V - 1), \end{aligned}$$

and the results in Section 4.3, the following procedures yield the conic-fixing 19-maps. For the 2-parameter case, substitution gives

$$\begin{aligned}
f_{64}(y) &= \frac{(\alpha F(z)^{25})^{17}}{|\tau_z|^{13}} \tau_z \left\{ \left[ 10 \left( \frac{T_Y}{\alpha} \right)^4 F_Y^6 \Phi_Y + 100 \left( \frac{T_Y}{\alpha} \right)^3 F_Y^4 \Phi_Y^2 + \right. \right. \\
&\quad 45 \left( \frac{T_Y}{\alpha} \right)^2 F_Y^2 \Phi_Y^3 + 156 \left( \frac{T_Y}{\alpha} \right) \Phi_Y^4 + 39 \left( \frac{T_Y}{\alpha} \right)^2 F_Y^3 \Psi_Y + \\
&\quad \left. \left. 51 \left( \frac{T_Y}{\alpha} \right) F_Y \Phi_Y \Psi_Y \right] \cdot \psi_Y(w) - 27 \Psi_Y \cdot \phi_Y(w) + 54 \Phi_Y^2 \cdot f_Y(w) \right\} \\
&= \frac{\alpha F(z)^{25}}{|\tau_z|^{13}} \tau_z (f_{64})_Y.
\end{aligned}$$

Thus,

$$\begin{aligned}
f_{19}(y) &= \frac{f_{64}(y)}{X(y)} \\
&= \frac{(\alpha F(z)^{25})^5}{|\tau_z|^4} \tau_z \frac{(f_{64})_Y(w)}{X_Y(w)} \\
&= \frac{(\alpha F(z)^{25})^5}{|\tau_z|^4} \tau_z (f_{19})_Y(w).
\end{aligned}$$

As for the special 19-map, let  $f_Y(w) = (f_{19})_Y(w)$  so that

$$\begin{aligned}
h_{19}(y) &= 1620 F(y)^3 \cdot y + f_{19}(y) \\
&= 1620 (\alpha F(z)^{25} F_Y(w))^3 \cdot \tau_z(w) + \frac{(\alpha F(z)^{25})^5}{|\tau_z|^4} \tau_z f_Y(w) \\
&= \frac{(\alpha F(z)^{25})^5}{|\tau_z|^4} \tau_z \left( 1620 \frac{|\tau_z|^4}{(\alpha F(z)^{25})^2} F_Y(w)^3 \cdot w + f_Y(w) \right) \\
&= \frac{(\alpha F(z)^{25})^5}{|\tau_z|^4} \tau_z \left( 1620 \left( \frac{T_Y}{\alpha} \right)^2 F_Y(w)^3 \cdot w + f_Y(w) \right).
\end{aligned}$$

The family of conic-fixing maps on  $\mathbf{CP}^2$  is

$$\begin{aligned}
h_Y(w) &= 1620 \left( \frac{T_Y}{\alpha} \right)^2 F_Y(w)^3 \cdot w + f_Y(w) \\
&= (2^{12} \cdot 3^{16} \cdot 5) T_Y^2 F_Y(w)^3 \cdot w + f_Y(w)
\end{aligned}$$

Similar calculations in the 1-parameter setting yield

$$\begin{aligned}
h_V(w) &= 1620 \left( -\frac{64}{9} \right) V^4 (V-1)^2 F_V(w)^3 \cdot w + f_V(w) \\
&= 11520 V^4 (V-1)^2 F_V(w)^3 \cdot w + f_V(w).
\end{aligned}$$

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